

SPACE DEPENDENT MODEL
FOR THE SLOWING DOWN OF FAST NEUTRONS

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THESIS

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THE SLOWING DOWN OF FAST NEUTRONS

by

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June 1971

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the Slowing Down of Fast Neutrons

by

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Submitted in partial fulfillment of the
requirements for the degree of

MASTER OF SCIENCE IN PHYSICS

from the

NAVAL POSTGRADUATE SCHOOL
June 1971

10/15/19

ABSTRACT

The slowing down of fast neutrons was analyzed by a multi-group method of discrete time and energy states coupled with a spatial harmonic expansion method to determine the neutron density in a homogeneous, isotropically scattering slab. Five neutron source geometries were studied for both a fissioning and a non-fissioning system.

Numerical results were obtained for the neutron flux, mean neutron energy and the neutron spectra for the one dimensional system using a harmonic mode expansion of up to six terms to determine the time-energy-space dependence.

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ACKNOWLEDGEMENTS

I wish to thank Professor Theodore J. Williamson for his guidance, support and assistance in the completion of this work.

I. INTRODUCTION

Existing and future experimental work in the study of pulsed fast-neutron assemblies require the ability to determine the time dependent solution of the neutron transport equation accurately in the time period extending from the nanosecond to the millisecond range. This work allows the experimenter to compare both measured response and the analytical predictions for the overall improvement of each. Safe operation and control of fast nuclear reactors demands a fast, accurate mathematical simulation.

Current research on the response of fast assemblies to pulsed neutron sources considers that the time dependence can be represented by an exponential decay of the form:

$$T(t) = \sum_i A_i \exp(-\lambda_i t)$$

where the constants A_i and λ_i can be related to the composition and geometry of the system. A common assumption is that all terms above the first or fundamental will decay very quickly and that all measurable quantities can be described by a single fundamental or pseudo-fundamental value for λ_i . Experimental effort is being directed to answering the question of whether or not such a pseudomode does in fact exist, and if not, to understand why not.

This work will investigate several of the aspects of this problem.

In their movement through a material, neutrons undergo interactions, scattering or absorption, with the atoms and suffer changes in direction and velocity as the result of both elastic and inelastic collisions. While this is a random process for an individual neutron, the overall process for considering the time history and motion of the bulk of the neutrons is amenable to a statistical analysis. The Boltzman time-dependent transport equation:

$$\frac{\partial N}{\partial t}(\underline{r}, v, \underline{\Omega}, t) = -v\underline{\Omega} \cdot \nabla N(\underline{r}, v, \underline{\Omega}, t) - v\Sigma_t(v)N(\underline{r}, v, \underline{\Omega}, t) + \int_{\underline{\Omega}} \int_v v' \Sigma(v') N(\underline{r}, v', \underline{\Omega}', t) \cdot f(v', \underline{\Omega}' \rightarrow v, \Omega) d\underline{\Omega}' dv' + S(\underline{r}, v, \underline{\Omega}, t)$$

can describe the net processes occurring in any elemental volume of a physical system to account for all changes in the density of neutrons, subject to the usual conventions:

- \underline{r} Position of the volume element dV .
- v Speed.
- $\underline{\Omega}$ Direction. (The vector velocity would be defined as $\underline{v} = v\underline{\Omega}$).
- $N(\underline{r}, v', \underline{\Omega}', t)$ Density of neutrons at time t having velocity $v' \underline{\Omega}'$.
- $\Sigma(v')$ Speed dependent cross section for neutrons with speed (v')
- $f(v', \underline{\Omega}' \rightarrow v, \Omega)$ Probability that a neutron having velocity $(v' \underline{\Omega}')$ will interact to become one velocity $(v \underline{\Omega})$.

The neutron density variation with time considers:

(a) leakage of neutrons:

$$-v\underline{\Omega} \cdot \nabla N(\underline{r}, v, \underline{\Omega}, t),$$

(b) losses due to interactions (scattering or absorption) in the volume element:

$$-v\Sigma_t(v) N(\underline{r}, v, \underline{\Omega}, t),$$

(c) scattering of neutrons into the element phase space (position-velocity-direction) due to interaction of neutrons having other initial parameters:

$$\iint v' \Sigma_t(v') N(\underline{r}, v', \underline{\Omega}', t) C_t' f(v', \underline{\Omega}' - v, \underline{\Omega}) dV' d\underline{\Omega}'$$

neutrons introduced into the element of phase space from a source in that element:

$$S(\underline{r}, v, \underline{\Omega}, t).$$

Complete derivations of the neutron transport equation can be found in Tait [1] and Davison [2] and is discussed in detail in many other texts on nuclear reactor physics.

During the past twenty years, numerous methods and approximations have been formulated to solve this integro-differential equation, either in a closed analytical form, or by a numerical approximation, for both the time dependent and time-independent cases. While a major portion of this effort has been in obtaining the solutions applicable to thermal reactors, uranium-235 systems, the need to conserve

fuels into the distant future requires fast breeder reactors. Adequate understanding of the complete time, space and energy dependence of fast neutron populations requires the development of new analytic and simulation models for the fast neutron transport and fission problem.

Analytical solutions are often based on extension of the existing models for thermal systems. These solutions usually assume $1/v$ or constant neutron cross sections, with the age-diffusion theory [3] considering the steady state, time independent cases and with perturbation analysis for time dependent solutions [4]. These provide representations of the steady state systems, and simple kinetic models for fast reactors.

(1) Numerical methods generally deal with the time independent solutions of the transport equation by two principal methods:

Multi-group theory accounts for all energy dependent parameters by dividing the range of neutron energies into some finite number of intervals and defining average values of the energy dependent terms for each energy group. This method is frequently used to study the energy and time relations and the energy and position dependence [5]. This method has been used to consider up to several hundred groups between fission energies and the thermal condition [6].

(2) Monte Carlo is a statistical method to consider the interactions of individual neutrons to determine principal integral or system values for several spatial regions using multiple energy groups. The study of the individual

neutron "life history" from its birth to its eventual loss, allows determination of such parameters as the effective multiplication factor (K_{eff}). Computational requirements for this method limit its usefulness mainly to time independent calculations, and to a limited extent [7], the time dependent problem using very few velocity groups.

The description of the slowing down and thermalization of neutrons as a probabilistic Markov chain process was proposed by Perkel [8] in 1960, and expanded to consider the thermal energy region (0.00002 to 1.034 eV) by Ohanian and Daitch [9] for the study of the time dependent neutron spectra. Jenkins and Daitch [10] have extended this technique to formulations of the study of a pseudo-fundamental mode decay for time-energy response of pulsed fast systems, assuming that the system response can be represented by the lowest Fourier spatial mode of the time dependent diffusion equation for two energy regions: .001 to 1.0 eV and 1.0 to 1000 Kev.

Williamson and Albrecht [11] have extended the Markov chain method of treatment for the slowing down problem to cover the continuous range from 10.5 MeV to thermal energies.

Menzel, *et al*, [12] have employed the method of Ohanian and Daitch to provide a complete harmonic expansion method for the study of the space-time-energy response of pulsed systems in the interval 0.01 to 1.0 eV.

The purpose of this thesis is to apply the numerical methods of the Markov chain process to describe the energy

and time response of pulsed fast systems, with the harmonic expansion techniques for the spatial dependence similar to Menzel, in the interval 10.5 Mev to thermal energies.

Current work on the pulsed neutron response of fast assemblies, considers that the time decay can be represented by some exponential function of the form:

$$T(t) = \sum_i A_i \exp(-\lambda_i t)$$

where the constants λ_i and A_i can be related to the composition and geometry of the assembly.

In this study, the fundamental and higher harmonic space modes are combined with energy-time expressions to determine the effectiveness of this technique is simulating the experiments conducted in the laboratory. Experiments suggest that the decay of the neutron population is often not best described by the pseudo-fundamental mode. Hopefully, this work will aid in determining if this is due to the effects of a constantly varying velocity spectrum, or to the persistence of higher spatial modes of the time response remaining for longer periods of time.

II. THE SLOWING DOWN EQUATION

A. THE BOLTZMAN TRANSPORT EQUATION

Starting with the general form of the time dependent transport equation, reasonable assumptions which will permit this integro-differential equation to be restated as the slowing down equation are:

(1) The neutron density $N(\underline{r}, v, \underline{\Omega}, t)$ is isotropically distributed in $\underline{\Omega}$ at all points in the physical media.

(2) All sources and scattering kernels are isotropic.

(3) No delayed neutron sources.

(4) The physical media is composed of locally homogeneous isotropic materials.

(5) Spatial dependence is included within the bounds of the diffusion approximation.

(6) Boundry conditions are that the neutron density is zero at the extrapolated boundry and that the extrapolation distance is constant for all energies.

Conditions 1, 2 and 4 permit the equivalence of the expression for the local streaming or leakage of neutrons from a volume element to be described as:

$$-v\underline{\Omega} \cdot \nabla N(\underline{r}, v, \underline{\Omega}, t) = D(v) \nabla^2 N(\underline{r}, v, \underline{\Omega}, t) = -B^2 D(v) N(\underline{r}, v, \underline{\Omega}, t)$$

where B^2 is called the geometric buckling.

Within this analysis, a source that is a delta function in time will be considered as the mathematical equation of

the pulsed neutron source with a very small time duration, which can be written as:

$$S(\underline{r}, v, \underline{\Omega}, t) = S'(\underline{r}, v, \delta(t))/4\pi.$$

Adjusting all constant terms to account for the isotropic conditions of the geometry under study, and to match the primary condition that the neutron density at a time ($t=0$) will be equivalent to the value of the source strength at that time we have:

$$N(\underline{r}, v, 0) = S(\underline{r}, v, 0) = \sum_n R_n(\underline{r}) F_n(v, 0),$$

where solutions of the form $N(\underline{r}, v, t)$ serve as Green's functions for the pulsed conditions with an arbitrary time and spatial parameter.

Starting with the basic assumption that the neutron density at any time can be described as an infinite sum of solutions of the form:

$$N(\underline{r}, v, t) = \sum_n R_n(\underline{r}) F_n(v, t),$$

it is possible to transform each of the terms of the general transport equation to the slowing down equation in the following manner:

$$(1) \quad \frac{\partial N(\underline{r}, v, \Omega, t)}{\partial t} = \sum_n R_n(\underline{r}) \frac{\partial F_n(v, t)}{\partial t}.$$

$$\begin{aligned}
 (2) \quad -\underline{v}\underline{\Omega} \cdot \nabla N(\underline{r}, \underline{v}, \underline{\Omega}, t) &= -D(\underline{v}) \nabla^2 N(\underline{r}, \underline{v}, t) \\
 &= -D(\underline{v}) \sum_n \nabla^2 R_n(\underline{r}) F_n(\underline{v}, t).
 \end{aligned}$$

$$\begin{aligned}
 (3) \quad \int \underline{v}' \Sigma_t(\underline{v}') \int N(\underline{r}, \underline{v}', \underline{\Omega}', t) C_t, f(\underline{v}', \underline{\Omega}' \rightarrow \underline{v}, \underline{\Omega}) d\underline{v}' d\underline{\Omega}' &= \\
 \int \underline{v}' \Sigma_t(\underline{v}' \rightarrow \underline{v}) \sum_n R_n(\underline{r}) F(\underline{v}', t) d\underline{v}' &
 \end{aligned}$$

which reduces to the now modified form of the slowing down equation:

$$\begin{aligned}
 \sum_n R_n(\underline{r}) \left\{ \frac{\partial}{\partial t} F_n(\underline{v}, t) + \Sigma_t(\underline{v}) F_n(\underline{v}, t) - \int \underline{v}' \Sigma_t(\underline{v}' \rightarrow \underline{v}) \cdot \right. \\
 \left. F_n(\underline{v}', t) d\underline{v}' \right\} + \sum_n D(\underline{v}) F_n(\underline{v}, t) \nabla^2 R_n(\underline{r}) = \\
 S'(\underline{r}, \underline{v}, t)
 \end{aligned}$$

substituting for the spatial-velocity dependent delta function source, the eigenfunctional expression of the slowing down equation is obtained:

$$\begin{aligned}
 \sum_n \left\{ R_n(\underline{r}) \left\{ \frac{\partial}{\partial t} F_n(\underline{v}, t) + \Sigma_t(\underline{v}) F_n(\underline{v}, t) + \int \underline{v}' \Sigma_t(\underline{v}' \rightarrow \underline{v}) \cdot \right. \right. \\
 \left. \left. F_n(\underline{v}', t) d\underline{v}' \right\} + D(\underline{v}) F_n(\underline{v}, t) \nabla^2 R_n(\underline{r}) \right\} = \sum_n R_n(\underline{r}) F_n(\underline{v}, 0).
 \end{aligned}$$

Consolidating all expansion functions of the index-n in the single summation,

$$\sum_n \left[R_n(\underline{r}) \left\{ \frac{\partial F_n(v,t)}{\partial t} + \Sigma_t(v) F_n(v,t) + \int v' \Sigma_t(v'-v) \cdot F_n(v',t) dv' - F_n(v,0) \right\} + D(v) F_n(v,t) \nabla^2 R_n(\underline{r}) \right] = 0$$

it is now possible to separate the spatial dependence from the energy time relations via the harmonic buckling factor, B_n^2 , to two expressions:

$$\frac{\nabla^2 R_n(\underline{r})}{R_n(\underline{r})} = + B_n^2$$

and:

$$-D(v) B_n^2 F_n(v,t) = \frac{\partial}{\partial t} F_n(v,t) + \Sigma_t(v) F_n(v,t) + \int v' \Sigma_s(v'-v) F_n(v',t) dv' - F_n(v,0).$$

Final clearing and rearrangement of all terms yield the decoupled expressions:

$$\nabla^2 R_n(\underline{r}) - B_n^2 R_n(\underline{r}) = 0$$

$$\frac{\partial}{\partial t} F_n(v,t) + B_n^2 D(v) F_n(v,t) + \Sigma_t(v) F_n(v,t)$$

$$+ \int v' \Sigma_s(v' \rightarrow v) F_n(v',t) dv' = 0$$

with $R_n(\underline{r})$ a time independent function and $F_n(v,t)$ a space independent function.

B. MOD-5: THE DISCRETE STATE APPROACH

In the separated form of the slowing down equation, the velocity-time harmonic term must satisfy the same form of the integro-differential equation as the space independent neutron density function.

$$\frac{\partial}{\partial t} N(v,t) = -vD(v) B^2 N(v,t) - v\Sigma_t(v)N(v,t) + \int v' \Sigma_s(v' \rightarrow v) N(v',T) dv' + S(v,t).$$

Any method that provides a solution to the space independent slowing down equation can be applied to solve the velocity-time harmonic equation.

Williamson and Albrecht [11,13] have developed a stochastic model for neutron moderation that provides a numerical solution of the slowing down of fast neutrons in a finite media. Since the slowing down process is a continuous time process in which collisions may occur at any time, having a continuous range of energy transition possibilities, it may be classified as a continuous time, continuous state Markov process. If the neutron cross sections remain constant with time, the slowing down may be classed as a stationary Markov system.

Application of the formalism of the Markov process calculations provides the capability to perform straightforward computer solutions for a discrete time-velocity model of the space independent equation. The probability that a neutron will be in a finite width velocity state, v_i , at a time, t_j ,

is calculated by a computer program, MOD-5, which follows the slowing down of neutrons by determining a probability density parameter $F(v_i, t_j)$ which is the probability that a neutron will be in the velocity state v_i at the discrete time $t_j = n\Delta t$.

The Markov process describes the variation of a probability density vector, $s(n\Delta t)$, which describes the probability that a neutron will be located at some state in the system during its entire lifetime from birth to death. This state vector is given as:

$$\bar{s}(n\Delta t) = s_1, s_2, s_3, s_4, \dots, s_N$$

where s_i is the probability that the neutron is in the state bounded by velocities v_i and v_{i-1} (where $v_i < v_{i-1}$) at the time $n\Delta t$. This method presumes that the initial state vector, neutron source, $\bar{s}(0)$, can be described for time $t=0$.

The problem description necessary for application of a MOD-5 type of approximation are:

(a) The velocity range of interest is divided into "M" discrete states.

(b) Collision physics is applied to construct an array, $P(i,j)$, of the discrete Markov transition probabilities of a neutron going from velocity state- i to state- j in the discrete time period Δt .

While the details of the Markov transition matrix are discussed in several sources [14], the main considerations are that the array for the slowing down approximation is

upper triangular with only positive elements, with the elements of a given row summing to unity, which states that all possible outcomes of the process are considered.

If the initial Markov transition matrix describes the probability for a transition within a specified time interval, Δt , the probabilities for a transition at a time $\{n\Delta t\}$ can be determined as:

$$\bar{P}_{n\Delta t} = (\bar{P}_{\Delta t})^n.$$

The probability of a neutron being in any state at a time $(-n\Delta t)$ can be determined from the $(n-1)$ th step in the matrix product as:

$$\bar{s}(n\Delta t) = \bar{s}((n-1)\Delta t) \cdot \bar{P}_{\Delta t}.$$

Both the program MOD-5 and the harmonic expansion method employ several series of notations and can result in some conflict with existing literature. In this work, the following definitions and terms will be used:

- $\bar{s}(n\Delta t)$ the probability density vector of a neutron being in any state in the system at the time $n\Delta t$.
- $S_{i,j}$ the probability that a neutron will be in a state "i" at a time t_j ,

where the time and total density vector can be written as:

$$t_j = n\Delta t$$

$$\bar{s}(t_j) = (S_{1,j}, S_{2,j}, \dots, S_{I,j}).$$

While this form of notation applies to the case of a single decay mode, the expansion to a multiple harmonic expansion method requires the addition of another indexing parameter- n to indicate the particular harmonic mode. This results in the probability density expressions being written as:

- $S_{n,i,j}$ - probability of the neutron fraction of the n -th harmonic mode being in a velocity state v_i at the time t_j .
- $F_n(v_i, t_j)$ - n -th harmonic term of the time velocity dependence of the neutron probability density.

The contribution of each harmonic term of the neutron density is the product of the spatial component and the probability density component:

$$N_n(x, v_i, t_j) = R_n(x) F_n(v_i, t_j)$$

where the space component, $R_n(x)$, is a measure of the fractional contribution to the neutron population of the n -th harmonic mode, and the term, $F_n(v_i, t_j)$ is the probability of a neutron being in the velocity state v_i at time t_j for the n -th mode of the time decay harmonic.

In the following discussions, the equivalence of the terms is described as:

$$S_{n,i,j} = F_n(v_i, t_j) \Delta v_i$$

where Δv_i is the width of the velocity state about the value v_i .

III. GEOMETRICAL MODELS

A. SOURCE FUNCTIONS

The main objective of this thesis is the investigation of the harmonic mode expansion method for the space, time and energy dependence of the neutron density in a physical system. Simple source functions representative of actual systems are desired to demonstrate this method. Three categories of source type were considered: a finite volume source and a first collision source as interior sources in the system, and a wide beam exterior source condition.

The test geometries are:

(a) Finite Volume Source - a finite volume element, located on an axis of symmetry and at an off axis position is considered to represent a hole or cavity in the slab. This might be envisioned as a beam port or localized neutron generator capable of providing a pulse of monoenergetic neutrons in a small volume in the slab. This is represented as a fourier square pulse on the x-axis of the system.

(b) First Collision Source - this exponential distribution of neutrons in the system is considered as a more realistic approach for the response of a physical system than the fixed source in volume condition. Rather than assume that the initial burst of neutrons would remain in the fixed volume element, it can be argued that they would initially travel throughout the system with an exponential form.

As with the finite volume type source, this source type is studied for both the symmetric and unsymmetric positions.

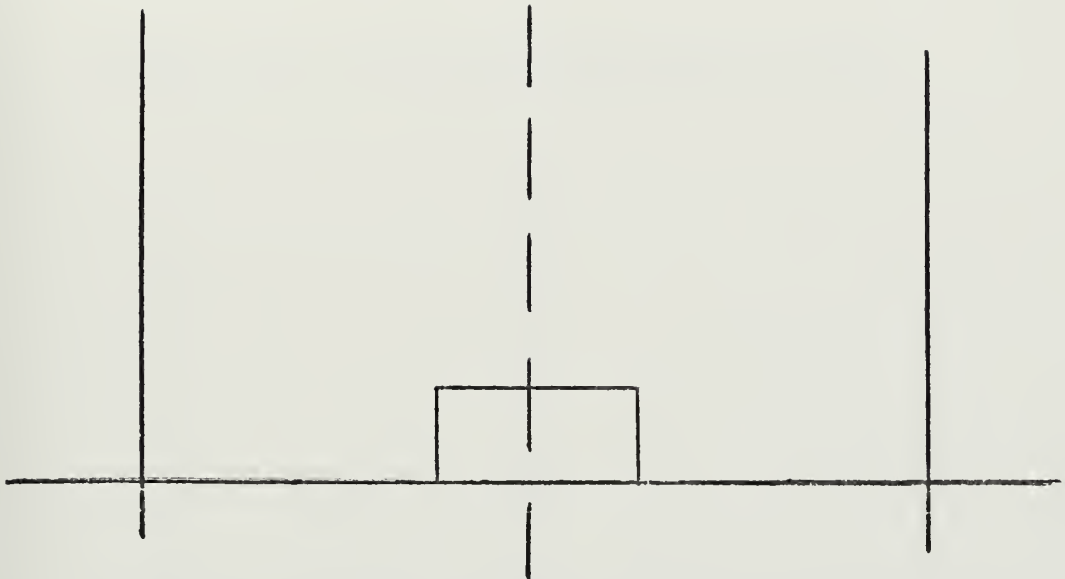
(c) Exterior Source Beam - a higher level of complexity in the source geometry selection process is the physical situation where the assembly under study is subject to a wide beam of fast neutrons from an external point. This can be considered as a very simple approximation to a first collision response from a plane source of neutrons on the surface of the slab. While this may be considered to follow an exponential type distribution through the slab, it is approximated as a straight line or ramp function within the confines of the slab.

B. CHOICE OF DIMENSIONALITY

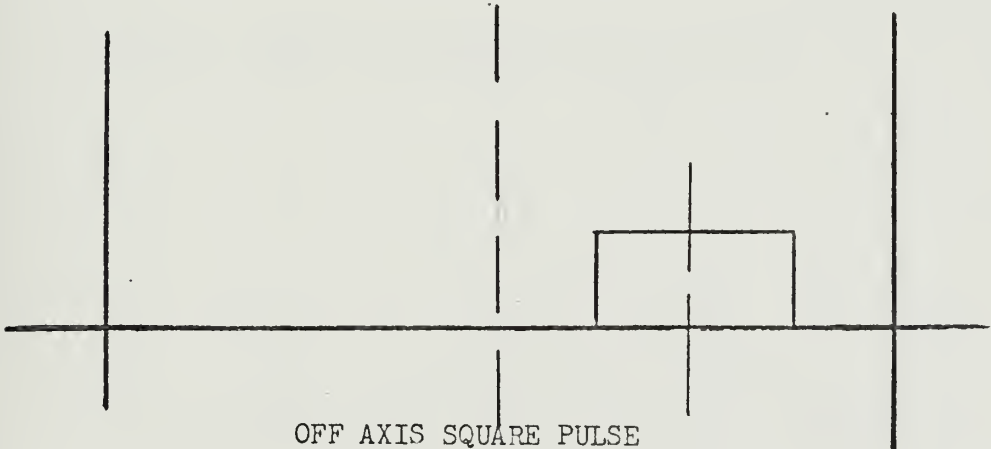
The principal goal of this work is the development of a useful mathematical model to determine the spatial dependent response of a physical system to a burst of fast neutrons. In this determination, it will be necessary to refer to the neutron density in terms of three main parameters - time, space and energy. The derivation of all equations and principal relationships will be written for the time, position and velocity of the neutrons. Equivalent expressions can easily be obtained using either the neutron energy (E) or lethargy (u), using the relationships:

$$E = \frac{1}{2} m_n v^2$$

$$u = -\ln(E_0/E) = -2 \ln(v_0/v).$$



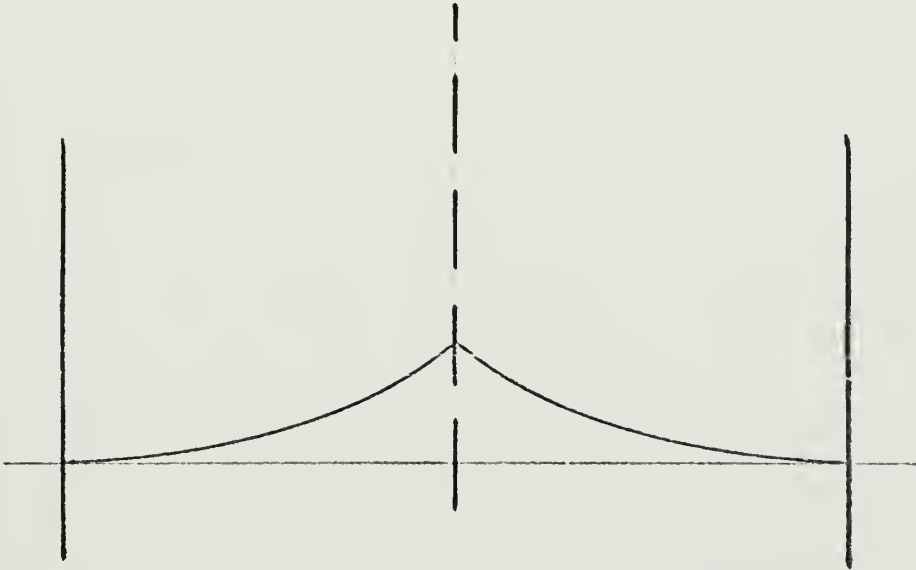
CENTRAL AXIS SQUARE PULSE



OFF AXIS SQUARE PULSE

FIGURE 1 SOURCE GEOMETRY SQUARE PULSE

CENTRAL AXIS FIRST COLLISION-EXPONENTIAL SOURCE



OFF AXIS FIRST COLLISION-EXPONENTIAL SOURCE



FIGURE 2 SOURCE GEOMETRY FIRST COLLISION SOURCE



FIGURE 3 SOURCE GEOMETRY EXTERIOR SOURCE RAMP FUNCTION

Employing the appropriate expressions, it can be easily shown that equivalent terms for the neutron density, $N(\underline{r},v,t)$, the time-space-velocity form can also be modified to give the time-space-energy or time-space-lethargy form. While the velocity dependent form is used for the purposes of demonstration within this paper, all calculations performed by the programs MOD-5 and MIL-6 employ the lethargy dependent forms.

In the determination of numerical solutions for the slowing down equation, reasonable methods and simple systems are desired to demonstrate the harmonic expansion method for the time, velocity and spatial expressions, where the method is general enough to be easily revised to handle more difficult and complex problems. As shown previously, the neutron density can be determined by the summation of a harmonic series whose terms are the product of a space function and a time-velocity function.

$$N(\underline{r},v,t) = \sum_n R_n(\underline{r}) F_n(v,t)$$

where the spatial dependent function, $R(\underline{r})$, satisfies the differential equation:

$$\nabla^2 R_n(\underline{r}) - B_n^2 R(\underline{r}) = 0.$$

In order to retain the flexibility to study and compare the response of physical systems to various source types, the selection of the dimensionality for problem selection posed some difficult choices.

The spatial functions must form an orthogonal set, a condition easily met by the Fourier expansion methods for the one dimensional problem using trigonometric functions for the cartesian geometry, but requires the use of regular Bessel functions for cylindrical geometry, and the spherical Bessel and spherical harmonic functions for the spherical geometry. A practical consideration was to obtain a balance demonstrating the harmonic expansion application with reasonable computing requirements. The more complex the spatial functions, the longer computing time and large core storage space required. The trigonometric functions can be determined quickly, while the existing methods available for the Bessel function calculations require approximately three times as much computing time.

In all three of the principal geometrical coordinate systems (cartesian, cylindrical and spherical), studies of sources located at an axis of symmetry resolve themselves into some form of the one dimensional problem. However, the study of sources located off the axis of symmetry of the system, while seeming reasonable conditions to analyze, can lead to some very messy mathematics, such that the utility of the method can become lost in the details. While these problems can be considered, the one dimensional case would allow easy comparison of all symmetrical and non-symmetrical cases without initially confusing the reader or user of this technique prior to demonstrating the method. After accomplishing that aim, the multi-dimensional cases (two or three dimensions) can be investigated.

A practical matter in the consideration of the harmonic expansion method involves the retention of data within the computing machinery while running the basic problem or one of the many variations that can be considered by this method.

An important consideration while developing the computer programming for these calculations involved the selection of methods to transfer data between the two computational routines. The program MIL-6 was written to process data input from punch cards, nine-track magnetic tape, or the IBM-2314 series of magnetic disk. The program MOD-5 was run sequentially to determine the velocity and time characteristics of the physical system for each harmonic mode of the buckling constant (B_n^2). This generally requires one minute of computing time for the IBM-360/67 with 70 velocity states, requiring 200k bytes of core storage for each harmonic mode. Thus, a six-mode expansion would require approximately six minutes of central processor time.

Assuming that the problem under study is going to deal with 150 discrete velocity states, 20 time step values and a minimum of six harmonic modes of the functional data, $F_n(v,t)$, as produced by program - MOD-5, one is faced with the retention of 18000 numbers which are required either in core storage or in an on-line direct access device. In the IBM-360/67, all real numbers require a 4 byte location, and a minimum problem would require 72k bytes (12 per mode)

available during the running of the problem. Once the primary calculation of the density vector $N(\underline{r},v,t)$ is done, data retention is required for the total system $N(\underline{r},v,t)$ which would be estimated as 4.0 (position points) (dimensions) (energy states) (time intervals) which can easily become a very large number when considering 40 or 50 points in a physical system, for cases in two or three dimensions.

The crux of the matter is that the one dimensional model will allow the utility of this approach to be demonstrated easily with relatively simple, easily compared Fourier series approximations in sine and cosine terms, rather than the functional forms required for solution in three dimensional cylindrical or spherical coordinate systems for sources not located at the principal axis symmetry.

C. SOURCE FUNCTION APPROXIMATIONS

Within the bounds of the one dimensional spatial model, the five source functions can be described by comparatively simple Fourier series approximations:

$$S(x, v_i, 0) = \sum_n (A_n \cos(k_n x) + C_n \sin(k_n x))$$

for a representation of the source geometry within the confines of the slab.

The following labeling conventions are used in defining the terms of the Fourier series representation:

S_0 source strength-total neutrons in the pulse.

x_0 half width of the slab.

- x_1 half width of the finite volume source.
 x_2 mid-point of the source for the off axis conditions
 L diffusion length for the source velocity neutrons.
 λ the extrapolation distance for the source neutrons.
 The source function approximations are subject to the conditions:

Square Pulse

$$\begin{aligned}
 \text{On-Axis} \quad S(x, v_j, 0) &= \begin{cases} 0 & x < -x_1 \\ S_0 & -x_1 \leq x \leq x_1 \\ 0 & x > x_1 \end{cases} \\
 \text{Off-Axis} \quad S(x, v_j, 0) &= \begin{cases} 0 & x < x_2 - x_1 \\ S_0 & x_2 + x_1 \leq x \leq x_2 - x_1 \\ 0 & x > x_2 + x_1 \end{cases}
 \end{aligned}$$

First Collision Source

$$\begin{aligned}
 \text{On-Axis} \quad S(x, v_j, 0) &= \begin{cases} S_0 e^{+ax} & x \leq 0 \\ S_0 e^{-ax} & x \geq 0 \\ 0 & x > x_1 \end{cases} \\
 \text{Off-Axis} \quad S(x, v_j, 0) &= \begin{cases} S_0 e^{-a|x-x_2|} & -x_0 \leq x \leq x_0 \\ 0 & x < -x_1 \end{cases}
 \end{aligned}$$

Exterior Source

$$S(x, v_j, 0) = \begin{cases} mx + b & -x_0 \leq x \leq +x_0 \\ m = \frac{S_0}{2x_0^2} & b = \frac{S_0}{2x_0} \end{cases}$$

Source Type	Series Representation	Fourier Coefficients
(1) Square Pulse on axis n = positive odd integers	$A_n \text{Cos} k_n x$	$k_n = \frac{n\pi}{2(x_0 + \lambda)}$ $A_n = \frac{S_0 \sin(k_n x_1)}{k_n x_0}$
(2) Square Pulse off axis n = positive odd integers	$A_n \text{Cos} k_n (x - x_2)$	$k_n = \frac{n\pi}{2(x_0 + x_1 + \lambda)}$ $A_n = \frac{S_0 \sin k_n x_1}{2n x_1}$
(3) First Collision Source on axis n = positive odd integers	$A_n \text{Cos} k_n x$	$k_n = \frac{n\pi}{2(x_0 + \lambda)}$ $A_n = \frac{S_0 (-1)^{\frac{n-1}{2}}}{n \left(\frac{1}{4} + (x_0/nL)^2 \right) (1 - \exp(-x_0/L))}$
(4) First Collision off axis source n = positive integers	$A_n \text{Cos} k_n x$ $+ C_n \text{Sin} k_n x$	$k_n = \frac{n\pi}{2(x_0 + x_2 + \lambda)}$ $A_n = \frac{S_0 (\text{Cos} k_n x_1 - k_n L \text{Sin} k_n x_1 - (-1)^{n+1} \sin(x_1/L))}{L(1 + (k_n L)^2) (1 - \exp(-(x_0 + x_1)/L))}$

TABLE I. Fourier Expansion Coefficients.

TABLE I. - Continued

$$C_n = \frac{S_0 (\text{Sink}_n x_1 - k_n L \text{Cos} k_n x_1 - (-1)^n \sinh(x_1/L))}{L(1+(k_n L)^2)(1-\exp(-(x_0+x_1)/L))}$$

(5) Exterior Ramp $A_n \text{Sink}_n x$

Function Source

n = positive integers

$$k_n = \frac{n \pi}{4(x_0 + \lambda)}$$

$$A_n = \frac{(-1)^{n+1} 2 S_0}{n \pi x_0}$$

D. ENERGY-VELOCITY-LETHARGY

The slowing down equation can be written in several forms to describe the variation of the neutron energy with time. For computational purposes, it is often convenient to describe the kinetic parameter in terms of the neutron velocity, while at other times, the kinetic energy of the neutron would provide a better description. There are three main terms that can be used interchangeably to describe the neutron during the slowing down process: velocity, energy and lethargy. Any expression in one of these parameters can be transformed to the appropriate form in the others by the relationships:

$$\text{Energy-velocity} \quad E = m_n v^2/2$$

$$\text{Energy-lethargy} \quad - \ln(E/E_0) = u$$

$$\text{Velocity-lethargy} \quad - 2\ln(v/v_0) = u.$$

By these expressions, the parameters, E_0 and v_0 , describe the neutron reference or source condition with the added convenience that as the neutron slows down, its lethargy increases, while for the velocity and energy expressions, the slowing down process obviously results in the neutron going to a state with a smaller value.

Simple relations between the derivatives of these expressions will allow the slowing down equation to be easily revised with the desired terms:

$$dE = m_n v \, dv$$

$$du = \frac{-2 dv}{v} = \frac{-dE}{E} .$$

In terms of a neutron density function, $N_1(v)$ and the differential relations, all three forms of the neutron density are:

$$N(v) dv = N(E) dE \quad N(v) = m_n v N(E) = (2m_n E)^{\frac{1}{2}} N(E)$$

$$N(E) dE = N(u) du \quad -E N(E) = N(u)$$

$$N(v) du = N(u) du \quad N(u) = \frac{-v}{2} N(v) .$$

IV. APPLICATIONS

A. NEUTRON DENSITY

The neutron density can be determined in the numerical model by the infinite series expression:

$$N(\underline{r}, v, t) = \sum_1^N R_n(\underline{r}) F_n(v, t)$$

such that the series representation converges to the analytical value as N increases to infinity. In the numerical method for this calculation, the series representation will be determined as a finite series, truncated to the first 3 to 6 terms. In this form of the approximation, the error introduced is principally that of under estimating the true value if an even number of terms are used, while slightly over estimating with an odd number of terms.

The truncation error of the finite series approximation can be estimated from the initial conditions of the system at time equal zero as the neutron density at that time must equal the initial source strength:

$$N(\underline{r}, v, 0) = S(\underline{r}, v, 0).$$

If the total source strength (S_0) is obtained by integrating over all velocities and the source volume, the source strength can be given by the series approximation:

$$S_0 = \int_v \int_r S(\underline{r}, v, 0) dv d\underline{r} = \int_v \int_r \sum_n R_n(\underline{r}) F_n(v, 0) dv d\underline{r}.$$

With the use of a truncated series approximation, the error introduced can be calculated as:

$$\text{Error} = S_0 - \iint \sum_1^N R_n(\underline{r}) F_n(v,0) d\underline{r} dv.$$

This normalization error is used to obtain a first order adjustment factor, Δ_0 , by the simple approximation:

$$\Delta_0 = \frac{\text{Error}}{S_0}$$

and a normalization adjustment is determined as:

$$\text{Source Normalization} = (1 + \Delta_0) S_0$$

such that the truncated series will give a numerical value equal to the initially defined parameter- S_0 . This now gives the series approximation as:

$$S(\underline{r},v,0) = \sum_1^N (1 + \Delta_0) R_n(\underline{r}) F_n(v,0)$$

where the artifice of the source strength normalization factor (Δ_0) would be a succeedingly smaller correction as more terms are included in the finite series representation.

This adjustment of the numerical value of the source strength has no effect on the velocity-time functions, $F_n(v,t)$, as this is a constant multiplier included in the spatial dependent expressions.

In the discrete state, discrete time calculations of MOD-5, the probability that a neutron will be within a

specified velocity interval (v_i) at the finite time step (t_j), transforms the neutron density expression to the form:

$$N(\underline{r}, v_i, t_j) = \sum_n R_n(\underline{r}) F_n(v_i, t_j)$$

and the total neutron population at any fixed time can be found by integration over the spatial coordinate and summation over the velocity intervals:

$$N(t_j) = \int_r \sum_i \sum_n R_n(\underline{r}) F_n(v_i, t_j) d\underline{r}.$$

B. NEUTRON FLUX

The neutron flux, a scalar quantity to describe the net flow of neutrons per unit time and unit area, is calculated as the product of the speed of the neutron and the number density of neutrons having that speed:

$$\phi(\underline{r}, v, t) = v N(\underline{r}, v, t)$$

and is approximated in this work as the discrete value:

$$\phi(\underline{r}, v_i, t_j) = v_i N(\underline{r}, v_i, t_j).$$

The total integrated flux can be analytically determined by an integration over all speeds, and in this work, is therefore approximated by a summation over all discrete states of the density parameter $N(\underline{r}, v_i, t_j)$ to yield the result:

$$\phi(\underline{r}, t_j) = \sum_i v_i N(\underline{r}, v_i, t_j) = \sum_i \sum_n v_i R_n(\underline{r}) F_n(v_i, t_j).$$

C. MEAN ENERGY OF NEUTRONS

The spatial variation of the neutron mean energy provides valuable information on their migration and diffusion following the initial burst. This additional insight into the phenomena of "diffusion cooling" considers the general trend of the higher velocity components of the neutron population to leak out the boundaries of the finite system. As such, the remaining neutrons, while still undergoing interactions which remove them or lower their velocity, have their total number decreased by this additional effect.

In the discrete velocity state, discrete time model, the mean energy of the neutrons at a point \underline{r} in the slab is determined as:

$$\bar{E}(\underline{r}, t_j) = \frac{\sum_i E(v_i) N(\underline{r}, v_i, t_j)}{\sum_i N(\underline{r}, v_i, t_j)} .$$

In comparing the results of this mean energy determination, two basic methods can be used: first, to study the variation in time of the mean energy at a fixed position, and secondly, to consider the ratio of the mean energy at all positions in the system, for all time periods, compared to a single reference position. Within this work, the mean energy is compared at each finite time period to the current value of the mean energy at the center of the slab, while the time variation of the mean energy at the reference point is followed in detail.

In the discrete time-velocity-space model, this is determined as:

$$\frac{\bar{E}(\underline{r}, t_j)}{\bar{E}(\underline{r}_0, t_j)} = \frac{\frac{\sum_i E(v_i) N(\underline{r}', v_i, t_j)}{\sum_i N(\underline{r}', v_i, t_j)}}{\frac{\sum_i E(v_i) N(\underline{r}_0, v_i, t_j)}{\sum_i N(\underline{r}_0, v_i, t_j)}}$$

where \underline{r}_0 is the reference point, $x=0.0$, and \underline{r}' is a position within the slab.

D. DETECTOR RESPONSE

Numerous methods [15] are currently utilized to determine the neutron spectra and flux in fast reactor test assemblies, for internal and external measurements. Fast neutron leakage is examined in time of flight tests to evaluate the time-varying spectra and neutron resonance absorption foils to study the internal spatial dependent spectra.

For this analytical model, no single existing device appeared to be comparable to cover the continuous time and energy response as calculated by the discrete state method. An analytical detector is therefore constructed that would respond to the time-velocity variations of the neutron density. For future applications of this model, a one dimensional version of an existing devices characteristics can be introduced.

The analytical detector is modeled as a proton recoil scintillation device capable of measuring the proton recoil of hydrogen-neutron collisions. The detector is considered to be filled with hydrogen at a molecular concentration as H_2 equivalent to an ideal gas at STP., with physical dimensions of a width $2 X_4$ about a midpoint $\underline{r}_D = X_3$.

The detector response (DR) is determined using the ABN-26 group cross section [16] to be:

$$DR(\underline{r}_D, v_i, t_j) = \int \Sigma_s(v_i) v_i N(\underline{r}', v_i, t_j) d\underline{r}'$$

where the integral over the spatial coordinate $d\underline{r}'$ refers to the total detector volume.

The total response of the detector both for the total integrated flux over all velocities and over all times can be determined in this discrete velocity-time model by numerical integration and summation:

$$DRV(\underline{r}_D, t_j) = \sum_i DR(\underline{r}_D, v_i, t_j)$$

$$DRT(\underline{r}_D, t_j) = \sum_j DRV(\underline{r}_D, t_j)$$

where the final summation over all times from $t=0$ to $t=t_j$ provides the neutron fluence at the detector.

E. MIL-SIX COMPUTER PROGRAM

A general purpose computer program, MIL-SIX, was written in FORTRAN IV to process the time-energy (lethargy) response

date produced by MOD-5 to determine within this version of the transport approximation, the neutron density function- $N(\underline{r},u,t)$ for the one dimensional infinite slab geometry. Several minor modifications were made in MOD-5 to provide the ability to calculate the time-energy response of a system with multiple harmonic modes and to provide the neutron probability density for each lethargy state at the same time steps. The only form of output system tested involved the use of punch card output from MOD-5, however, provisions have been made to employ nine-track magnetic tape and/or on-line disk storage capabilities (IBM-2314 units). Program MIL-SIX provides the general user the ability to accept any of the three forms output from MOD-5, for up to six (6) harmonic modes of data. The program can be modified easily to adjust for a higher order approximation using more than six modes with a simple change in the size of the storage arrays in the program.

Program MIL-SIX is composed of the following subprograms and routines:

<u>Subroutine</u>	<u>Function</u>
Main	Principal control portion of the program; handle all logical decisions on which routines the program will perform
INCON1	Establishes all principal default parameters for the program; reads all problem definition statements (punch cards) which determine the parameters that the program

will need to calculate the desired functions; provides the user with instructions and information concerning actions the program will perform for the stated problem run.

- READ10 Processes all punch card output from program MOD-5.
- READ11 Processes all tape (9-track) from program MOD-5.
- ESTIMATE1 Provides guidance information of the execution time to be expected for the running of a particular problem.
- MODEL 1 Calculates the non-harmonic terms of the fourier space dependent functional expression for the source geometry specified for the particular problem; estimates the error introduced by the use of a finite number of terms in the fourier infinite series approximation; adjusts the values of the fourier non-harmonic coefficients to obtain a closer approximation to the infinite series representation; and provide a simple space dependent plot of the fourier expansion of the space dependent function on the on-line printer.
- PHI1 Principal function sub-program of the program: to calculate the fourier space dependent harmonic terms in the calculation

of the terms of $R_n(x)$ (n =the harmonic mode, x =the space point) for the one dimensional model.

FOUT1(x_i)

Principal determination subprogram to calculate the series of values for the function - $N(x,u,t) = \sum R_n(x_k) F_n(u_i, t_j)$ for summed response due to all harmonic mode to give the neutron density for the single point - x_i .

FOUT2

This subprogram will calculate the two principal space dependent relationships: $\phi(x_k, t_j)$ and $\bar{E}(x_k, t_j)$ the integrated (total) neutron flux and the mean energy of the neutrons at the point (x_k) for a profile of 100 points across the slab for each time steps (t_j), and will provide a graphical plot of the functional forms:

$$\phi(x_k, t_j) \text{ vs } x_k$$

$$\bar{E}(x_k, t_j) / \bar{E}(0, t_j) \text{ vs } x_k$$

via a on-line printing of a rough plot of the results or via the Cal-Comp Plotter for each set of discrete time step date (t_j); also provides tabulated data on neutron flux, mean energy and mean energy ratio.

FOUT3

Detector response: this routine considers the simple model of a hydrogen recoil detector located inside the slab. This is not a true device, but one which determines the function:

$$\phi(x_k, t_j) = \sum_u v(u) * \Sigma_t * \sum_n R_n(x_k) F_n(u_i, t_j)$$

using the integrated position dependent flux (calculated at 5 points within the assumed detector position); a simple average is performed to obtain the flux at the detector, where the energy dependent microscopic cross section ($\Sigma_t(u_i)$) for STP hydrogen is based on the ABN-26 group cross section tables. Starting with $\phi(x_k, t_j)$ - flux data, a numerical integration is performed to determine the integrated detector response. Output of the two functional forms of the detector response are given both by the on-line printer, and the cal-comp plotter.

FOUT4

Provides a simple spectral plot of the data input from program MOD-5, to show the distribution $F_n(u_i, t_j)$ for each harmonic mode versus the lethargy state, showing the movement of the neutron distribution to states of lower energy (higher

lethargy) with time for each mode of harmonic buckling.

FOUT5

Provides a more detailed spectral response plot of the MOD-5 input data, by consideration of the fourier harmonic coefficient $R_n(x_k)$ for the individual position point - x_k weighting the harmonic data - to give the response spectra: $N_n(x_k, u_i, t_j) = R_n(x_k) F_n(u_i, t_j)$ as a single plot for each harmonic mode of input data, and a total spectral response function $N(x_k, u_i, t_j) = \sum_n R_n(x_k) F_n(u_i, t_j)$, summed over all modes and plotted versus lethargy state - u_k on the Cal-Comp plotter.

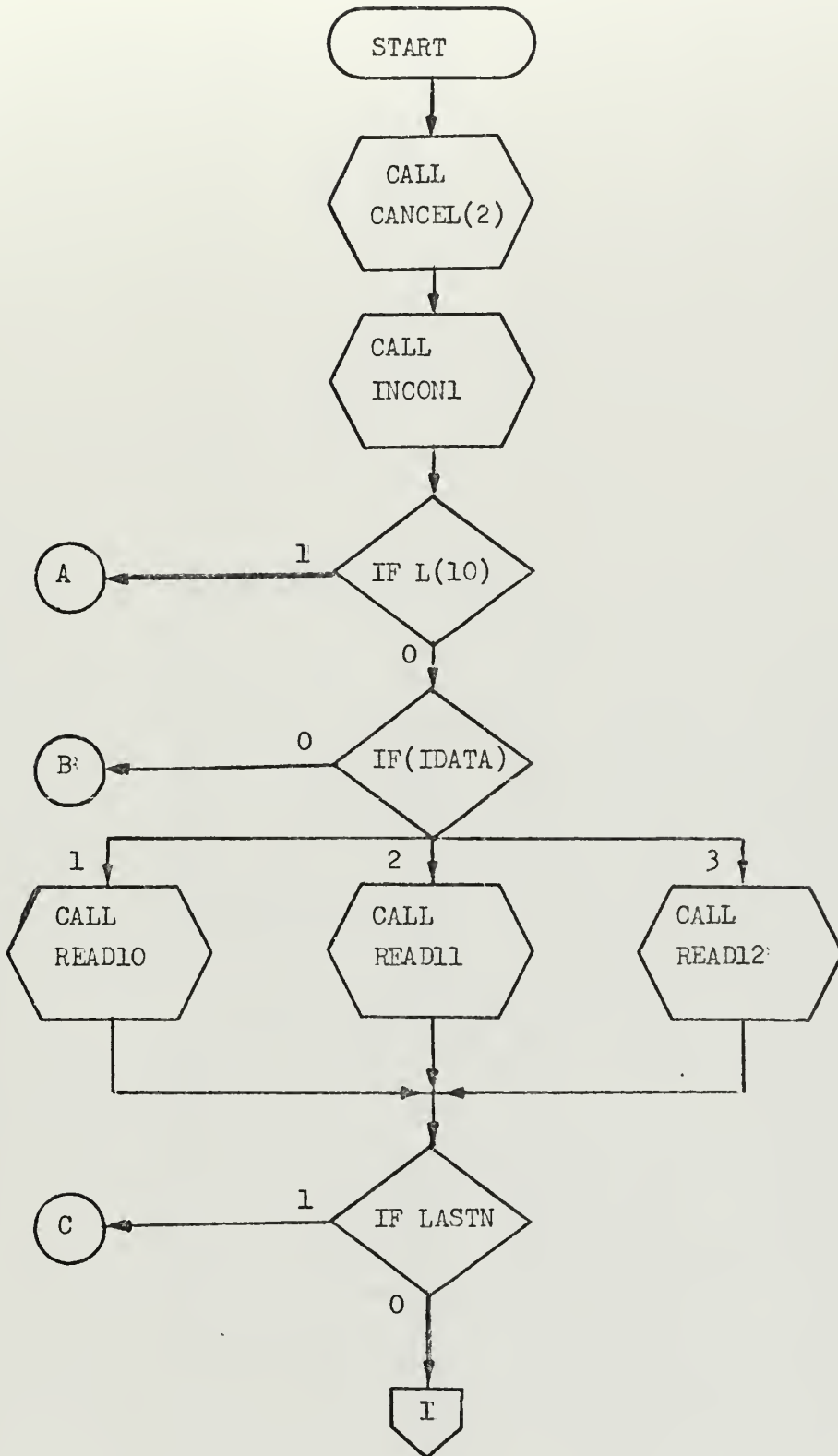
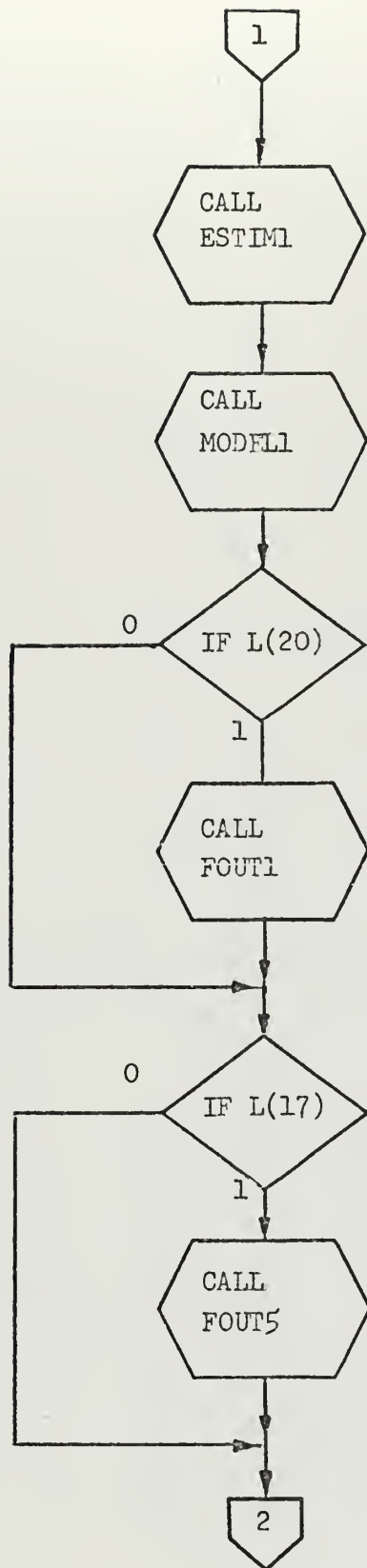
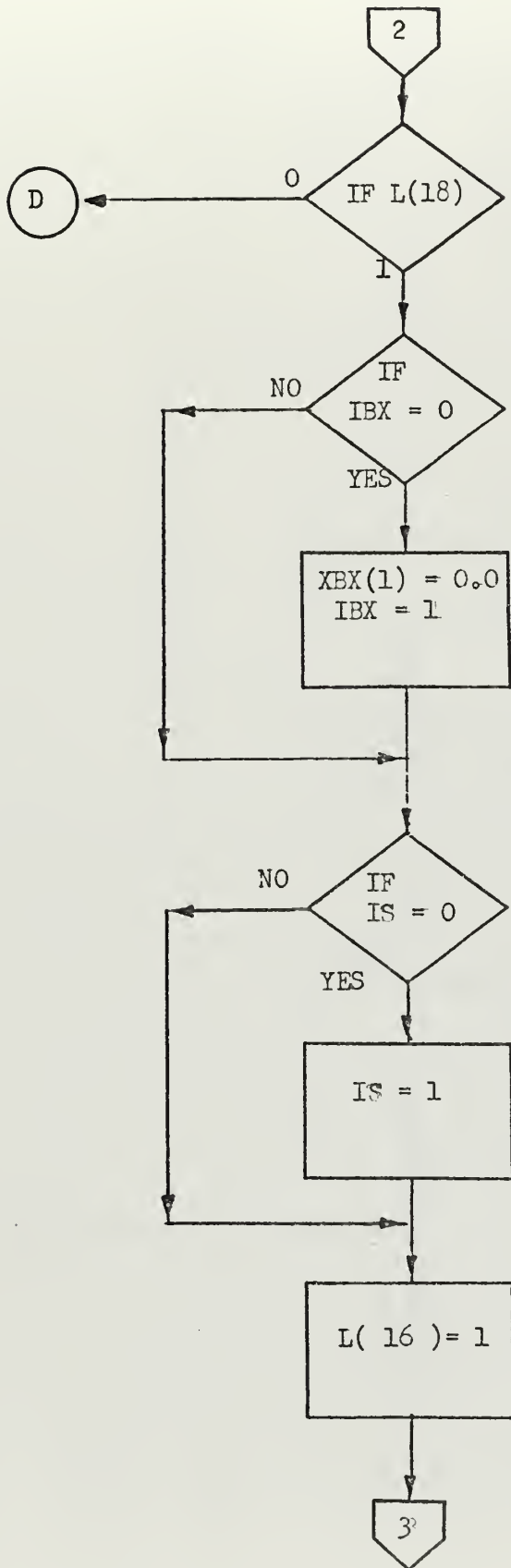
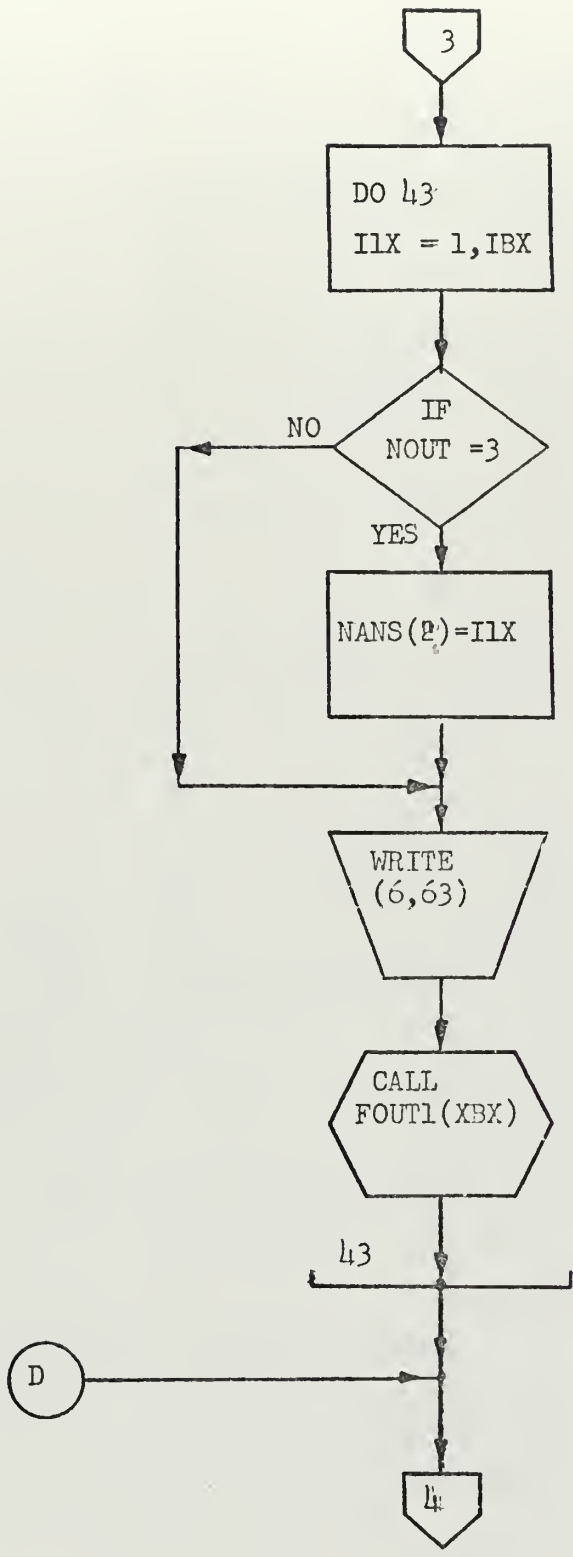
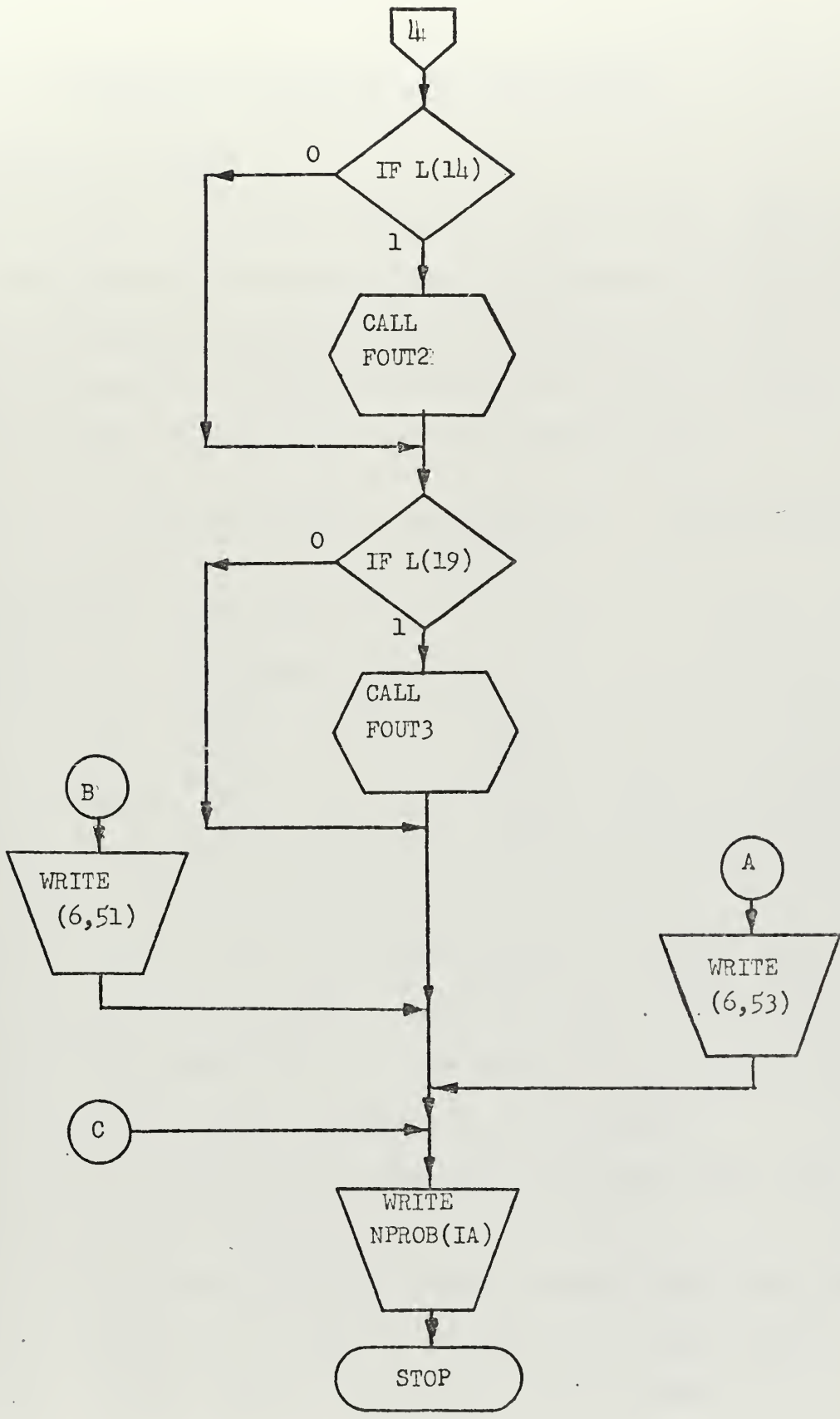


FIGURE 4. FLOW CHART COMPUTER PROGRAM MIL-SIX









V. DISCUSSION OF RESULTS AND CONCLUSIONS

A. PHYSICAL SYSTEMS

The harmonic expansion method was examined for application to multiplying and non-multiplying systems, a beryllium slab and a modified form of the ZPR-3 assembly 6F. The physical parameters of the two systems are given in Table II. The nuclear characteristics were assumed to be:

System	Constituents	Mass Density (gm/cc)	Total Modes
Beryllium Slab	Beryllium	1.84	6
ZPR-3 (6F)	Aluminum	0.848	3
	Iron*	0.965	
	U-235	2.622	
	U-238	3.016	

(*Iron substituted for stainless steel.)

A pulse of one 2.46 MeV neutron was fixed as a delta function in time to study the decay of each harmonic component of the discrete state velocity-time response $F_n(v_i, t_j)$. Two versions of the slowing down process were examined:

- (1) follow each harmonic mode to the same final time (t_j) as the fundamental mode,
- (2) follow decay of each harmonic mode to the time step at which the probability density component for that mode was below 0.001 of the neutron remaining in the system.

Both the beryllium and the assembly 6F systems were studied to determine the variations in the neutron spectra

	<u>ZPR-3 (6F)</u>	<u>Beryllium Slab</u>
Slab Half Width (x_0)	17.50	25.00
Source Half Width (x_1)	2.00	2.50
Diffusion Length	9.79	7.085
Mid Point or Off-Axis Source (x_2)	12.50	12.50
Detector Mid Point (x_3)	10.00	
Detector Half Width (x_4)	1.00	12.50
Fundamental		1.00
Buckling Mode (B_0^2)	<u>.00704</u>	<u>.00093</u>

All measurements in centimeters

TABLE II. Physical Parameters of Test Systems.

resulting from the inclusions of the higher harmonic modes. The persistence of the higher modes was observed in the early time history, but provided a negligible contribution toward the end of the decay time.

B. DATA PRODUCTION

The neutron probability density data for the lethargy-time response was calculated by MOD-5 for all harmonic modes using two limits on the systems considered. MOD-5 normally selects those discrete time steps to provide the probability density data under its internal selection criteria as though all problems were a fundamental mode type calculation. This freedom in selecting the output time steps was over ridden by modification of the existing control sequency in the basic MOD-5 program.

The fundamental mode calculation was done for a system to obtain the fixed time intervals and output time steps for the higher harmonic mode calculations. With the "locking in" of the data output time steps, the higher harmonic modes were run to the same final discrete time as the fundamental mode. Unfortunately, as the higher modes decay more rapidly, this resulted in the abnormal termination of several computer runs when the probability density dropped to values smaller than the IBM-360/67 is capable of processing, i.e., below 10^{-75} . As a result, all final computer runs were set to terminate when the probability density component for a mode dropped below an arbitrary limit of 0.001.

The time decay of the harmonic components is mainly affected by the leakage of neutrons from the system which is dependent on the buckling factor, DB_n^2 , which appears in the time dependent diffusion equation as a term:

$$vD(v)B^2N(x,v,t).$$

A summary of the MOD-5 calculations is given in Table III for the two basic systems where all calculations were carried to the point where the probability that a neutron had leaked from the system, been absorbed or had dropped below the lowest energy was .999. In all cases, the higher harmonics were still providing a small contribution to the total population at times approaching 300ns, but the major contribution was due mainly to the fundamental.

C. ANALYSIS OF EXPANSION METHODS

The adequacy of the truncated Fourier series to represent the various source geometry conditions was evaluated by the comparison of the source strength parameter which was determined via the form:

$$\Delta_o = 1.0 - \int_{-x_o}^{x_o} \sum_1^N (A_n \text{Cos } Kn x + C_n \text{Sin } Kn x) dx$$

$$\Delta_o = 1.0 - \sum_1^N (-) \frac{A_n 2}{kn} \text{Sin } Kn x_o - \frac{C_n}{kn} \text{Cos } Kn x_o.$$

The uncorrected source strength error (Δ_o) was evaluated for the first 100 terms of the series expansion; which stabilized to almost constant-values after the first 30 to 40 terms.

System	Mode	Time (ns)	Final Neutron Probability Slowing to Bottom of Spectrum**	Final Neutron Probability Components (Percentage)		
				Leakage	Capture non-fission	Fission
Beryllium	1	16620	83.9	7.4	7.9	-
	2	9834	44.7	48.0	6.8	-
	3	4983	14.2	80.3	5.4	-
	4	1952	3.1	92.6	3.4	-
	5	1315	.01	95.6	3.4	-
	6	91	.002	96.1	2.9	-
ZPR-3 (6F)	1	311	.95	51.1	8.7	39.2
	2	311	.001	89.2	1.2	9.5
	3	311	.0	88.6	1.2	10.2

** Bottom of Spectrum is defined as:

Beryllium Slab below 2.31 eV

Assembly 6F below .4 eV

TABLE III. Summary of MOD-5 Computer Runs.

Uncorrected Source Error (Δ_0) 100 Term

Source Type	ZPR-3 (6F)	Beryllium Slab
1- Square Pulse	0.0652	0.0265
3- Exponential	0.0231	0.190
5- Ramp Function	0.0090	0.0090

The Fourier series expansions are defined such that with an infinite number of terms, the error should go to zero. However, calculations show that for the source function approximations, this error goes to some non-zero value. It should not be any surprise that these source terms have a non-zero error, since the Fourier approximation is based on the interval, $x_0 + \delta$, where the distance " δ " represents the extrapolation distance correction for the boundary condition of forcing the neutron density to zero at that point.

In the source function testing, the physical limits of the slab are within a width $2x_0$, rather than boundaries of $2(x_0 + \delta)$. This additional value of twice the extrapolation distance is responsible for apparent constant error. When the integration is done within the bounds of the slab, a loss occurs at the end points in the intervals:

$$-(x_0 + \delta) \leq x' \leq -x_0 \quad \text{and} \quad +x_0 \leq x \leq (x_0 + \delta).$$

As the ratio δ/x_0 is greater than zero for all real (non-infinite systems), this apparent difficulty in the approximation will be almost negligible for systems with the slab width much greater than the extrapolation distance ($\delta/x_0 \ll 1.0$).

Each of the three main source types (the square pulse, exponential and ramp functions) were studied to the limit of

a 100 mode harmonic expansion for the spatial dependent functions. The value of the delta correction (Δ_0) are shown in Figures 5 and 6 for both systems. While all test calculations were done for a six mode expansion, the effect of the truncation to less than an infinite series approaches an asymptotic limit in each source type, the limits on the value of Δ_0 are a measure of the δ/x_0 ratio. In both of the systems studied, it can easily be seen that inclusion of more than 10 harmonic modes would not significantly change the spatial dependent values.

D. COMPARISON OF SOURCE TYPES

The development and testing of the time-space-energy-dependence for five source geometries was a task more formidable than originally anticipated. Numerous programming difficulties arose toward the end of the project for the two cases of the off-axis sources, which resulted in their not being included in the final calculations. Both of these sources require modification of the spatial harmonic functions than those proposed previously.

The three source conditions that are included in the detailed calculations were the symmetrical square pulse and the first collision source centered about the mid-point of the slab and the exterior wide beam or ramp function source.

The existing version of the computer program, MIL-6, demonstrated the effectiveness of the three main source conditions, but the off-axis square pulse and first collision geometry must be revised and retested prior to further analysis.

FIGURE 5
SOURCE FUNCTION ERROR
BERYLLIUM SLAB

- - SOURCE = 1
- x - SOURCE = 3
- + - SOURCE = 5

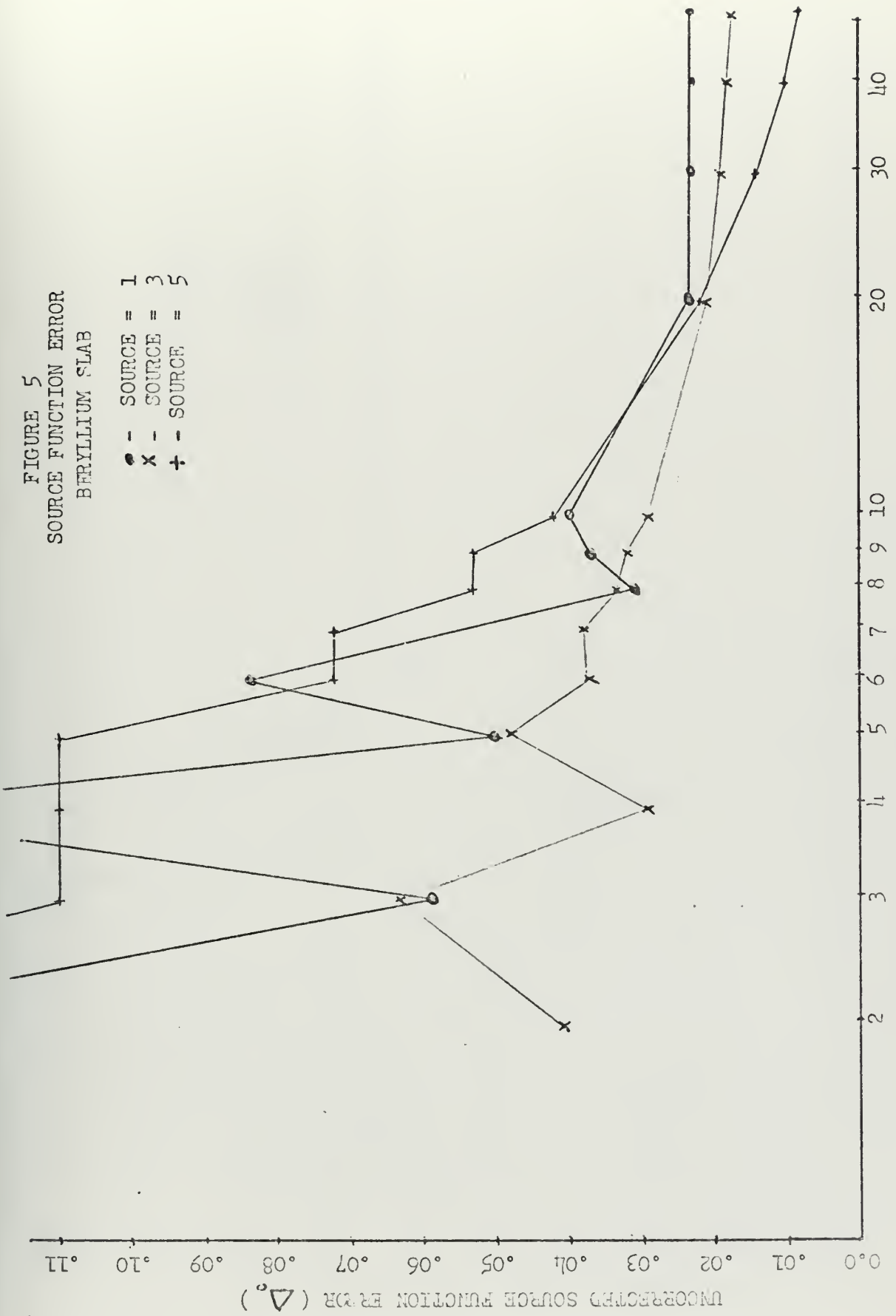
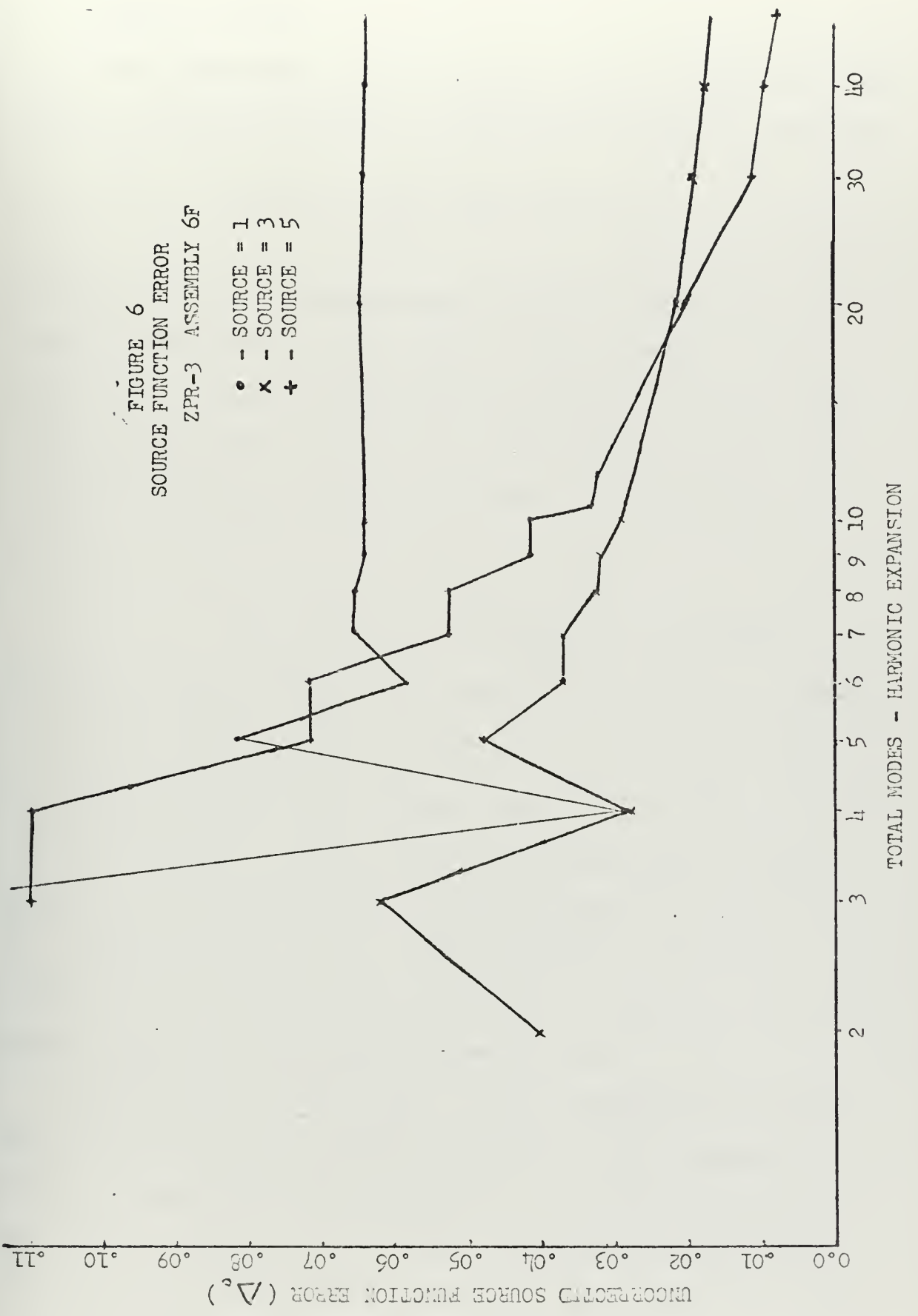


FIGURE 6
SOURCE FUNCTION ERROR
ZPR-3 ASSEMBLY 6F

- - SOURCE = 1
- x - SOURCE = 3
- + - SOURCE = 5



E. NEUTRON SPECTRA

The primary goal of this project was to demonstrate the spatial dependent response of the neutron spectral function:

$$N(x,v,t) = \sum_n R_n(x) F_n(v,t).$$

The contribution of each harmonic mode to the neutron spectra is determined by a Fourier space harmonic term, $R_n(x)$, and the MOD-5 time-lethargy probability density for the n-th time mode, $F_n(v,t)$. For each harmonic mode, the probability density function is subject to the basic condition:

$$F_n(v,t) = 1.0 - (\text{Absorption} + \text{Leakage})$$

and the initial neutron source spatial distribution provides the correct source contribution.

F. NEUTRON DENSITY AND NEUTRON FLUX

The results of neutron density and neutron flux calculations are presented at a limited number of points to be representative of the over all processes, however, inclusion of all points in the slab would result in an inordinate number of graphs and tables of data.

The neutron density, $N(x_k, v_i, t_j)$, was plotted only for the point $x_k=0.0$ and a complete spectral plot is given in Figures 13 and 16, to be representative of the information possible to determine. All other calculations made by the program MIL-6 required the evaluation of the density at all points in the slab, and similar response curves were prepared.

The total neutron flux,

$$\phi(x_k, t_j) = \sum_i v_i N(x_k, v_i, t_j),$$

is presented for one location half way between the center and the surface of the slab, the same positions selected for the detector response calculations. Each plot is for a single system, either the Beryllium or assembly 6F, follows the complete time decay of the total flux from one nanosecond to the time at which the neutron population had decayed to 0.1% of its initial value. In both cases, the first build up of the flux indicates the initial stages of the neutron motion within the slab, to the point at 20ns at which the leakage from the free surface becomes the dominant effect.

When the leakage terms are considered, the appearance of the decay curve now looks almost identical with the pseudo decay modes previously described. When looking at the neutron density curves, it can easily be noticed that it required several nanoseconds for the initial 2.46 MeV burst to develop into a smooth distribution.

G. MEAN ENERGY AND MEAN ENERGY RATIO

A convenient method of comparing the time and spatial effects on the neutron density is by comparing the average or mean energy of the neutrons at all points in the slab at the same time. A complete history of the neutron population can be concisely presented by observing the absolute variation of the mean energy at a single reference point in

the slab and then comparing the relative changes in the entire slab in terms of the reference point for each discrete time step.

The basic reference point for all comparisons is the center of the slab, $x=0.0$, for all geometries and systems. The decay of the mean energy with time is shown for both beryllium (Figure 19) and assembly 6F (Figure 20). Starting with the two reference cases, the spatial variation from the center to the surface of the slab starting at two nanoseconds after the pulse to the final decay of the fundamental mode, (beryllium - Figures 21 to 30--assembly 6F Figures 31 to 36).

After the initial pulse, a wave-like shape is observed to develop across the slab in the direction of the free surface, as one would expect the higher energy neutrons to migrate out of the system fastest. This is shown most vividly in Figures 24 to 27.

Near the end of the decay period, the last remaining neutrons in the system have almost a constant ratio of one, as the system tends toward an equilibrium condition. The leakage at the surface allows the final stages to show this condition (Figures 30 and 36).

H. DETECTOR RESPONSE

This computational routine did not work as well as intended. Minor computational and programming errors were encountered very late in the project and have not been completely resolved. These will require revision and further testing as follow on work.

Comparisons of the flux calculations and the mean energy determinations referred to in Sections F and G were compared to the "proton recoil" detector (Figures 37 and 38) which follow the same basic trends for the time dependent response, but the numerical differences in the time integrated response will require improvement of the space integral and time summation numerical methods.

VI. RECOMMENDATIONS

A. FOLLOW ON WORK

The program MIL-6 was prepared to be a general computational tool using the harmonic mode expansion method for the one dimensional systems. Principal follow on work is recommended to correct the minor computational difficulties encountered, write and test data transfer routines for both MOD-5 and MIL-6 and review and edit the programming for a more optimized execution.

Principal areas to have future work are:

- (1) re-evaluate the Fourier space functions for the two off-axis source geometries,
- (2) revise and test the numerical integration techniques of the detector response program for the summation over time, and consider methods to improve the space integration over the detector volume,
- (3) consider a detector response function that would be more realistic than the STP hydrogen case now included in the program,
- (4) complete and test the input/output options using nine-track magnetic tape and the IBM-2314 magnetic disc systems for data transfer between the programs MOD-5 and MIL-6,
- (5) revise the program MIL-6 to handle more than six harmonic modes by use of a smaller state structure of 70

to 100 states rather than the 150 states the program is currently written to process,

(6) provide more realistic source geometries for the one dimensional case than the five simple ones currently in the program,

(7) review the program for reduction in the core storage space requirements and to improve the speed of execution through program optimization.

B. IMPROVEMENTS TO PROGRAM MIL-6

The original program MIL-6 was written to deal with monoenergetic neutron sources and delta functions in time. This may require significant modification of the computational methods to considering the fission spectrum sources and sources that have a finite width time duration. These situations had not been tested in the MIL-6 work and would provide a more realistic model for fast neutron experiments.

A second item that appears worthy of consideration in the future will assist in modeling the wave-like motion of the neutron pulse from the finite volume square pulse and the exterior wide beam ramp function. This could be considered as a limiting correction factor to couple the time-space and energy dependence during the initial response to a neutron pulse.

C. THREE DIMENSIONAL SYSTEMS

The harmonic expansion method has been demonstrated to work for the one dimensional case, and definite action to

begin study of three dimensional systems is warranted. A recommended sequence of priorities would be to expand the MIL-6 work to three dimensions in the cartesian coordinates first, cylindrical geometry second and spherical geometry coordinate systems last.

APPENDIX A: NUMERICAL RESULTS

A. SPECTRAL RESPONSE

1. Beryllium Slab

a. Neutron density harmonic component $N_n(x, v_i, t_j)$ for harmonic modes 1 to 6 at point $x=0.0$. (Figures 7 to 12.)

b. Neutron density: $N(x, v_i, t_j)$: lethargy - time neutron population density versus lethargy for a 6 mode expansion at point $x=0.0$. (Figure 13.)

2. ZPR-3 Assembly 6F

a. Neutron density harmonic component $N_n(x, v_i, t_j)$ for harmonic modes - 1 and 2 at point $x=0.0$. (Figures 14 and 15.)

b. Neutron density $N(x, v_i, t_j)$ - lethargy - time neutron population density versus lethargy - for a sum of 3 harmonic modes at $x=0.0$. (Figure 16.)

B. MEAN ENERGY AND MEAN ENERGY RATIO

1. Beryllium

a. Mean energy versus time at $x=0.0$. (Figure 19.)

b. Plot of mean energy ratio versus position for 8 time periods. (Figures 21 to 30)

2. ZPR-3 (6F)

a. Mean energy versus time at $x=0.0$. (Figure 20.)

b. Plot of mean energy ratio versus position for 6 time periods. (Figures 31 to 36.)

C. DETECTOR RESPONSE

1. Beryllium - Figure 37.
2. ZPR-3 - Figure 38.

FIGURE 7
 NEUTRON DENSITY
 BERYLLIUM SLAB
 FUNDAMENTAL MODE
 $X=0$

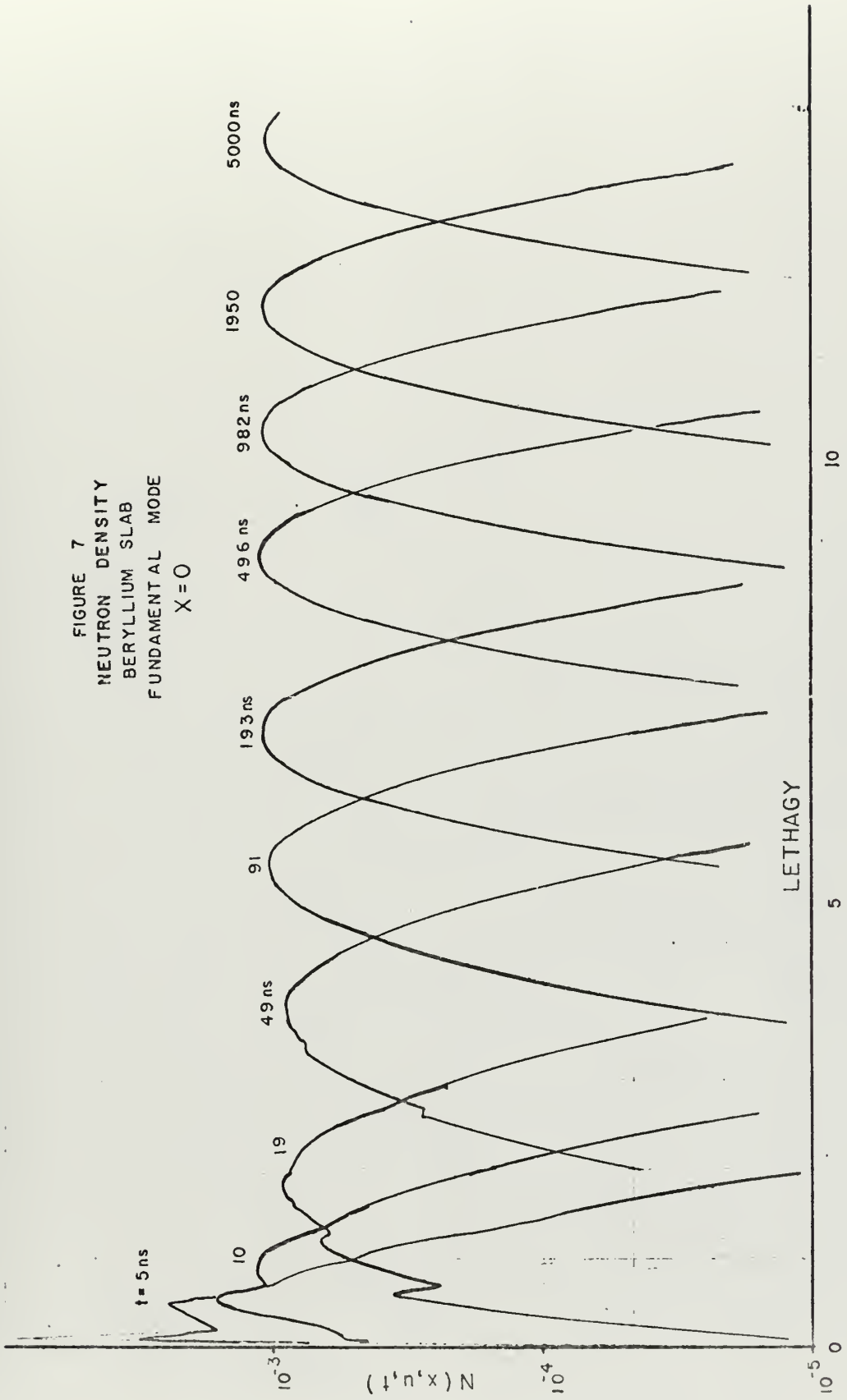


FIGURE 8
 NEUTRON DENSITY
 BERYLLIUM SLAB
 FIRST HARMONIC MODE
 $X=0$

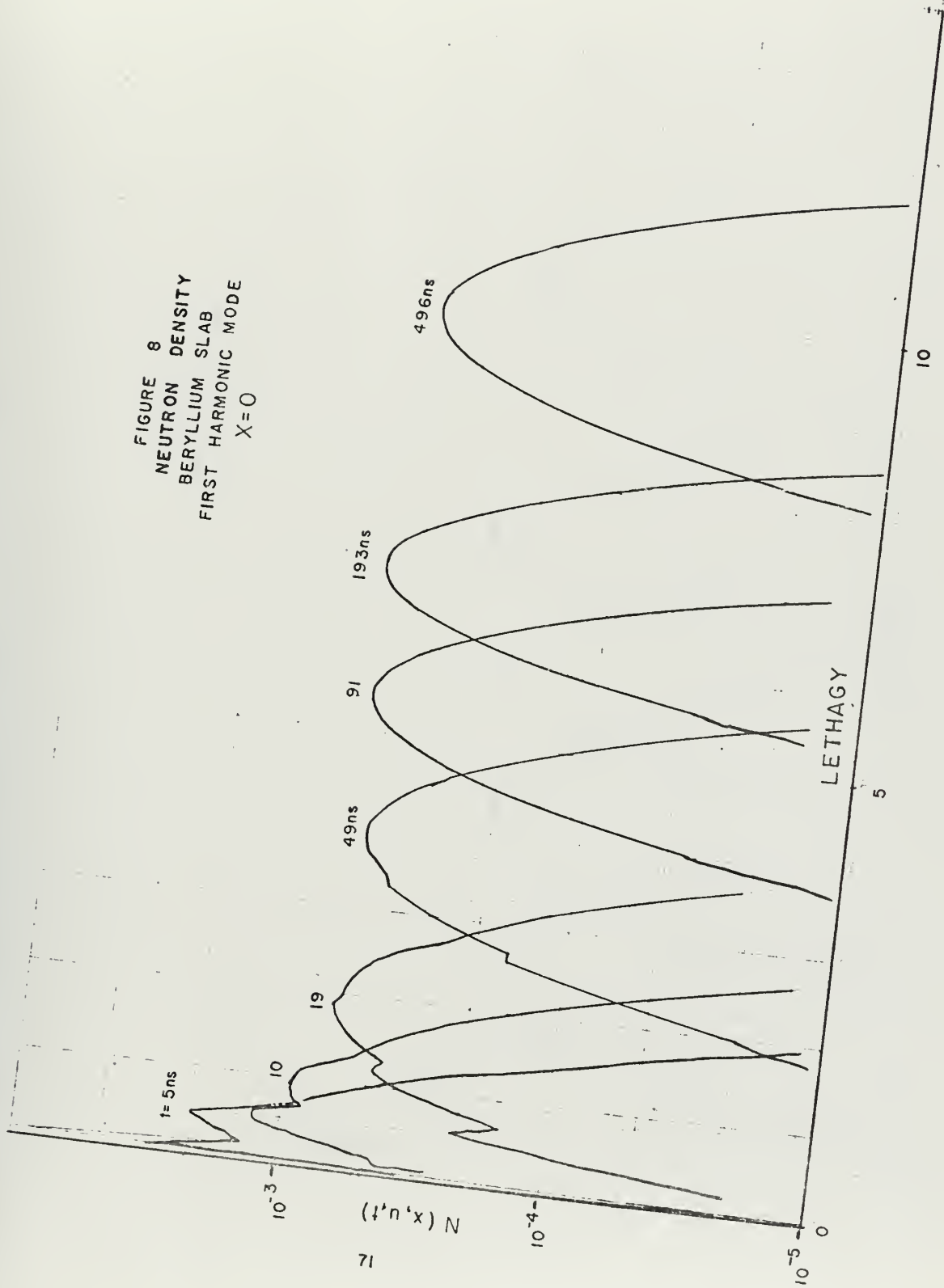


FIGURE 9
 NEUTRON DENSITY
 BERYLLIUM SLAB
 SECOND HARMONIC MODE
 $X = 0$

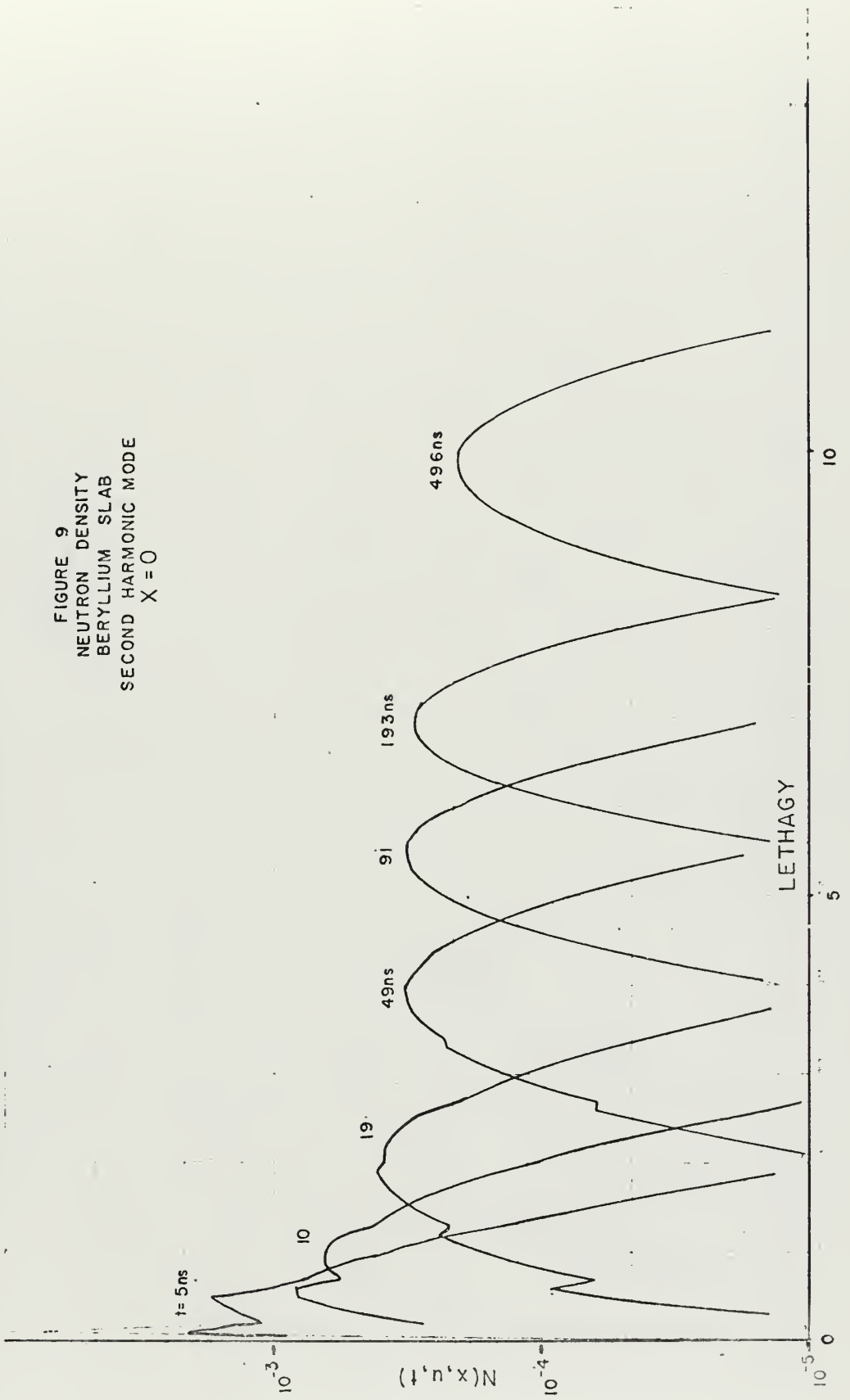


FIGURE 10
 NEUTRON DENSITY
 BERYLLIUM SLAB
 THIRD HARMONIC MODE
 $X = 0$

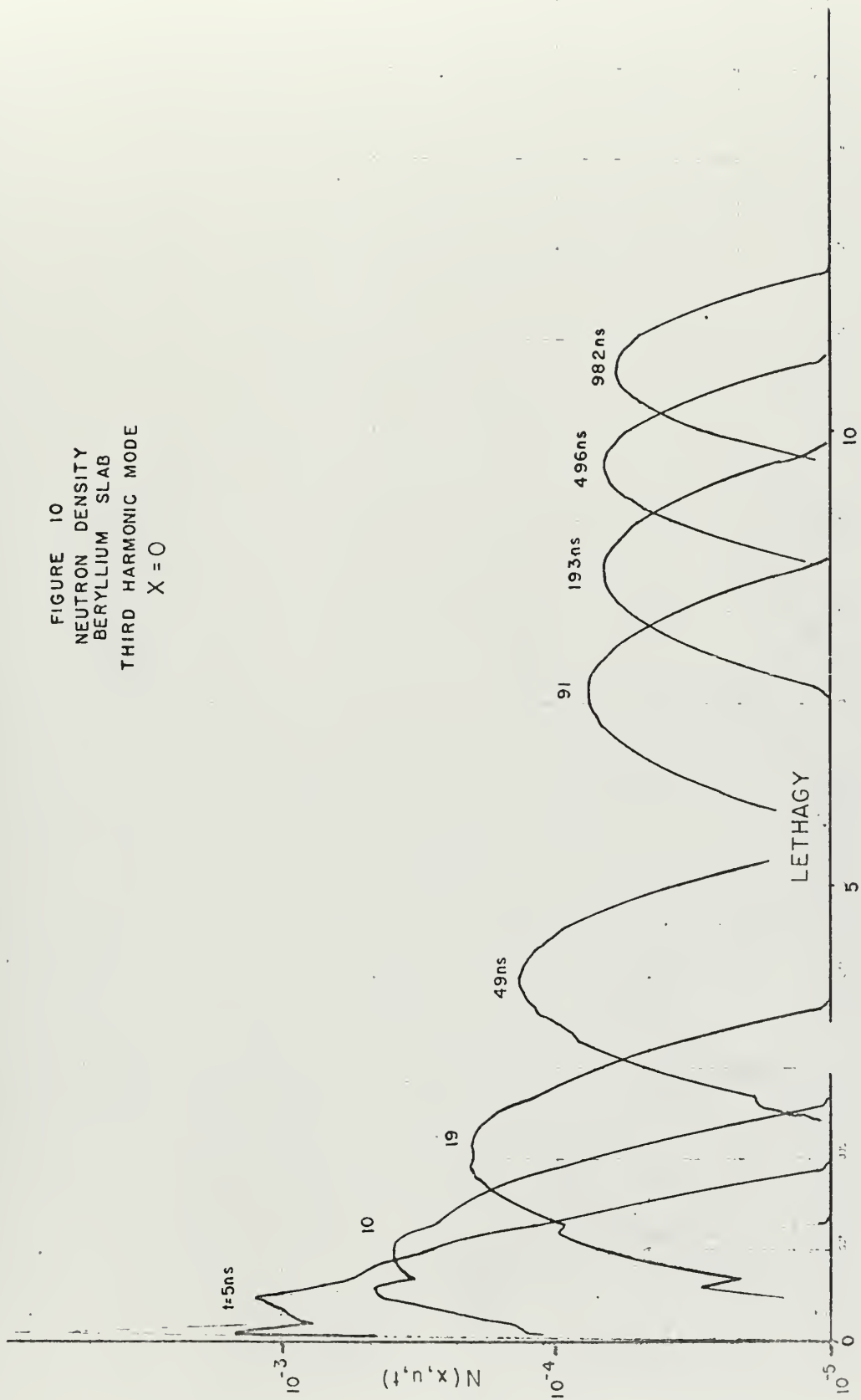


FIGURE 11
 NEUTRON DENSITY
 BERYLLIUM SLAB
 FOURTH HARMONIC MODE
 $X = 0$

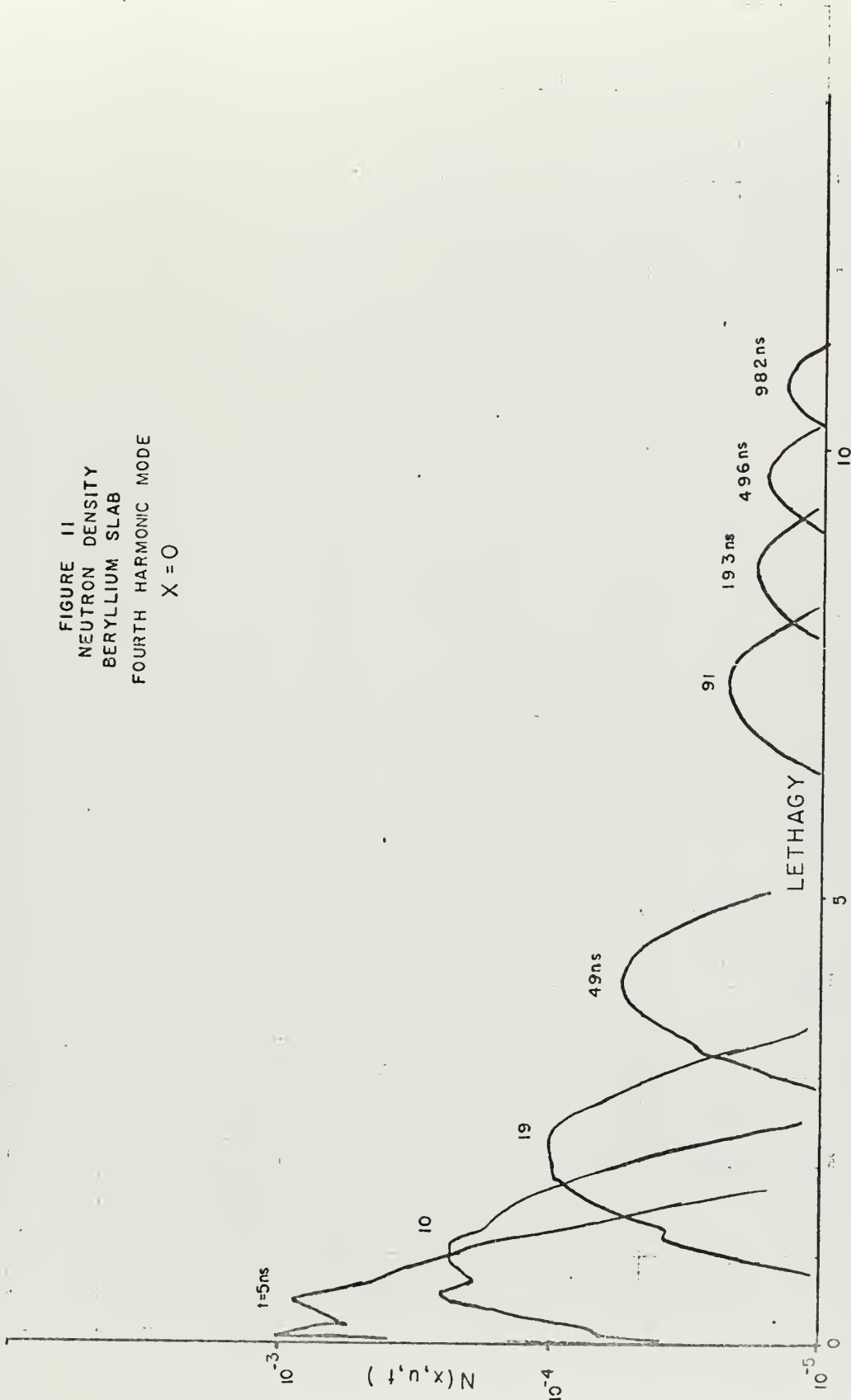


FIGURE 12
 NEUTRON DENSITY
 BERYLLIUM SLAB
 FIFTH HARMONIC MODE
 $X = 0$

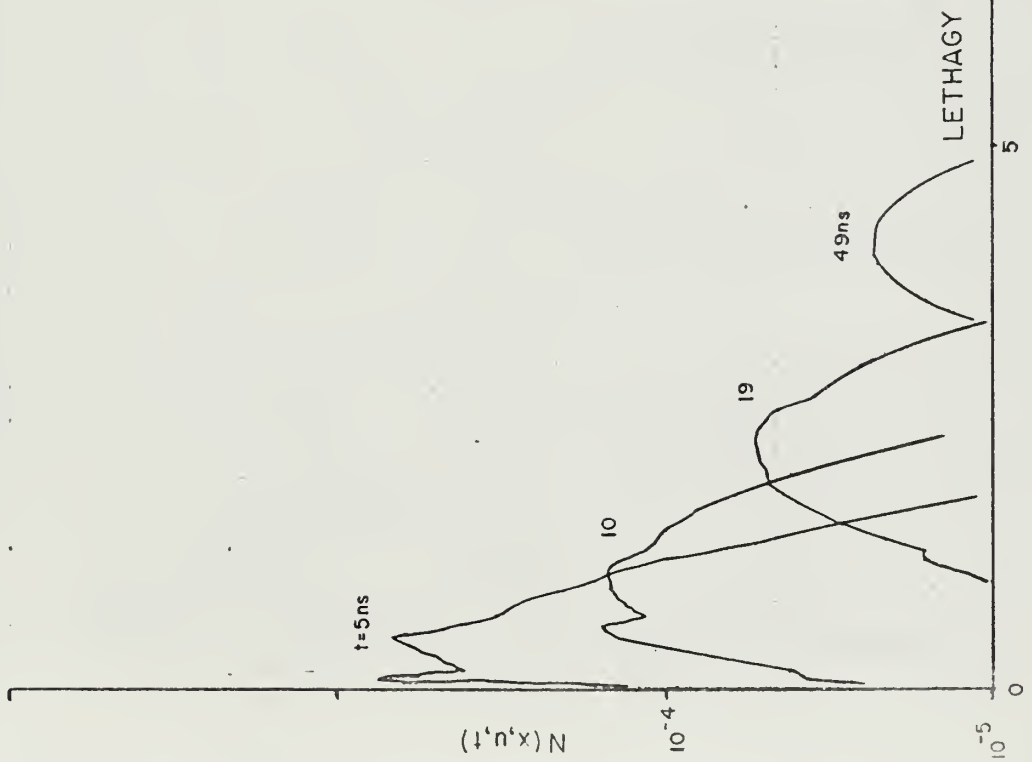


FIGURE 13
 NEUTRON DENSITY
 BERYLLIUM SLAB
 SUM OF ALL MODES

$X=0$

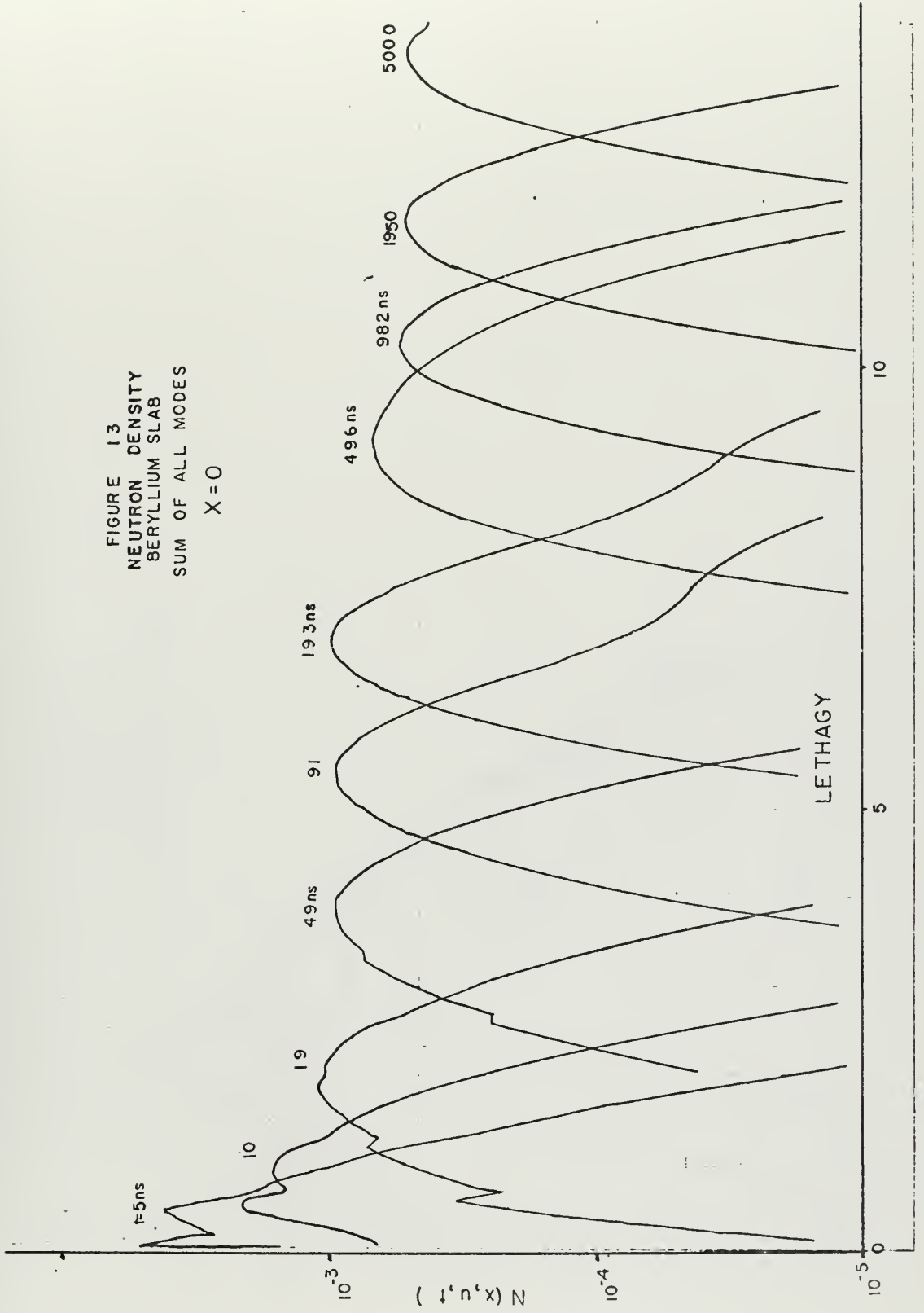


FIGURE 14
 NEUTRON DENSITY
 ASSEMBLY 6F
 FUNDAMENTAL MODE
 $X=0$

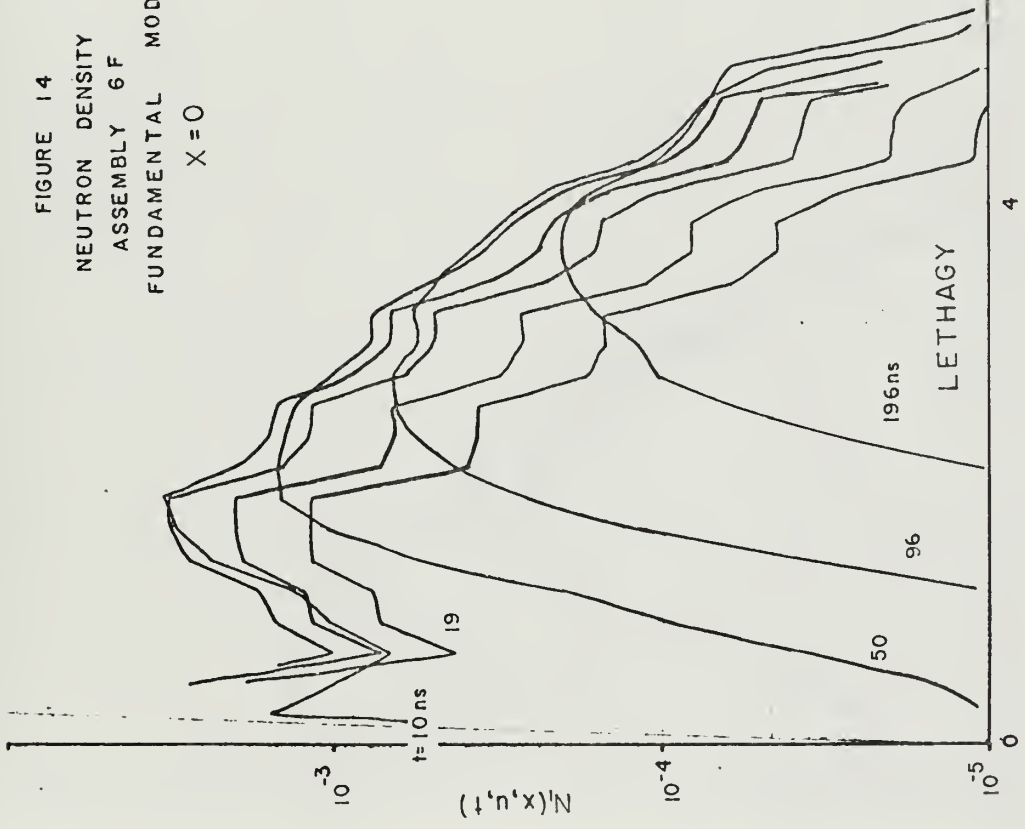


FIGURE 15
 NEUTRON DENSITY
 ASSEMBLY SF
 FIRST HARMONIC MODE
 $X = 0$

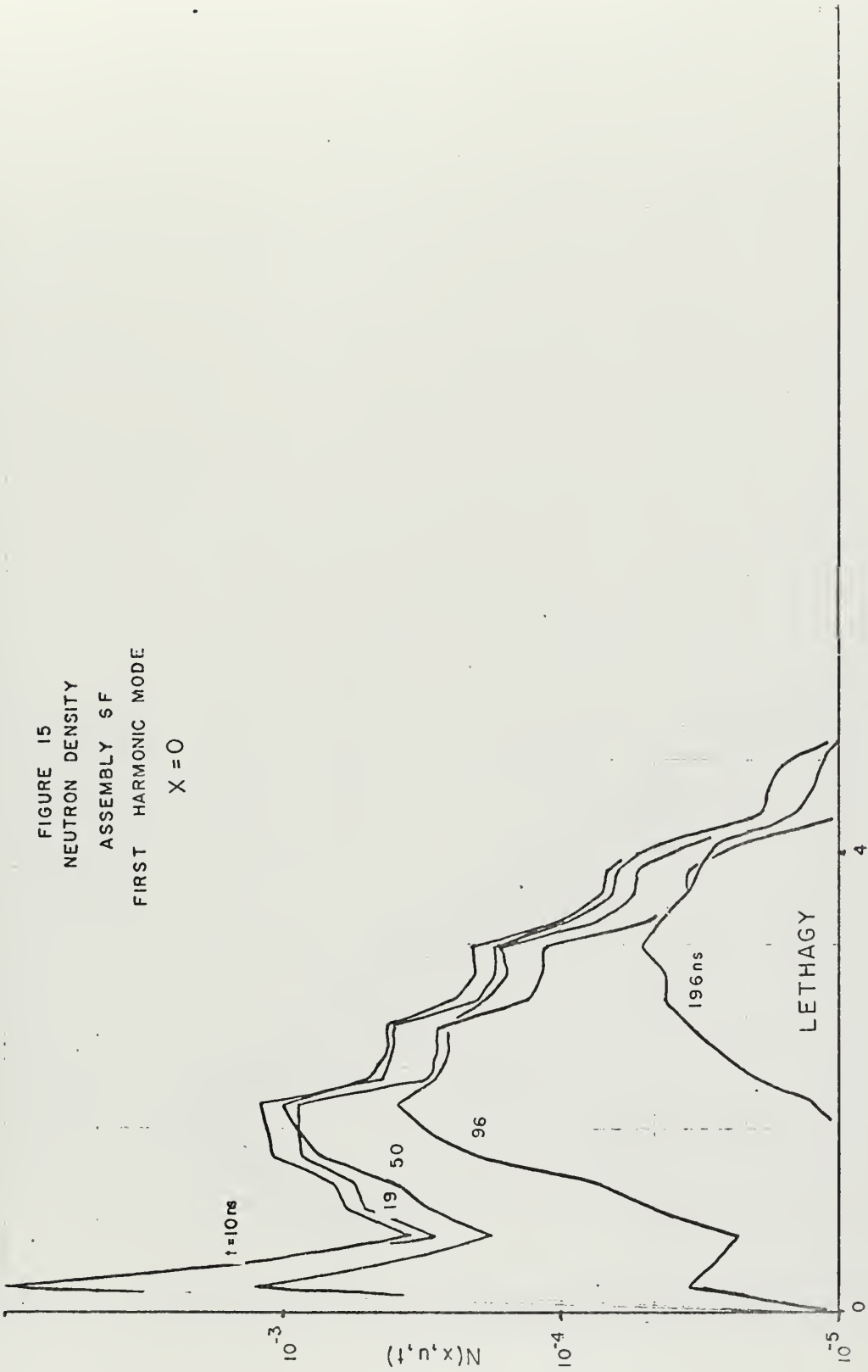


FIGURE 16
 NEUTRON DENSITY
 ASSEMBLY 6 F
 SUM OF ALL MODES
 X = 0

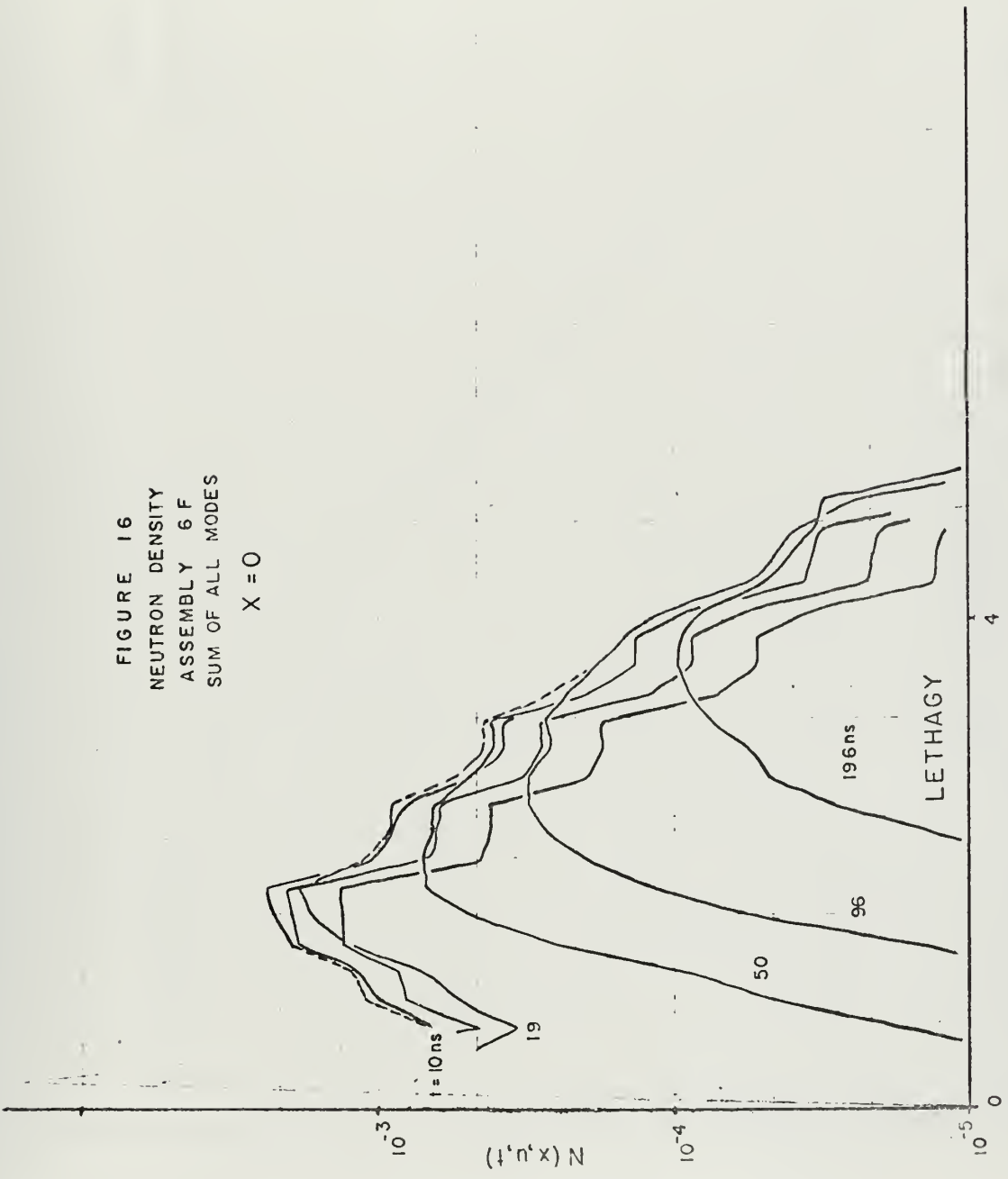


FIGURE 17
BERYLLIUM SLAB
NEUTRON FLUX AT $X = 0.0$

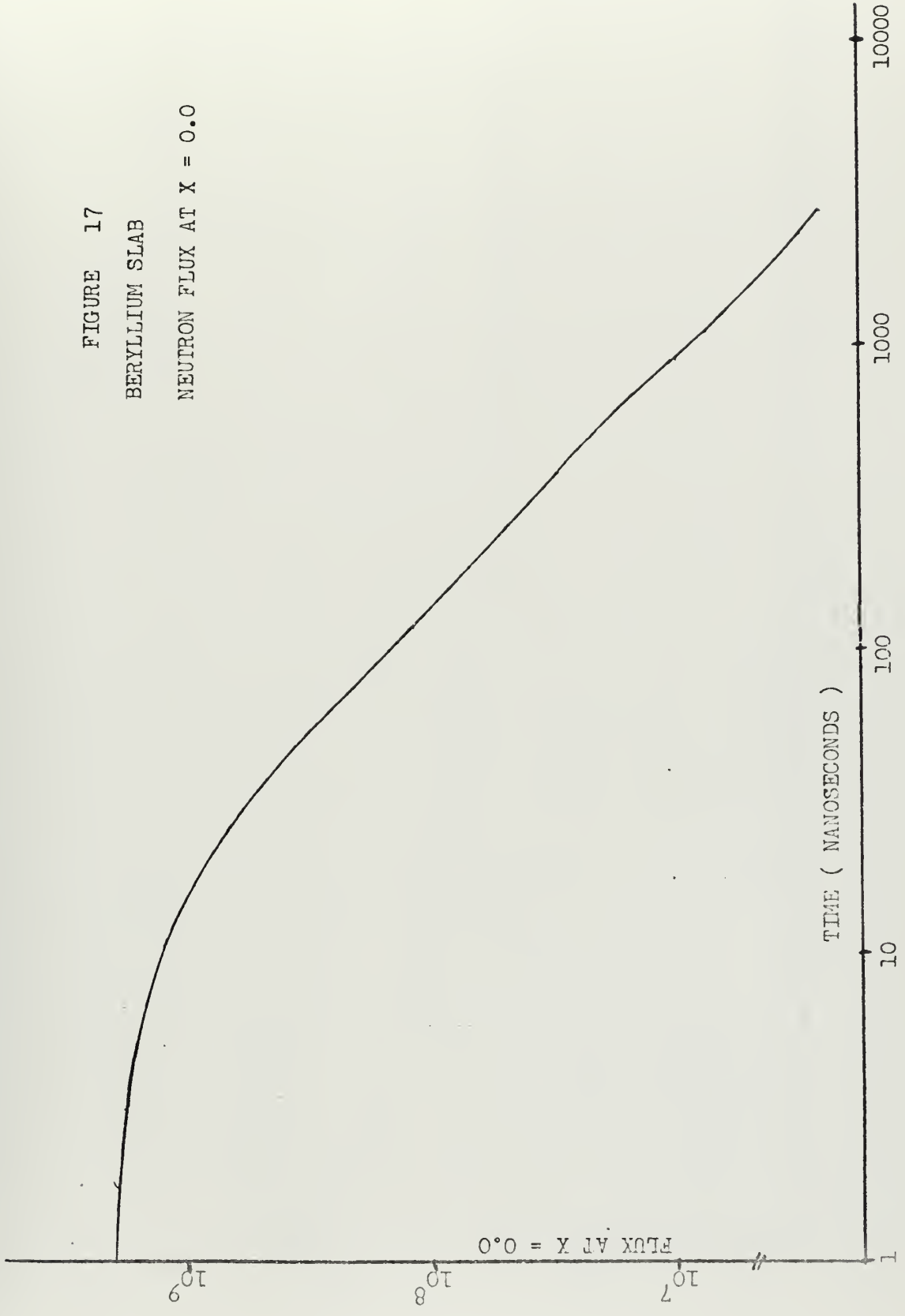


FIGURE 18
ASSEMBLY 6 F
NEUTRON FLUX AT X = 0.0

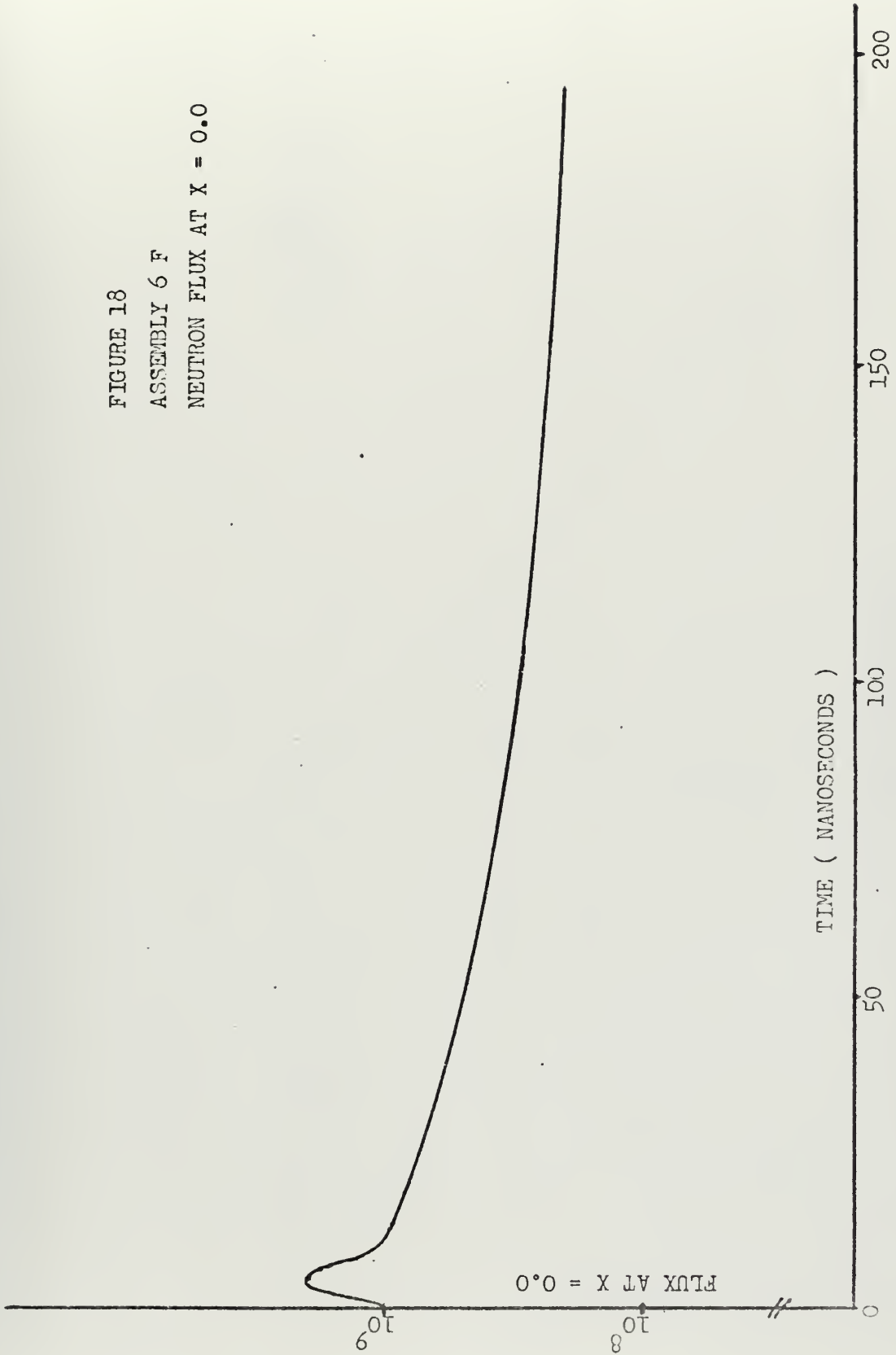


FIGURE 19
MEAN ENERGY AT $X = 0.0$
BERYLLIUM SLAB

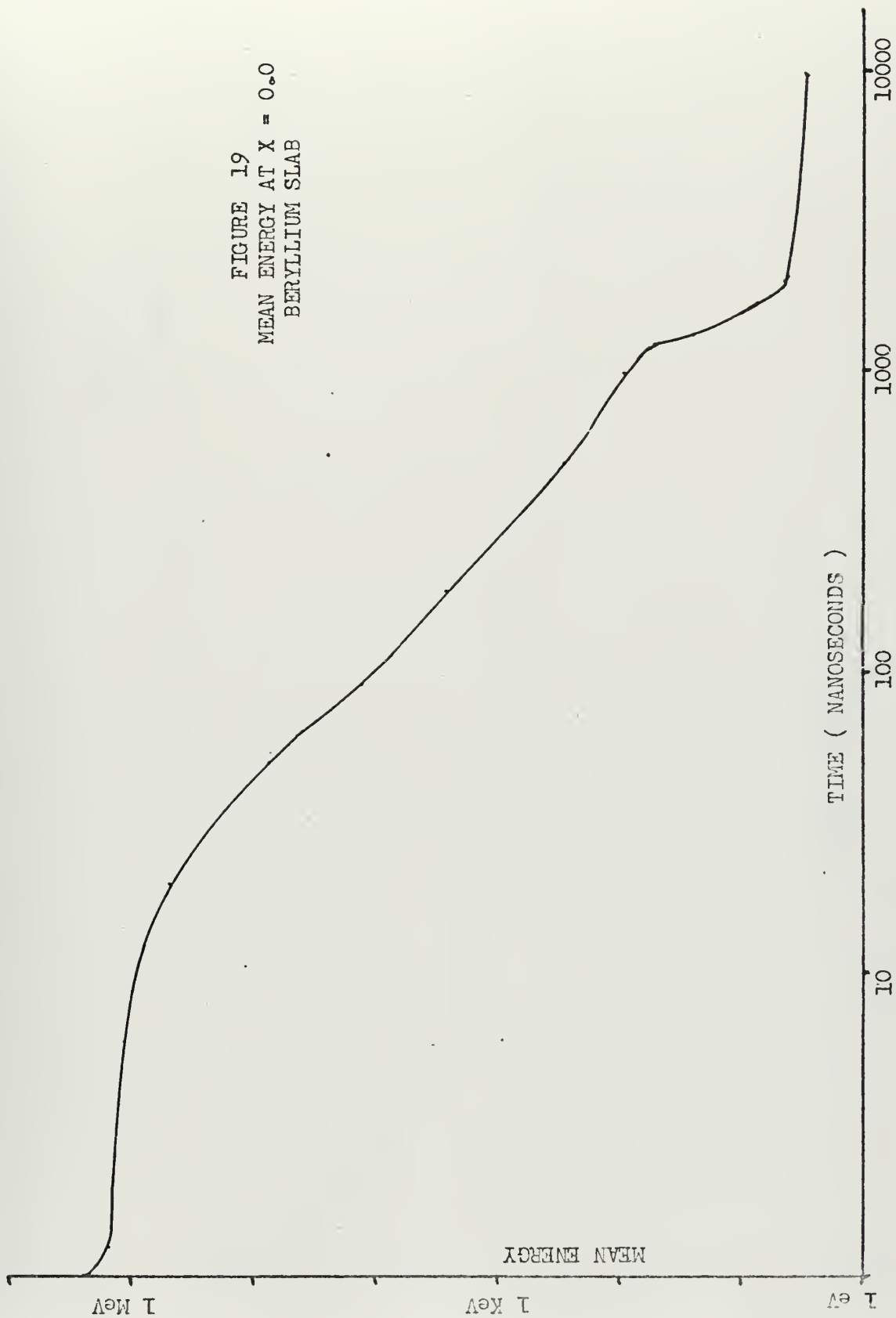
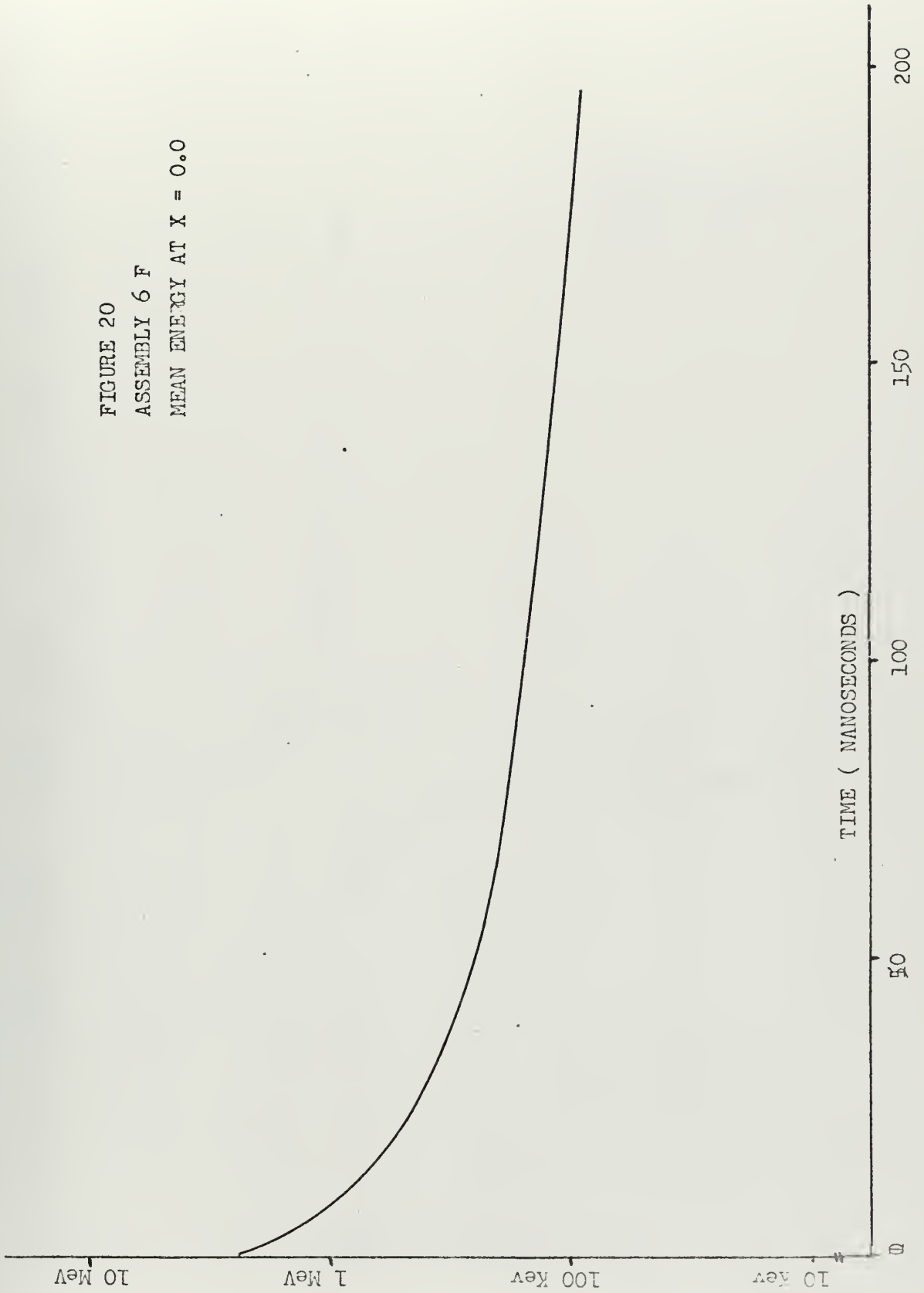


FIGURE 20
ASSEMBLY 6 F
MEAN ENERGY AT $X = 0.0$



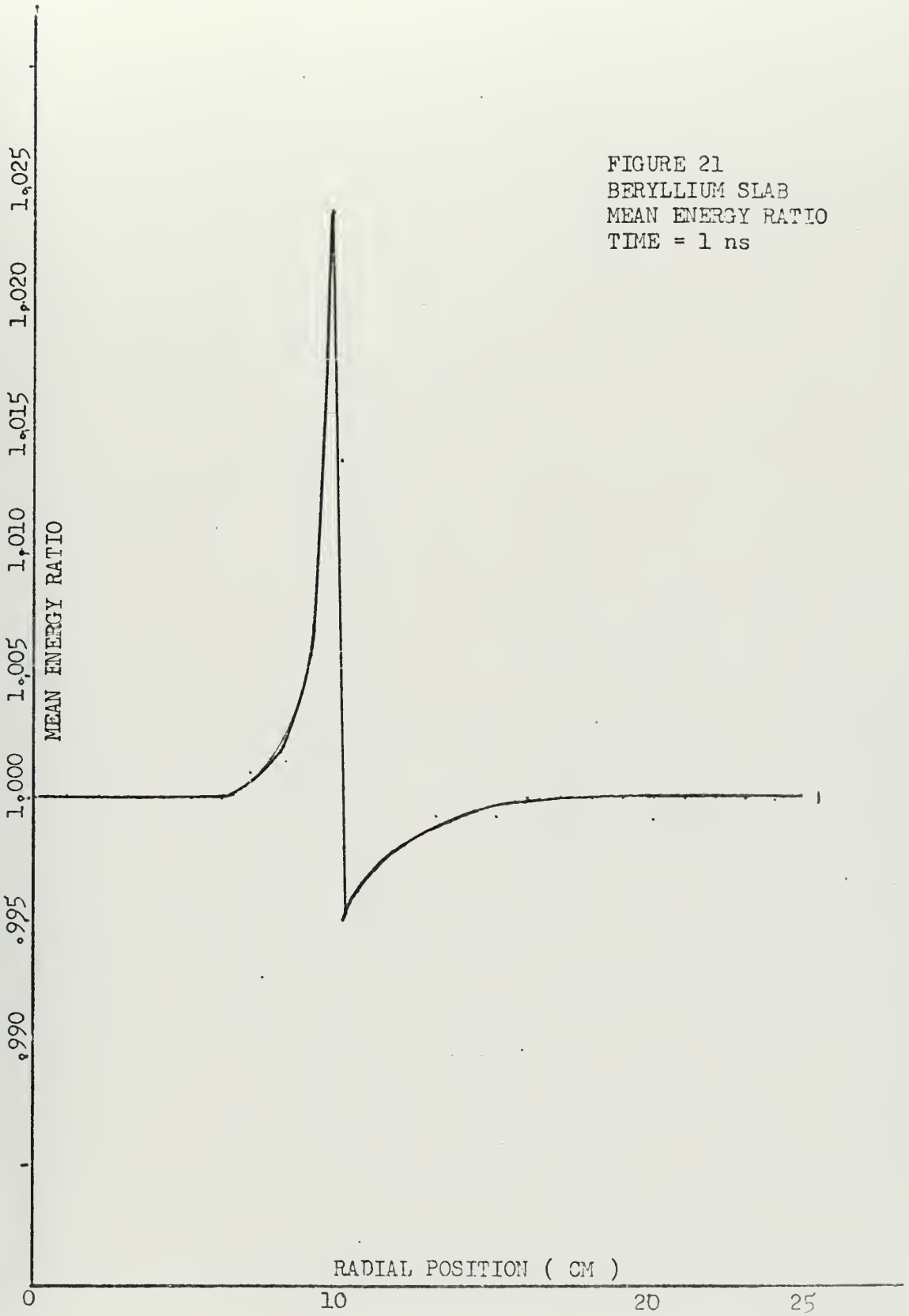
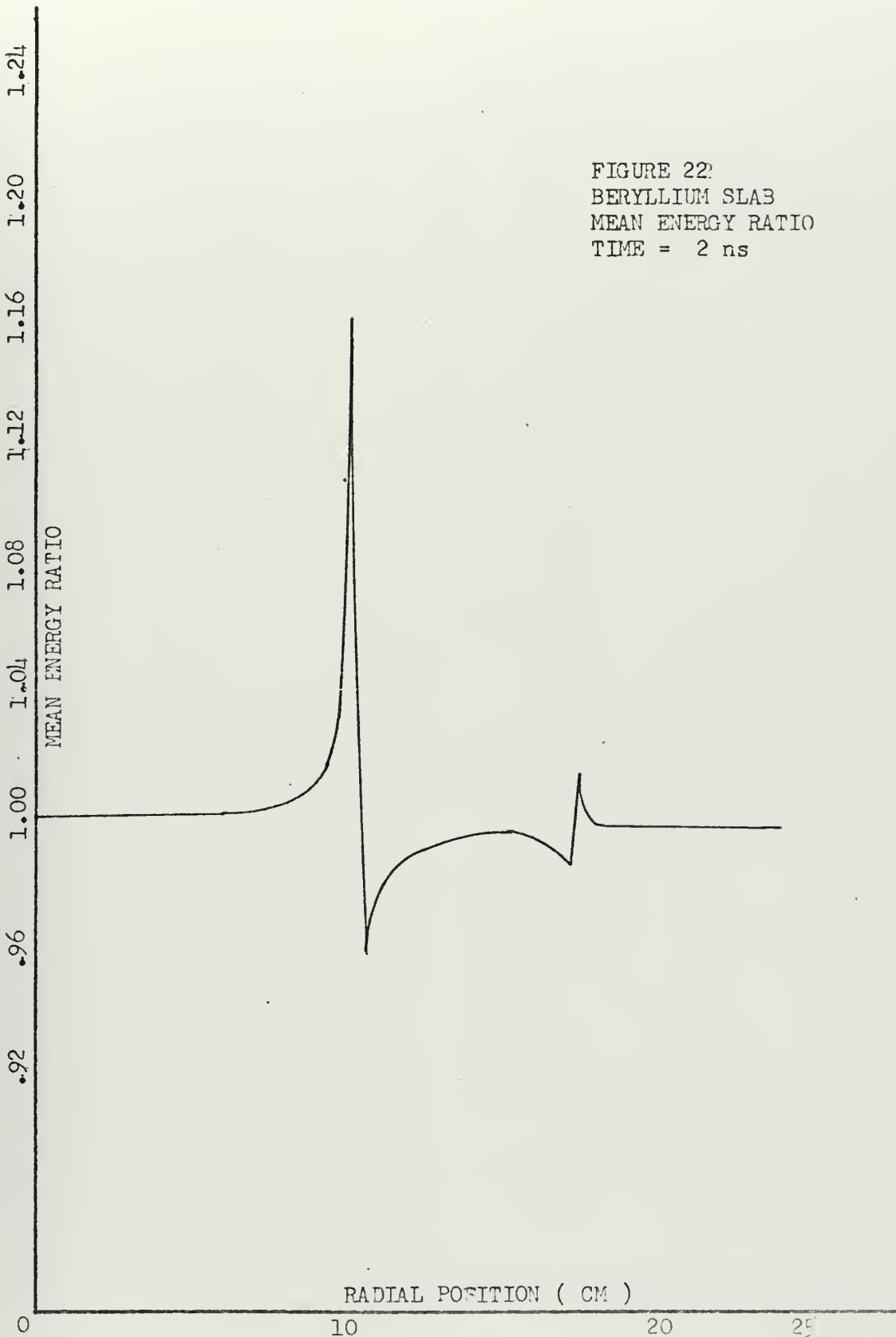


FIGURE 21
BERYLLIUM SLAB
MEAN ENERGY RATIO
TIME = 1 ns

FIGURE 22
BERYLLIUM SLAB
MEAN ENERGY RATIO
TIME = 2 ns



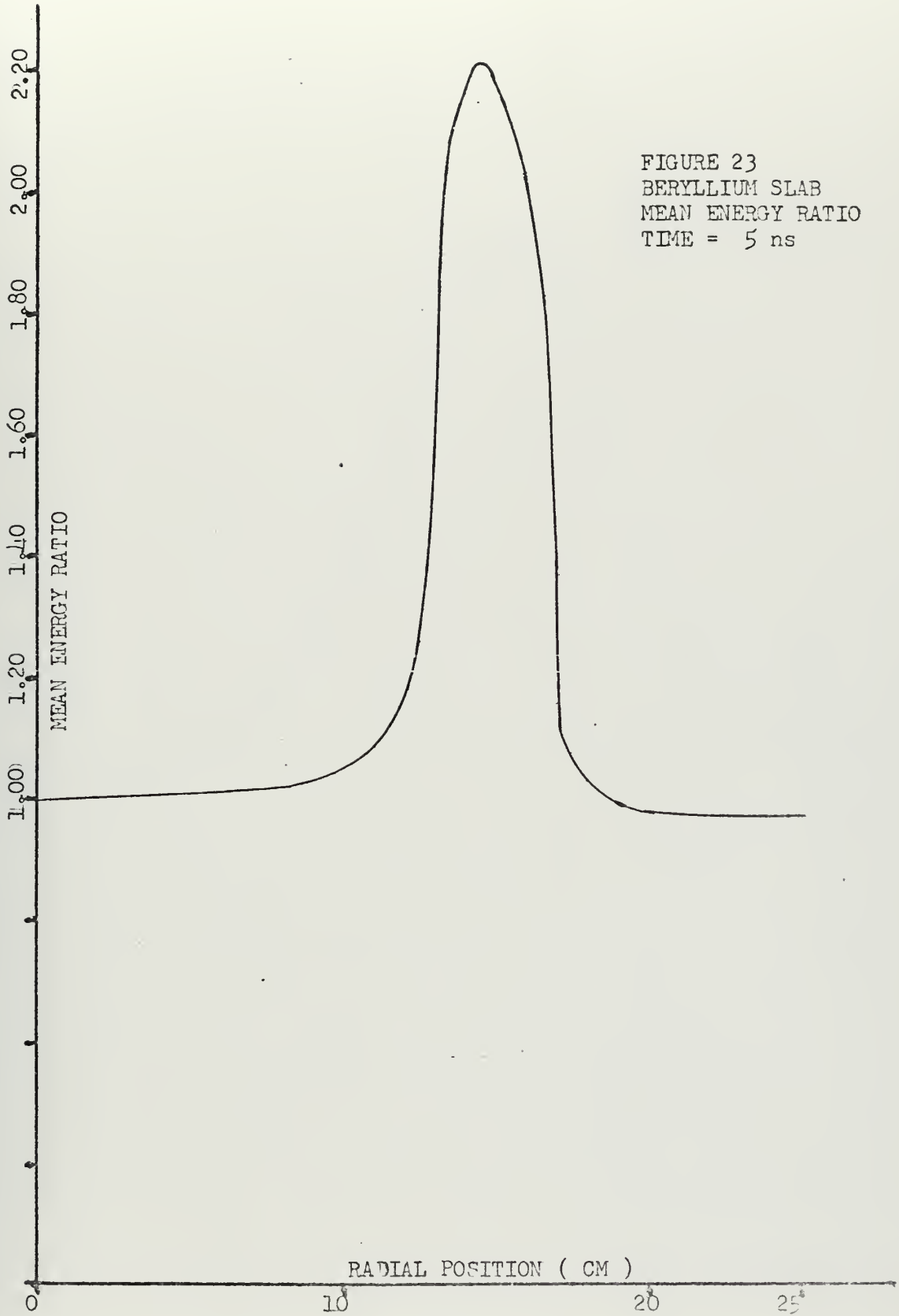


FIGURE 23
BERYLLIUM SLAB
MEAN ENERGY RATIO
TIME = 5 ns

FIGURE 24
BERYLLIUM SLAB
MEAN ENERGY RATIO
TIME = 10 ns

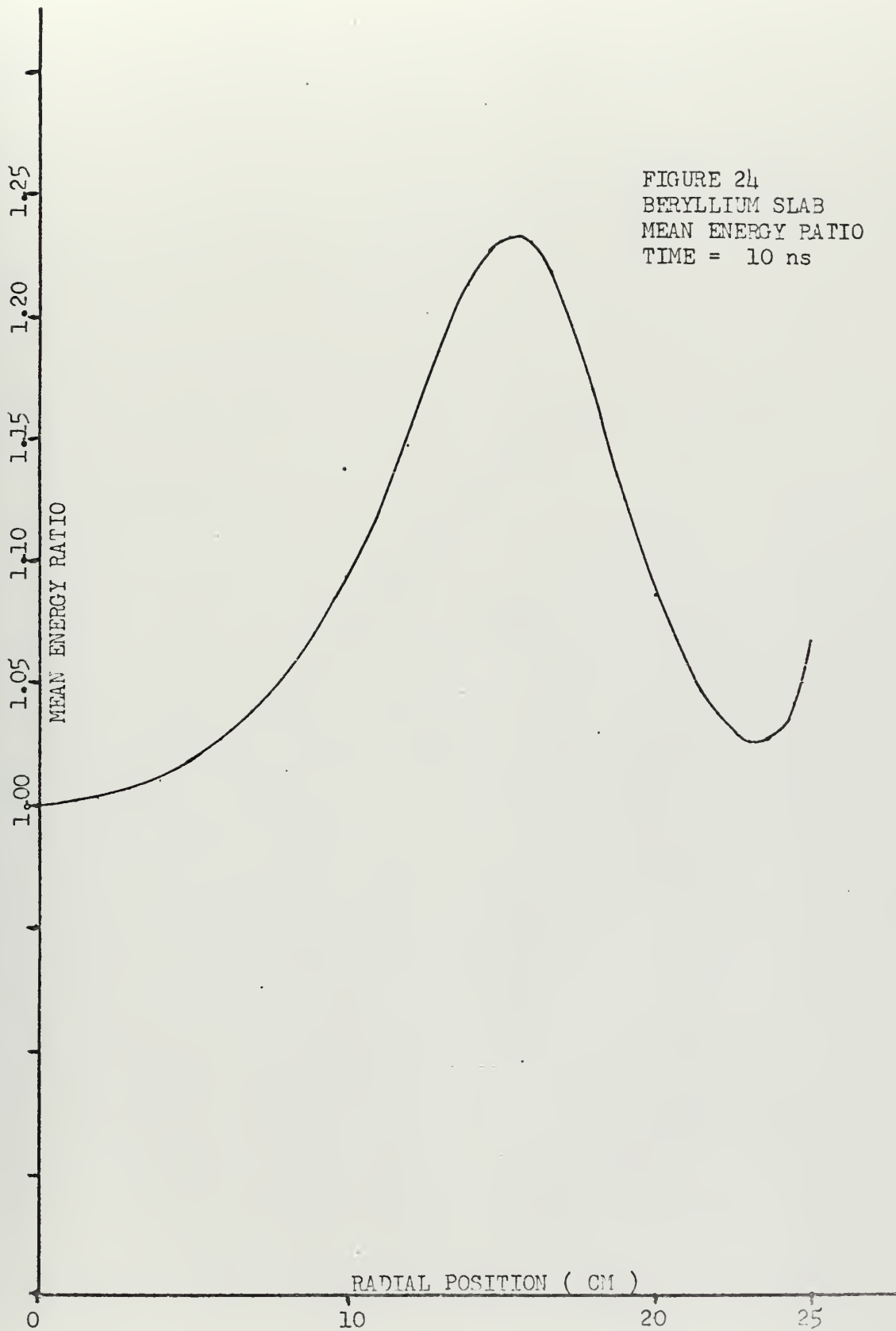


FIGURE 25
BERYLLIUM SLAB
MEAN ENERGY RATIO
TIME = 19 ns

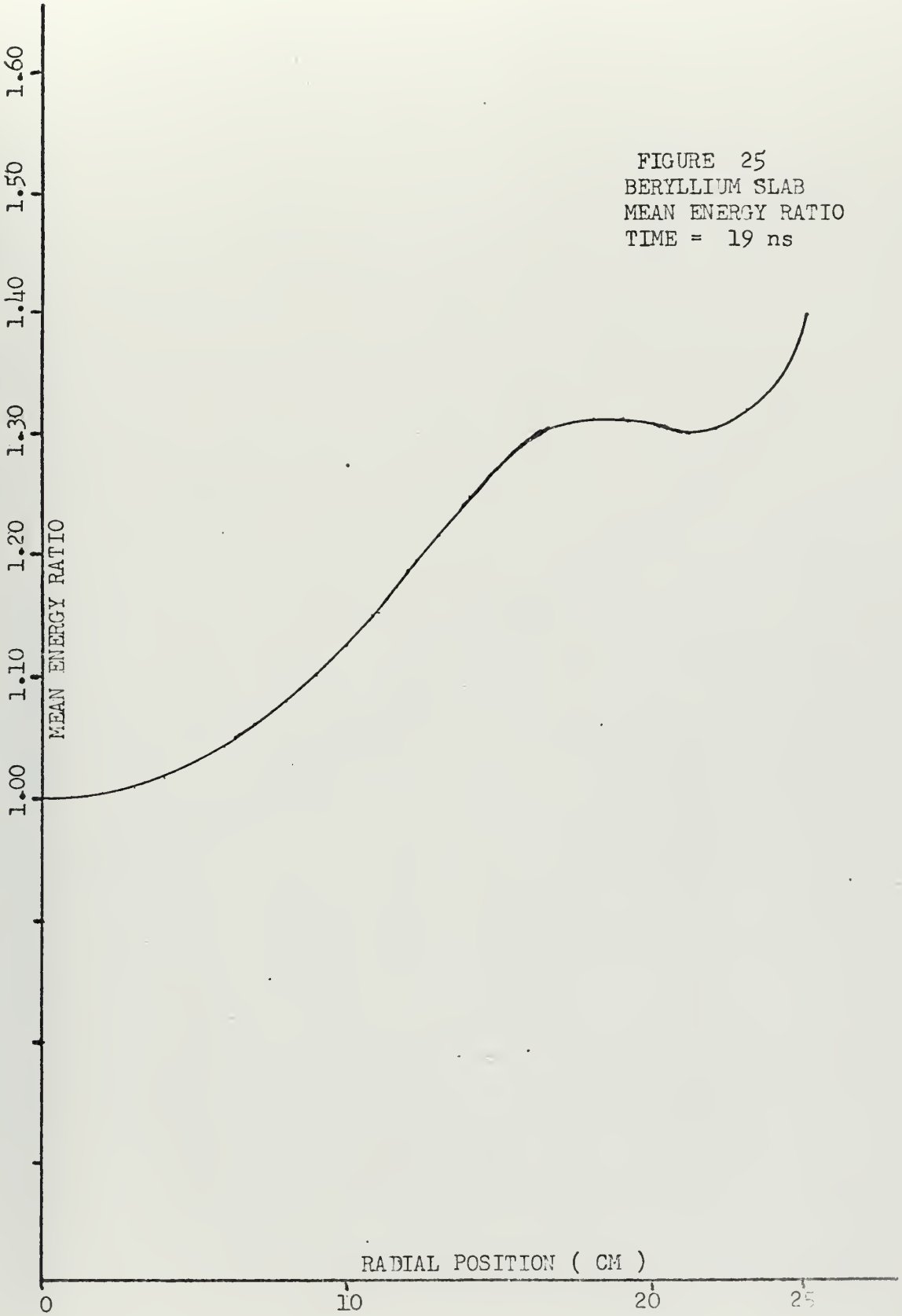


FIGURE 26
BERYLLIUM SLAB
MEAN ENERGY RATIO
TIME = 49 ns

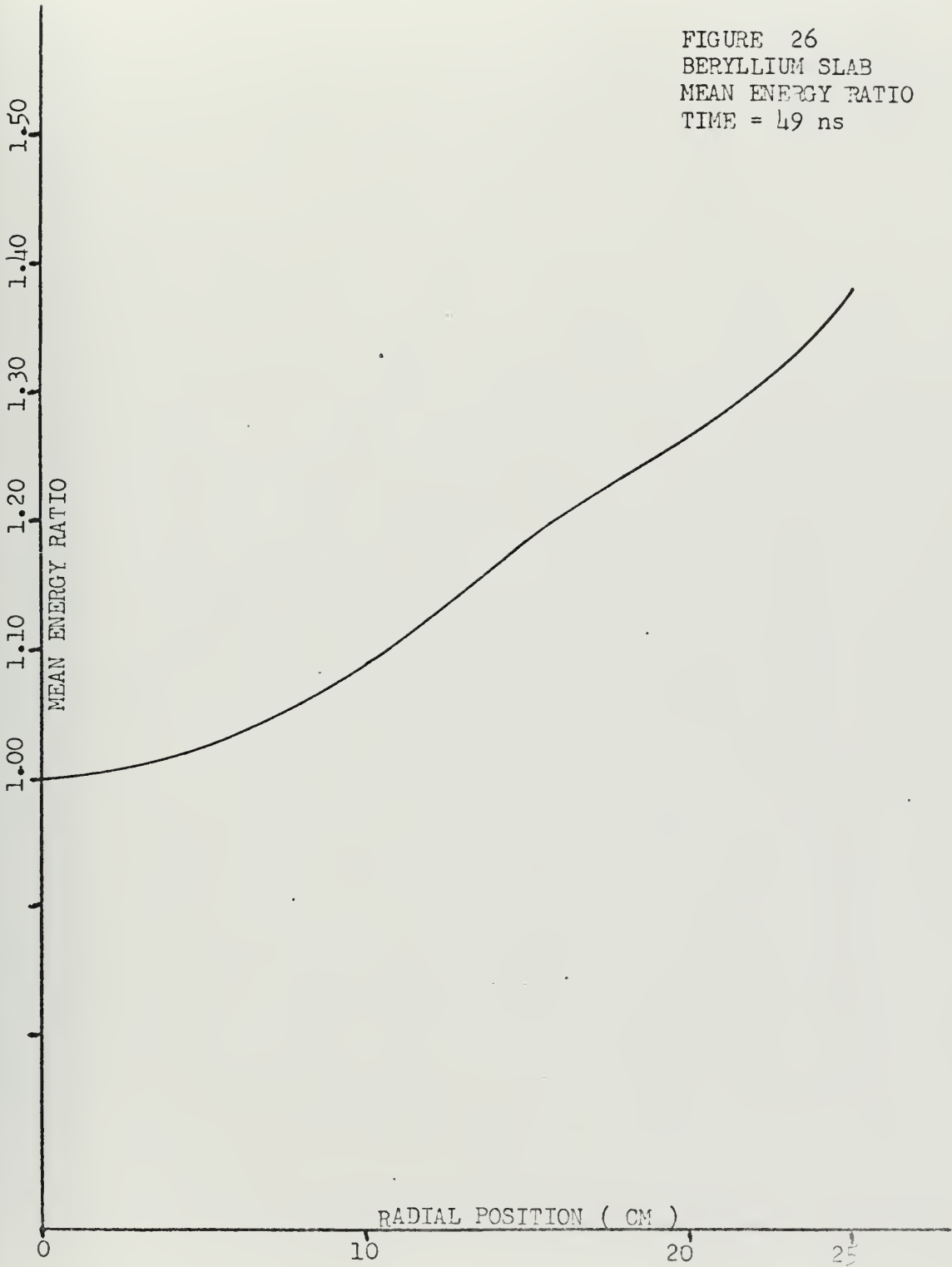


FIGURE 27
BERYLLIUM SLAB
MEAN ENERGY RATIO
TIME = 91 ns

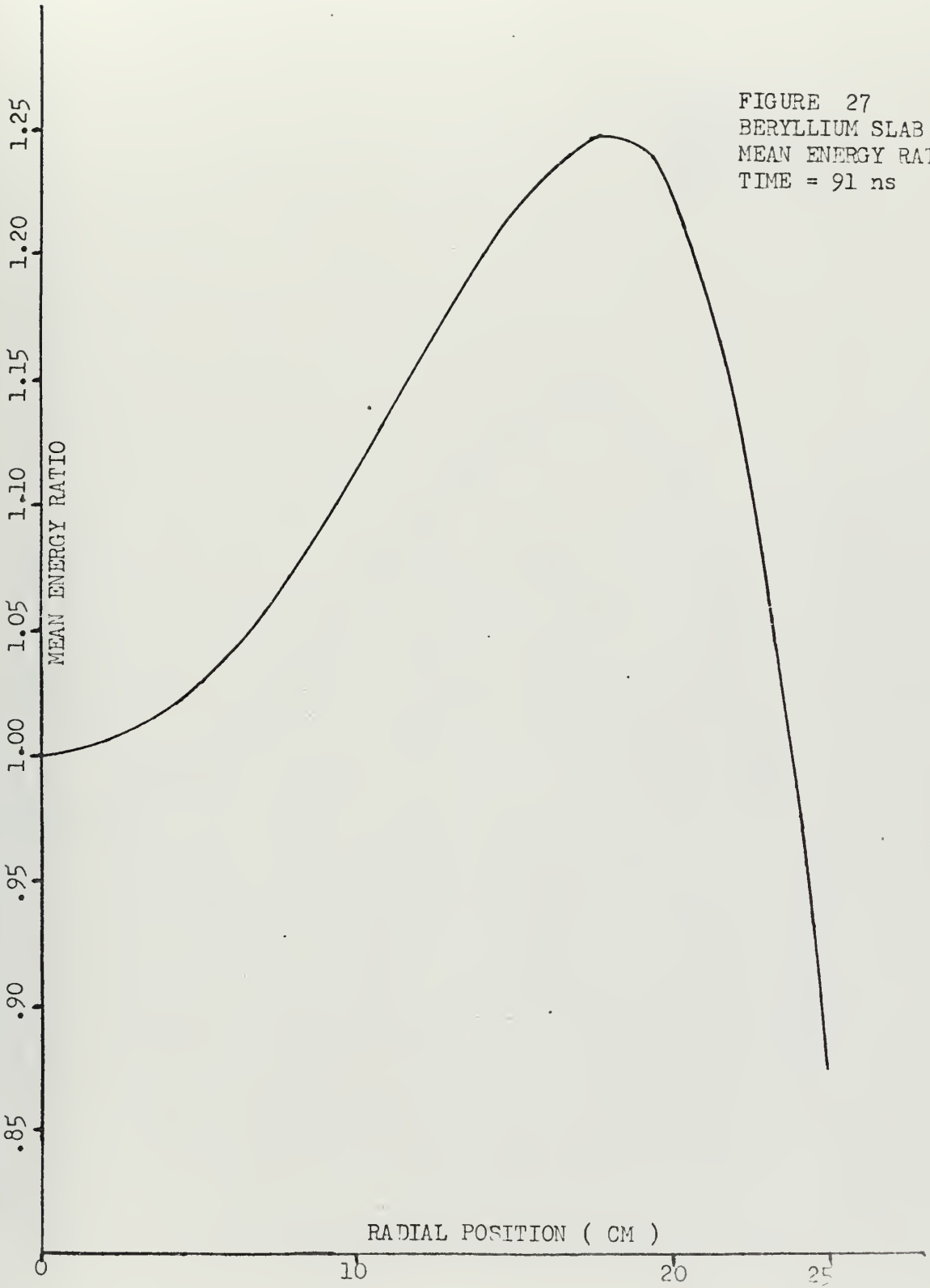


FIGURE 28
BERYLLIUM SLAB
MEAN ENERGY RATIO
TIME = 193 ns

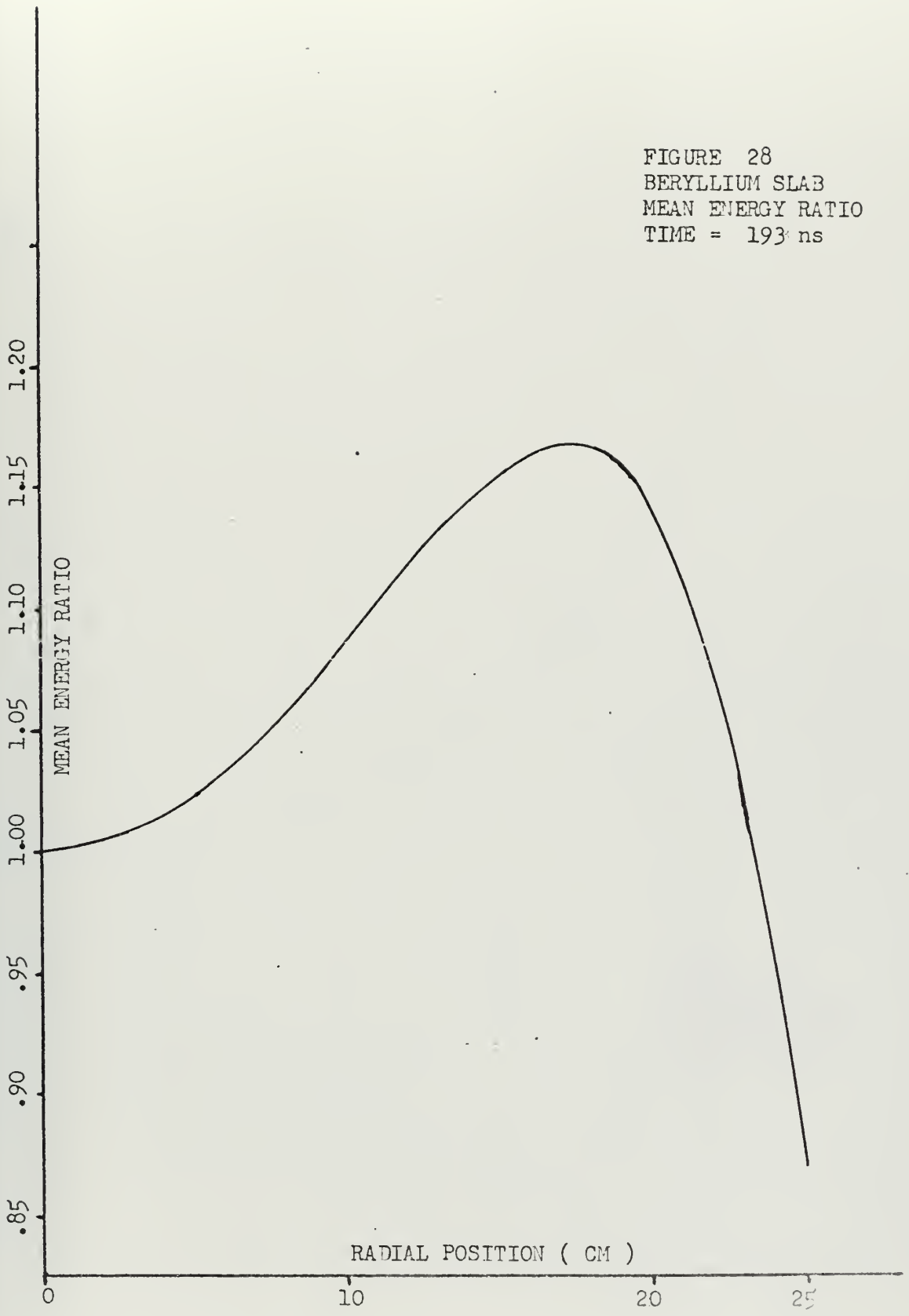




FIGURE 29
BERYLLIUM SLAB
MEAN ENERGY RATIO
TIME = 496 ns

RADIAL POSITION (CM)

FIGURE 30
BERYLLIUM SLAB
MEAN ENERGY RATIO
TIME = 982 ns

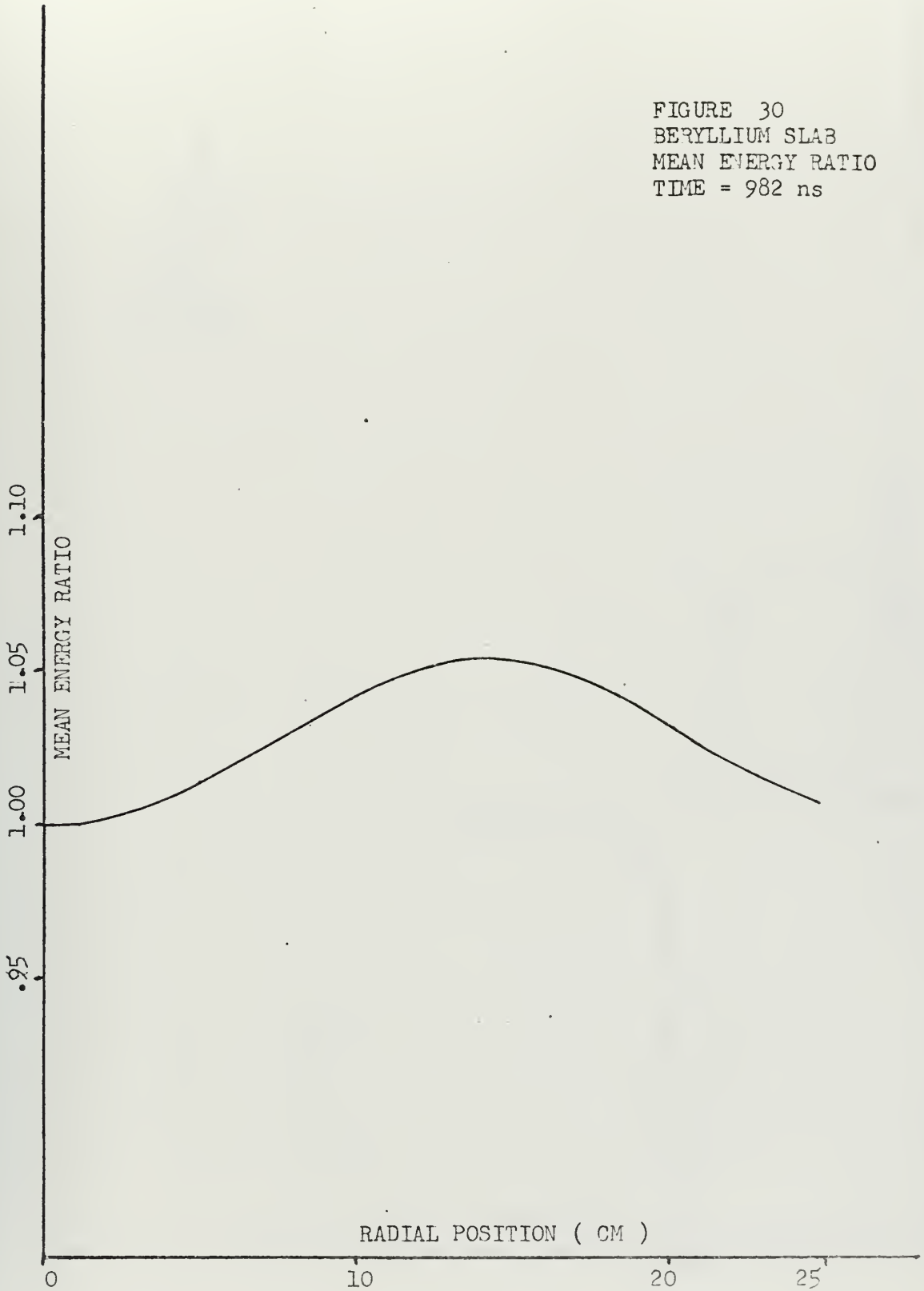


FIGURE 31
ASSEMBLY 6 F
MEAN ENERGY RATIO
TIME = 2 ns

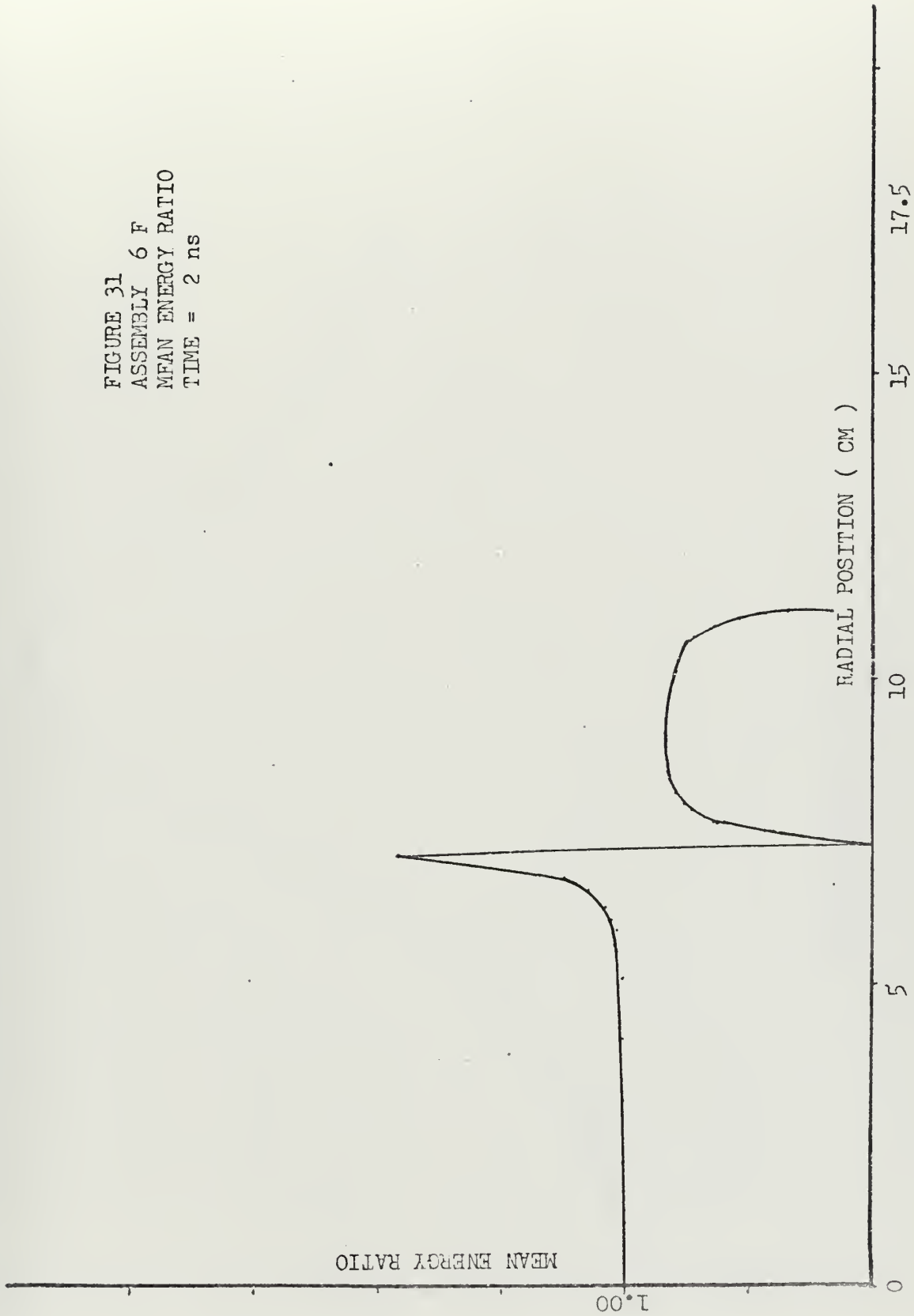


FIGURE 32
ASSEMBLY 6 F
MEAN ENERGY RATIO
TIME = 10 ns

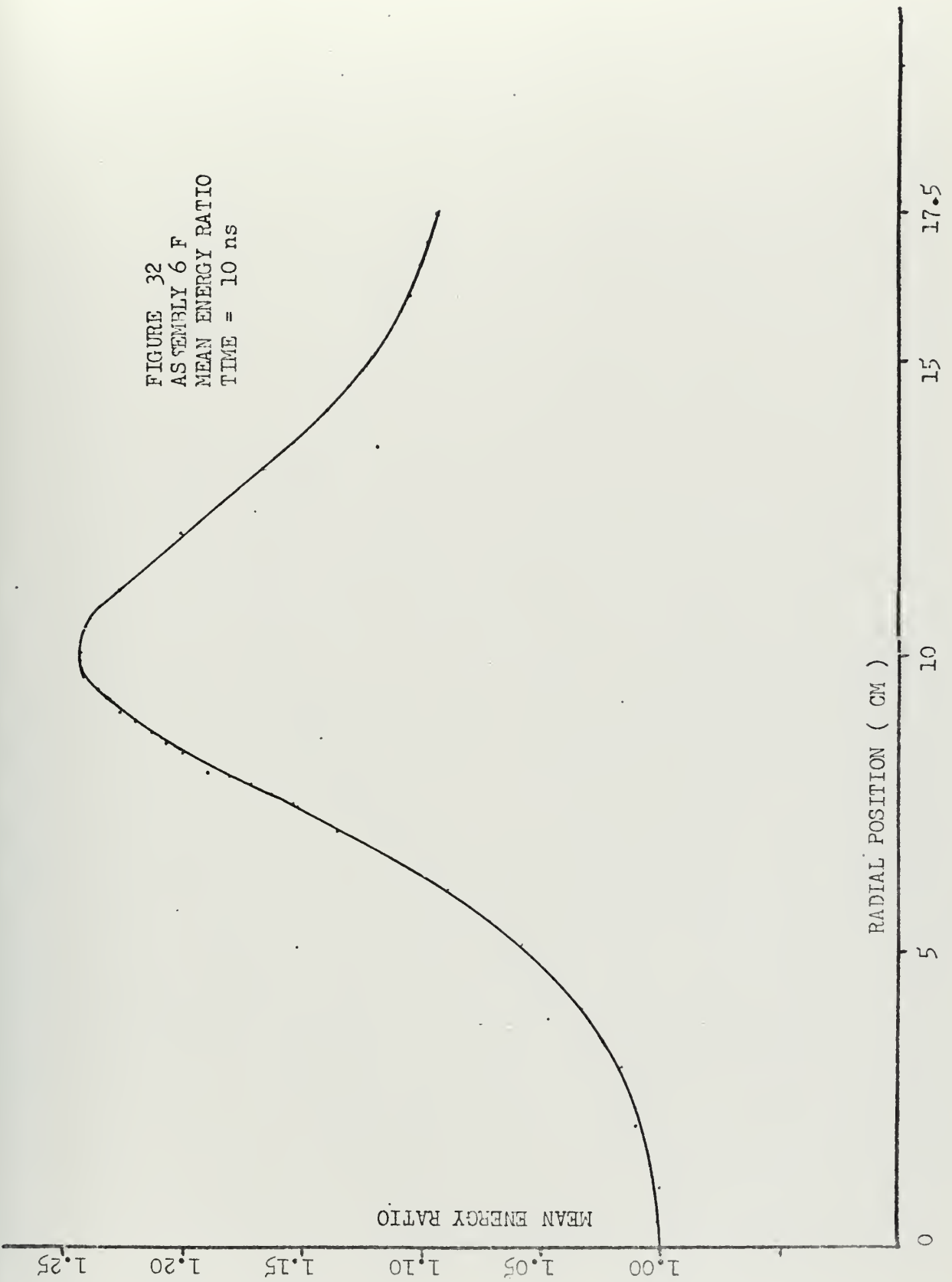


FIGURE 33
ASSEMBLY 6 F
MEAN ENERGY RATIO
TIME = 19 ns

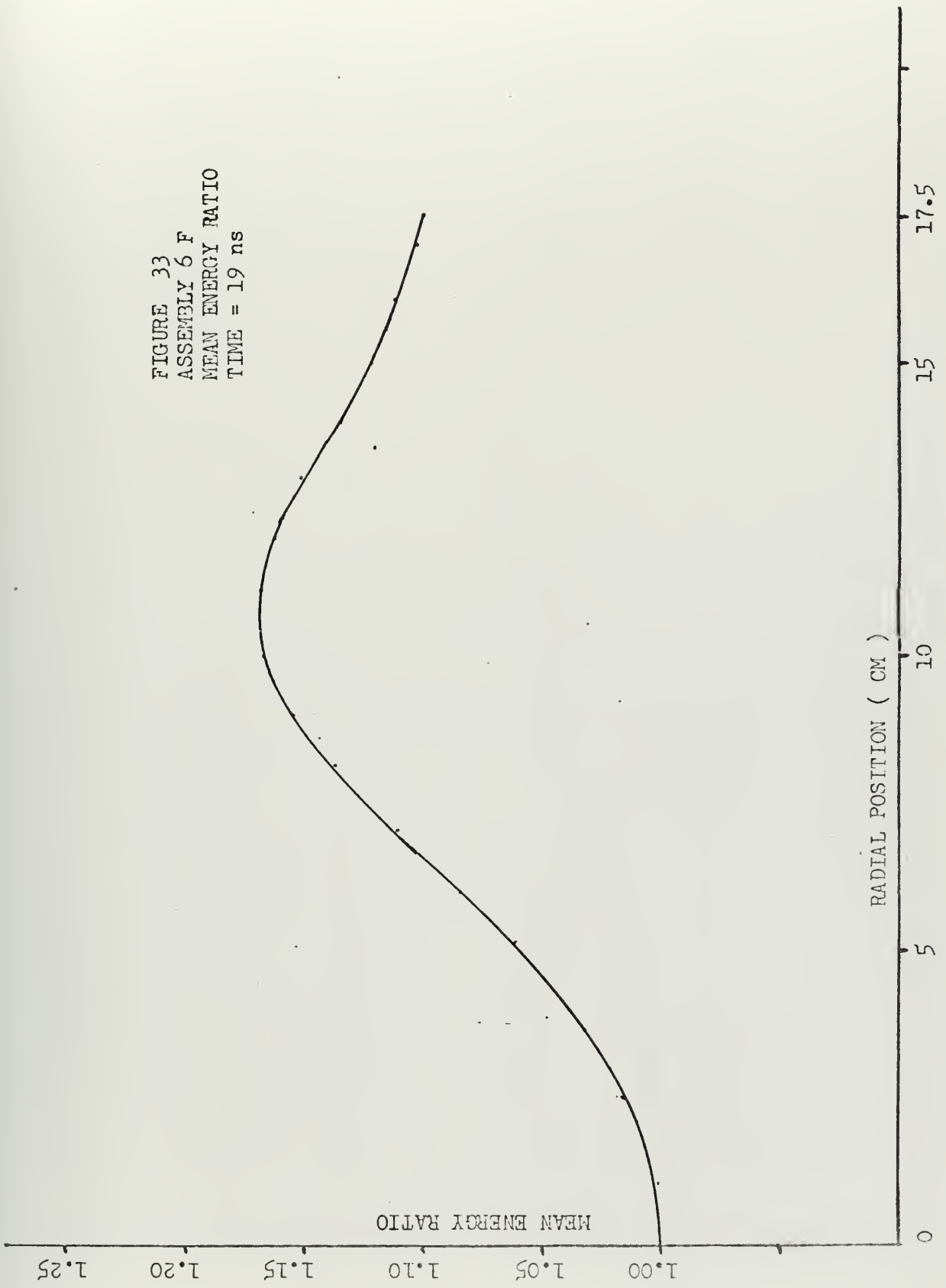


FIGURE 34
ASSEMBLY 6 F
MEAN ENERGY RATIO
TIME = 50 ns

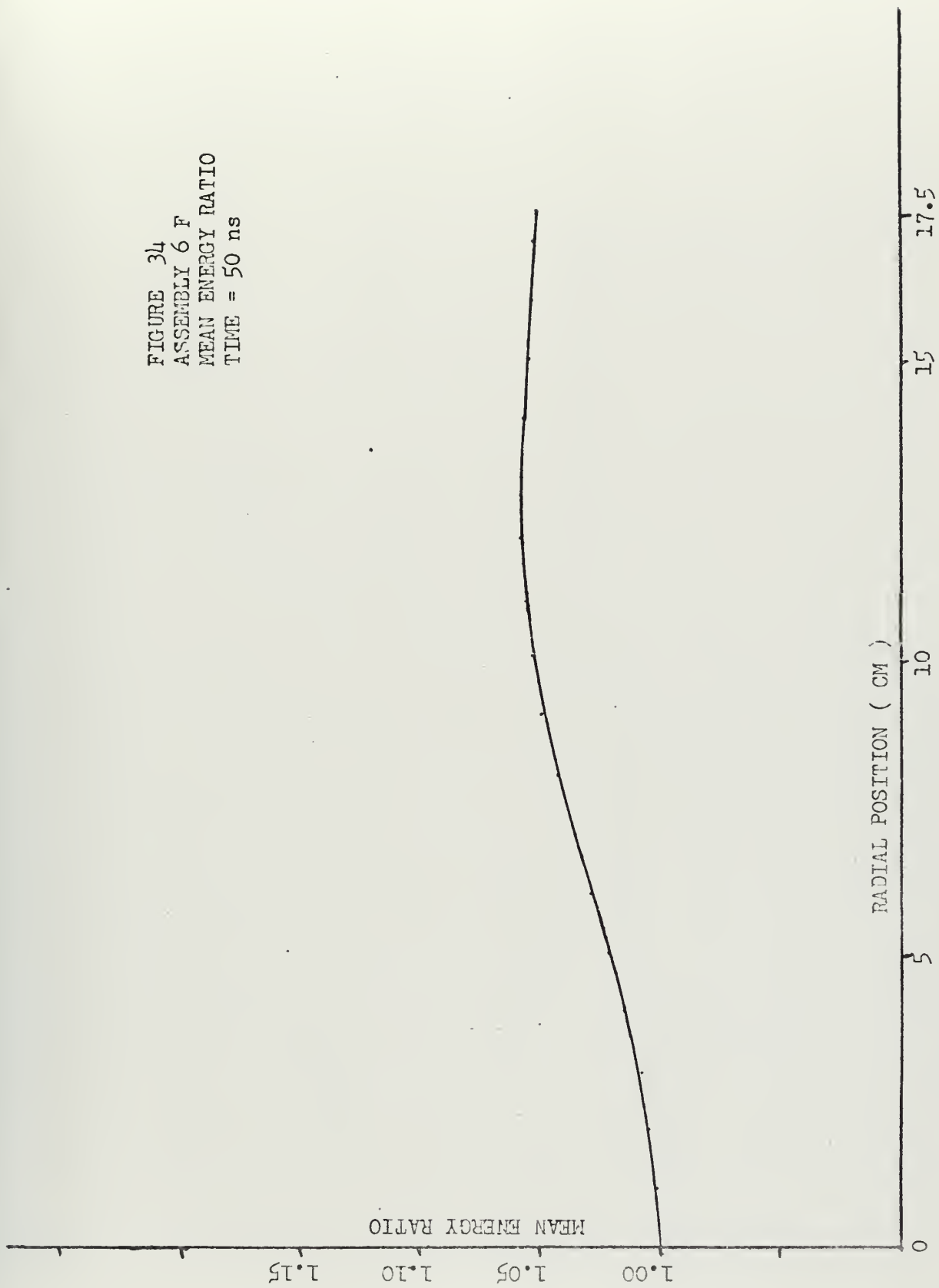


FIGURE 35
ASSEMBLY 6 F
MEAN ENERGY RATIO
TIME = 96 ns

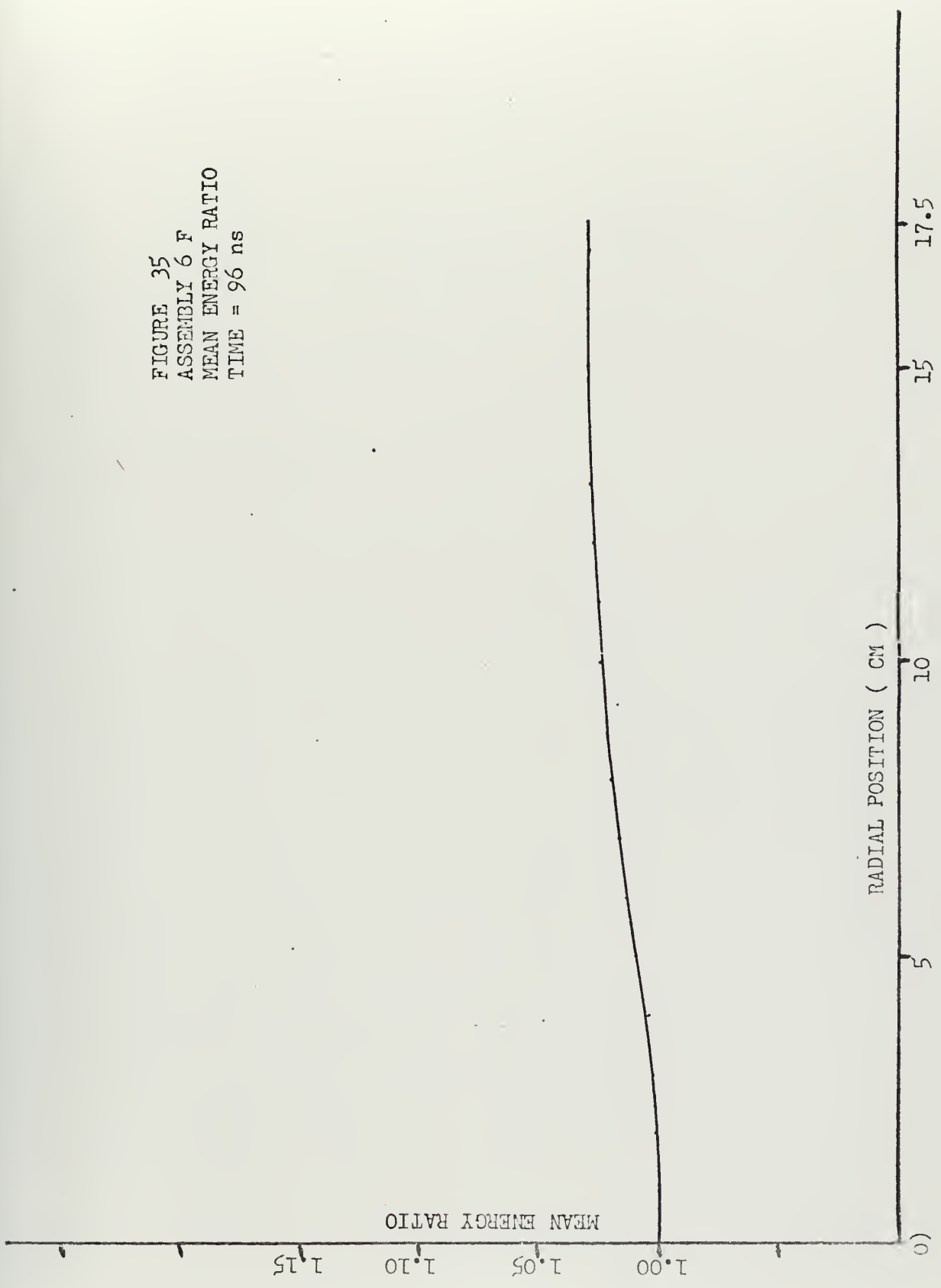


FIGURE 36
ASSEMBLY 6 F
MEAN ENERGY RATIO
TIME = 196 ns



FIGURE 37
BERYLLIUM SLAB
DETECTOR RESPONSE
 $X = 12.5$ CM

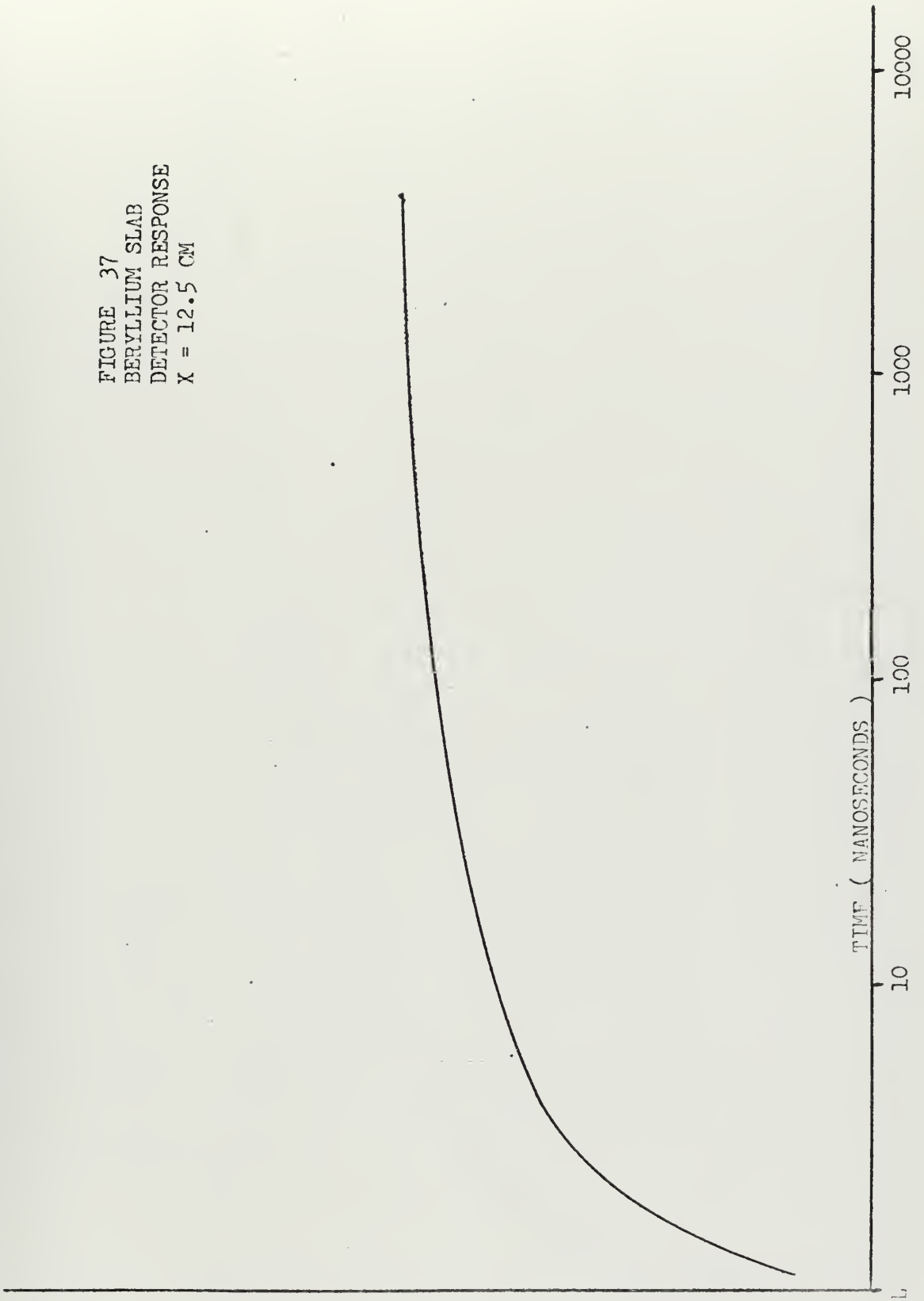
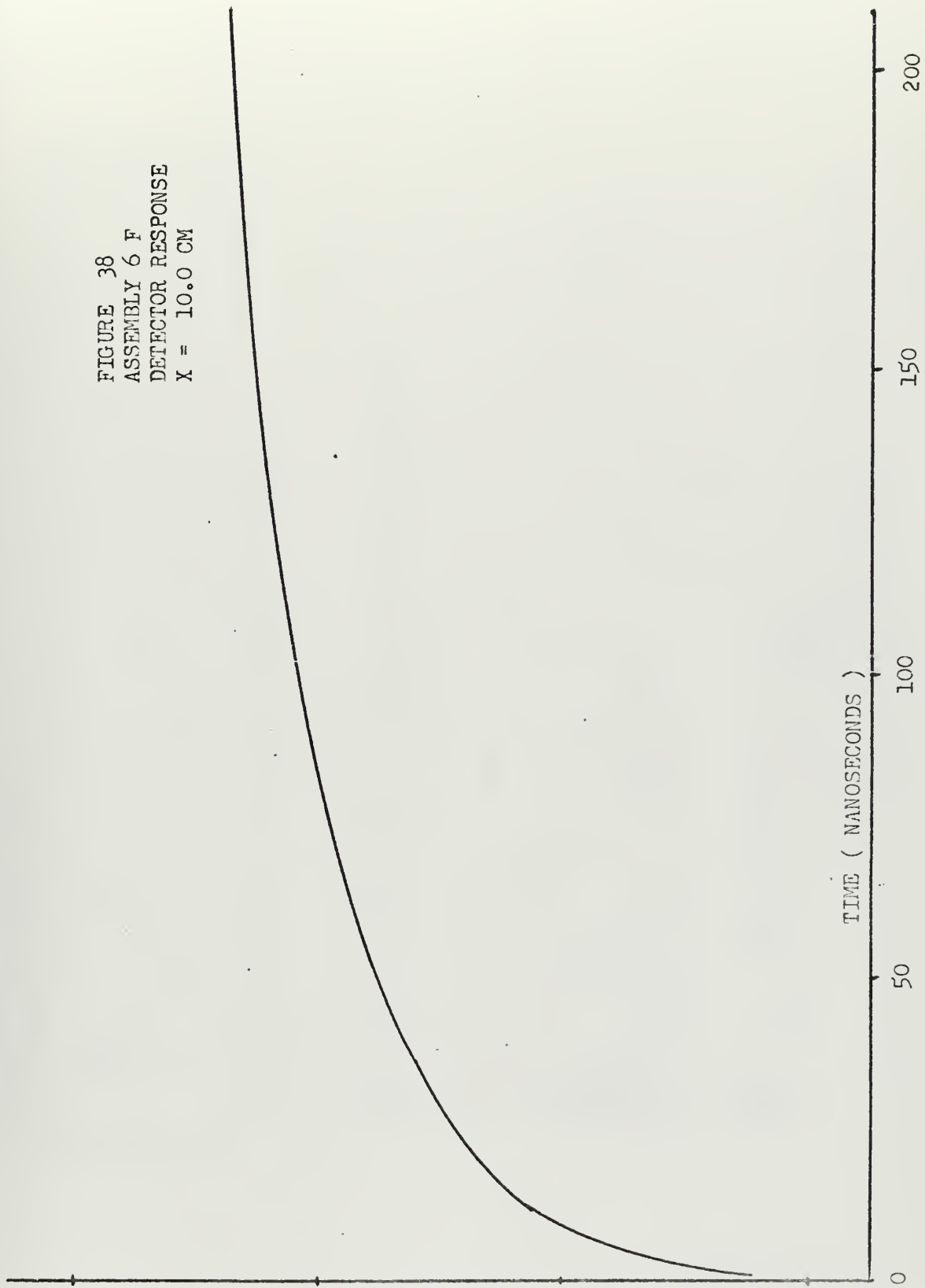


FIGURE 38
ASSEMBLY 6 F
DETECTOR RESPONSE
X = 10.0 CM




```

SOH = HALF WIDTH OF THE SLAB
SA = HALF WIDTH OF SOURCE
S1 = MID POINT OF OFF AXIS SOURCE
S3 = MID POINT LOCATION OF DETECTOR
S4 = HALF WIDTH OF DETECTOR
DIFLEN = 1.0 / SIGMAT - INVERSE TOTAL MACROSCOPIC CROSS SECTION
MODES = 1.0 / NN - TOTAL MODES FOR PROGRAM TO FOLD IN
NSTAT = M - DELTA FUNCTION IN LETHARGY
          2 STATE NUMBER
          3 FISSIION SPECTRUM SOURCE
IBX = NUMBER OF POINTS ON X AXIS TO CALCULATE DATA ABOUT
XBX(NN) INPUTS THE POSITIONS TO DO CALCULATIONS FOR X-AXIS
      TO A MAXIMUM OF 100 DATA POINTS ALONG THE X-AXIS
EMAX - MAXIMUM ENERGY THAT A DETECTOR CAN MEASURE
EMIN - MINIMUM ENERGY THAT THE DETECTOR CAN MEASURE

USE OF THE LOGICAL CONTROL VARIABLES
L( 4) - IN SUBROUTINE FOUT5 SPECTRAL
      - SPECTRA AT POINT OTHER THAN X=0.0
      - SUBTRA AT POINT TOTAL POINTS TO DO.
      USE INPUT VARIABLES NANS(6) FOR TOTAL POINTS TO DO.
      USE POSITION VARIABLES XBX(1) TO XBX(NN)
      EXTERNAL CONTROL - BYPASS THE PROFILE VS X ON CALCOMP
      PLOT MEAN ENERGY - RATIO VS X ON CALCOMP IN FOUT2
      EXTERNAL CONTROL - PERFORM SIMPLE PLOT OF SOURCE
      FUNCTION ON THE PRINTER IN ROUTINE MODEL1
      IF ERROR IN PROBLEM DEFINITION OR CONTROL MODES
      CALL FOUT2-CALC MEAN ENERGY-E(X,T) AND NEW F(X,T)
      INTERNAL CONTROL - BY PASS FROM ANY OTHER ROUTINE IN
      SUBROUTINE FOUT1 ON CALLS FROM
      THE REMAINDER OF THE PROGRAM
      CALL SPECTRAL PLOT ROUTINE FOUT5
      CALL FOUT1 FOR INPUT DATA POINTS
      CALL SUBROUTINE FOUT3 - DETECTOR RESPONSE
      CALL SUBROUTINE FOUT4 FOR SPECTRAL PLOT WITHOUT
      ADJUSTMENT BY THE FOURIER COEFFICIENT
PROGRAM OUTPUT CONTROL VARIABLES
NANS( 1 ) - EXTERNAL CONTROL - LIMIT TO TOTAL GRAPHS TO BE
          ( 2 ) - PLOTTED IN SUBROUTINE FOUT1 - DEFAULT OF 1 GRAPH
          ( 3 ) - INTERNAL CONTROL - LABELING OF FOUT1 GRAPHS ON
          ( 3 ) - CALLS FROM MAIN PROGRAM AND IN ROUTINE FOUT2
          ( 4 ) - EXTERNAL CONTROL - LIMIT TO TOTAL GRAPHS TO BE
          ( 4 ) - PLOTTED BY FOUT2 - DEFAULT TO OPTION = 1
          ( 4 ) - EXTERNAL CONTROL - LIMIT TO ALLOWABLE NUMBER
          ( 5 ) - OF INTERACTIONS IN ROUTINE MODEL1
          ( 5 ) - NOT USED
          ( 6 ) - NUMBER OF POINTS FOR SPECTRAL RESPONSE PLOTTING

```

CC


```

30 CONTINUE
C
C LOGICAL PARAMETER L(18) = 1 RESULTS IN THE POSITION FUNCTION
C F(U,T,X) BEING CALCULATED FOR THE POSITIONS XBX(I)
C AS STATED IN THE INPUT DATA
38 IF (L(18).EQ.0) GO TO 65
C
C CALLING ROUTINE TO CALCULATE THE FUNCTION--F(U,T,X) FOR POSITIONS
C DESIRED AS LISTED IN THE INPUT DATA FOR THAT POINT
C IF ANY OUTPUT IS DESIRED FROM ROUTINE FOUT1 VIA INITIAL
C CALLING SEQUENCE FROM THE MAIN PROGRAM -- DEFINE L( 16 ) = 0
C IN THE PROBLEM INPUT STRUCTURE
40 IF (IBX.EQ.0) XBX(1) = 0.
C IF (IBX.EQ.0) IBX=1
C IF (IS.EQ.0) IS=1
C PROTECTION STEP TO ALLOW FOUT1 TO PRINT OUT THE VALUES
C OF THE FUNCTION -- F(U,T,X) ON THE CALLS FROM MAIN PROGRAM
C L(16) = 0
C DO 43 I1X = 1, IBX
C IF (NOUT.EQ.3) NANS(2) = I1X
C WRITE (6,63)
C CALL FOUT1(XBX(I1X))
C CONTINUE
43 CONTINUE
65 IF (L(14).EQ.1) CALL FOUT2
C
C 1 SUBROUTINE FOUT2 CALC MEAN ENERGY AND PLOT SUM
C OVER ALL LETHARGY STATES OF F(U,T,X) AND PLOT
C BOTH FUNCTIONS ON THE SAME PLOT FOR EACH TIME STEP
C
C IF(L(19).EQ.0) GO TO 48
C PARAMETER L(19) = 1 AS DEFINED IN THE PROBLEM DEFINITION
C TO CALCULATE THE DETECTOR RESPONSE IN SUBROUTINE FOUT3
C IF L(19) = 0 --- WE WILL BY-PASS THIS ROUTINE
C
C WRITE (6,64)
C CALL FOUT3
48 CONTINUE
49 CONTINUE
GO TO 60
C
C THIS WILL ALLOW A BYPASS AND STOP IF ERROR
50 WRITE (6,51)
GO TO 60
52 WRITE (6,53)
60 CONTINUE
WRITE (6,61) (NPROB(IA),IA=1,54)
51 FORMAT (//,10X,'INPUT ERROR IN DATA FILE - F(U,T) FROM MOD5',/,

```



```

110X, 'PROGRAM TERMINATED')
53 FORMAT (///, 'ERROR IN PROBLEM STRUCTURING', /,
110X, 'RUN TERMINATED')
61 FORMAT (///, 3(10X, 18A4, /), 10X, 4(3X, '**END**', 3X))
63 FORMAT ('1', 10X, 'EXECUTION BEGINS WITH CALL TO FOUT1')
64 FORMAT ('1', 10X, 'EXECUTION CONTINUES WITH CALL TO FOUT3')
STOP
END

```

```

SUBROUTINE INCON1
REVISED 2 JUNE 1971

```

```

1 THIS ROUTINE DOES THE FOLLOWING
2 SETS ALL ARRAY LOCATIONS TO ZERO
DEFINES ALL PRINCIPAL DEFAULT OPTIONS, I.E.

```

```

SOH = 100.0 CM
SA = 1.00 CM
S1 = 0.00 CM
S2 S3 S4 = 0.0
DIFLEN = 25.0
IS = 1
IDATA = 1
MODES = 1
NSTAT = 1 2
IBX=0
SZERO = 1.0

```

```

3 READS ALL PROBLEM DEFINITION STATEMENTS FROM PUNCH CARDS
AND ECHO PRINTS ALL INPUT DATA
4 INFORMS THE USER WHAT ACTIONS THE PROGRAM WILL EXECUTE
REAL, LTAB(20)/, T1, T2, T3, T4, T5, T6, T7, T8, T9,
1, T10, T11, T12, T13, T14, T15, T16, T17, T18, T19, T20 /
DIMENSION NDATAC(18)
NAMELIST/DATAIN/ IS, SZERO, SOH, SA, S1, S2, S3, S4, IDATA, NOUT, IBX, XBX,
1 DIFLEN, MODES, NSTAT, EMAX, EMIN
OPTION FOR THE MAX AND MIN ENERGYS THAT THE DETECTOR RESPONSE
ROUTINE WILL HAVE TO CONSIDER
COMMON/COM1/ XBX(100), U(150), E(150), TIM(20), DUMMY(150), TXBX(100)
COMMON/COM4/ SOH, SA, S1, S2, S3, S4, SZERO, BUCKLE, EMAX, EMIN, DIFLEN, X
COMMON/COM7/ SO1, SO2, SU3, SU4, AK11, AK12, AK13, AK14, AN3A, AN3B,
1 AN3C, AN4, ANA4, ANB4, AN2X, AK15, AK5X
COMMON/COM2/ NT(20), NPROB(54), NANS(6), L(20), LTAB
COMMON/COM5/ MODES, IS, IDATA, N, NVIR, NI, NF, NCALL, NOUT, IBX,
1 NSTAT, MODE
COMMON/COM3/ FXT(100, 20), FUTX(150, 20), BUCK(6), FUT(6, 71, 20)
COMMON/COM6/ NTUP(6, 20), NM(6, 20), IKNT(6)

```



```

WRITE (6,28)
READ (5,29) NCARD
DO 27 I=1,NCARD
READ (5,30) (NDATAC(IA),IA=1,18)
WRITE (6,31) (NDATAC(IA),IA=1,18)
CONTINUE
27 END OF DEBUG DATA PRINT OUT
C START BY READING THE PROBLEM TITLE
C READ (5,9) (NPROB(IA),IA=1,54)
C READ IN THE NAMELIST DATA
C READ (5,DATAIN)
C INPUT OF PROGRAM LOGICAL CONTROLS
C READ (5,11) (NANS(I),I=1,6)
C DUPLICATE DATA PRINT OUT OF THE PROBLEM DEFINITION CARDS
C READ (5,12) (L(I),I=1,20)
C OUTPUT OF PROBLEM STATE STRUCTURE
C START PROBLEM LISTING ON A NEW PAGE
C WRITE (6,28) (NPROB(IA),IA=1,54)
C WRITE (6,61) IS, IDATA, NOUT, IBX, SZERO, DIFLEN, SOH, SA, S1, S2, S3, S4
C WRITE (6,63) MODES, NSTAT
C IF (IBX.EQ.0) GO TO 68
C
C PROTECTION STEPS TO TERMINATE PROGRAM EARLY IN EXECUTION IF
C SOME ERRORS EXIST IN THE BASIC PROBLEM STATEMENT
C COMPARE ALL PERTINENT DATA DEPENDING ON THE SUBROUTINES TO
C BE DONE TO INSURE THAT ALL NECESSARY INPUT HAS BEEN PROVIDED
C LOGICAL L(10) = 0 INITIALLY-- IF INPUT DATA IS IN ERROR,
L(10) SET = TO 1 PROBLEM WILL BE TERMINATED IN MAIN PROGRAM
ON RETURN FROM THIS SUBROUTINE.
L(10)=0
IF ((IBX.LT.0).OR.(IBX.GT.100)) L(10)=1
IF (IBX.GT.100) IBX=100
IF (IBX.GT.0) WRITE (6,62) (XBX(J),J=1,IBX)
GO TO 69
68 WRITE (6,67) XBX(1)
69 WRITE (6,70) (NANS(IA),IA=1,6)
WRITE (6,72) (L(I),I=1,20)
IBX=1
C COMPLETE TEST SERIES TO EXAMINE MAIN INPUT CONTRAL DATA
IF ((NOUT.GT.3).OR.(NOUT.LE.0)) L(10)=1
IF ((IS.GT.5).OR.(IS.LE.0)) L(10)=1
IF ((SOH.LE.SA).OR.(SOH.LE.S1).OR.(SOH.LE.S3)) L(10)=1
IF ((NSTAT.GE.4).OR.(NSTAT.LE.1)) L(10)=1
IF ((MODES.LT.1).OR.(MODES.GT.6)) L(10)=1
TEST FOR TOTAL MODES MUST BE REVISED IF PROGRAM IS MODIFIED
TO PROCESS MORE THAN 6 MODES OF DATA

```



```

C THIS TEST QUESTION MUST BE CHANGED TO REFLECT THE STORAGE
C ARRAY CAPABILITIES OF COMMON BLOCKS COM3 & COM6
C
C PROGRAM WILL WRITE THE USER A SERIES OF MESSAGES TO DESCRIBE
C WHAT ACTIONS THE PROGRAM WILL EXECUTE IN THE PERFORMANCE
C OF THIS SPECIFIC JOB
200 WRITE (6,220)
201 GO TO (201,202,203,204,205),IS
201 WRITE (6,221) SOH,SA
202 GO TO (204,204,204,204,204),SI,SA,SOH
202 WRITE (6,222) SI,SA,SOH
203 GO TO (204,204,204,204,204),SOH,DIFLEN
203 WRITE (6,223) SOH,DIFLEN
204 GO TO (204,204,204,204,204),SI,SOH,DIFLEN
204 WRITE (6,223) SI,SOH,DIFLEN
2045 GO TO (204,204,204,204,204)
2041 WRITE (6,2250)
2041 CONTINUE
205 GO TO (205,206,207),IDATA
205 WRITE (6,224)
206 GO TO (208,208,208,208,208)
206 WRITE (6,225)
207 GO TO (208,208,208,208,208)
207 WRITE (6,226)
208 GO TO (209,210,211),NOUT
209 WRITE (6,227)
210 GO TO (212,212,212,212,212)
210 WRITE (6,228)
211 GO TO (212,212,212,212,212)
211 WRITE (6,229)
212 CONTINUE
213 WRITE (6,230) MODES
PROGRAM EXECUTION LISTING TO ADVISE THE USER THE ROUTINES
THAT WILL BE EXECUTED AND THE PRINCIPAL ITEMS OF DATA THAT
WILL BE USED BY EACH OF THEM
THIS WILL ASSIST THE USER IN ANY CASES WHERE SOME UNEXPECTED
EXECUTION ERROR MAY OCCUR
STATEMENTS ARE LISTED IN THE ORDER IN WHICH THE PROGRAM WILL
EXECUTE THEM
FOUT4 -- SPECTRAL DATA PLOT
IF (L(20).EQ.0) GO TO 301
WRITE (6,320) MODES
FOUT5 -- FOURIER ADJUSTED SPECTRAL DATA PLOT
IF (L(17).EQ.0) GO TO 302
WRITE (6,321) MODES.
NPXT5=NANS(6)
IF (L(4).EQ.1) WRITE(6,321) NANS(6), (XB(IJ), IJ=1,NPXT5)
FOUT1 - CALCULATE F(U,T,X) = F(U,T)*PHI(X) FOR INPUT DATA

```



```

302 IF (L(18).EQ.0) GO TO 303
    WRITE (6,322) IBX,NOUT,NANS(1),(XBX(IJ),IJ=1,IBX)
    FOUT2 -- MEAN ENERGY AND FLUX PROFILE
    IF (L(14).EQ.0) GO TO 304
    WRITE (6,323) NANS(3)
    IF (L(5).EQ.1) WRITE (6,3231)
    IF (L(7).EQ.1) WRITE (6,3232)
    IF (L(8).EQ.1) WRITE (6,3233)
    FOUT3 -- DETECTOR RESPONSE
    IF (L(19).EQ.0) GO TO 305
    WRITE (6,324) S3,S4,EMAX,EMIN,NOUT
305 CONTINUE
    FORMAT (18A4)
    9 FORMAT (6I2)
    11 FORMAT (20I2)
    28 FORMAT (11,/,/,20X,'INPUT DATA FILE')
    29 FORMAT (6X,14)
    30 FORMAT (18A4)
    31 FORMAT (10X,18A4)
    60 FORMAT (10X,10X,IS = ,6X,I3,/,10X,ISDATA = ,3X,I3,/,10X,ISOUT = ,
    14X,I3,/,10X,ISPOINTS = ,14,/,10X,ISOURCE STRENGTH = ,2X,F6.2,
    2/,10X,DIFFUSION LENGTH = ,F5.2,/,10X,SOH = ,F6.2,3X,SA = ,
    43X,S4 = ,2X,F5.2)
    62 FORMAT (10(5X,10(F8.3,2X),/))
    63 FORMAT (/,10X,MODES = ,2X,I5,/,10X,NSTAT = ,3X,I5)
    67 FORMAT (10X,F8.5)
    70 FORMAT (/,10X,LOGICAL CONTROL VARIABLES - NANS(I),/,
    11X,6(5X,I3))
    72 FORMAT (/,10X,'PROGRAM LOGICAL CONTROL VARIABLES',/,
    16X,4(3X,5I3,3X))
    220 FORMAT (11,10X,3(2X,'**',2X),/,10X,'MIL SIX FOLDING',
    11X,3(2X,'**',2X),/,10X,'PROGRAM WILL',
    21X,'EXECUTE THE FOLLOWING')
    221 FORMAT (/,10X,DELTA FUNCTION SOURCE ON AXIS AT X = 0.0,/,
    110,'SLAB HALF THICKNESS = ,1X,F8.2,1X,'SOURCE HALF THICKNESS',
    21X,= ,1X,F6.2)
    222 FORMAT (/,10X,DELTA FUNCTION SOURCE OFF AXIS -CENTERED AT X = ,
    12X,F8.2,1X,'SOURCE HALF THICKNESS = ,1X,F8.2,/,
    210X,'SLAB HALF THICKNESS = ,1X,F8.2)
    223 FORMAT (/,10X,'FIRST COLLISION SOURCE CENTERED ABOUT X = 0.0',
    1/,110,'SLAB HALF THICKNESS = ,2X,F8.2,/,
    210X,'DIFFUSION LENGTH = ,2X,F8.2)
    2231 FORMAT (/,10X,'FIRST COLLISION SOURCE CENTERED ABOUT X = ,2X,
    1F8.2,2X,/,10X,'SLAB HALF THICKNESS = ,2X,F8.2,/,
    210X,'DIFFUSION LENGTH = ,2X,F8.2)
    224 FCRMAT (/,10X,'INPUT DATA IS PUNCHED CARDS:')

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225 FORMAT (/ , IGX, ' INPUT DATA FROM MAGNETIC TAPE' )
226 FORMAT (/ , IGX, ' INPUT DATA FROM DISK' )
227 FORMAT (/ , IGX, ' OUTPUT WILL BE LISTED ON THE PRINTER' )
228 FORMAT (/ , IGX, ' OUTPUT PLOTS WILL BE PLOTTED ON THE PRINTER' )
229 FORMAT (/ , IGX, ' GRAPHICAL OUTPUT WILL BE PLOTTED USING THE',
1 IX, 'CAL COMP PLOTTER' )
230 FORMAT (/ , IGX, ' MODES OF DATA TO PROCESS =', 2X, I5)
320 FORMAT (/ , IGX, 'L(2U)=1, T25, 'SUBROUTINE-FOUT4 WILL BE',
1 IX, 'EXECUTED', / , VS U, / , T35, I5, 3X, 'SPECTRAL PLOT OF DATA',
22 X, '---> F(U, T) ' / , T25, 'SUBROUTINE-FOUT5 WILL BE EXECUTED', / ,
321 / , T25, 'SPECTRAL PLOT OF MOD-5 OUTPUT ADJUSTED BY THE FOURIER', / ,
25 X, 'FOR EACH OF', 1X, I3, 1X, 'MODES', / ,
4 T25, 'AND A SUM OF ALL MODES WILL BE PLOTTED' )
322 FORMAT (/ , IGX, 'L(18)=1, T25, 'SUBROUTINE-FOUT1 WILL BE EXECUTED',
2 / , T25, 'CALCULATION OF THE FUNCTION-- F(U, T)*PHI(X)',
31 X, I3, / , T25, 'LIMIT OF', DATA I3, 2X, 'PLOTS FOR GRAPHICAL OUTPUT', / ,
4 T25, 'DATA POINT POSITIONS ARE', / , 10(IGX, I3(F6.2, 4X), / ))
323 FORMAT (/ , IGX, 'L(14)=1, T25, 'SUBROUTINE-FOUT2 WILL BE EXECUTED',
2 / , T25, 'CALCULATION OF THE MEAN ENERGY AND FLUX-DENSITY PROFILE',
3 T25, 'EMEAN(X, T) FOR EACH TIME STEP', / ,
4 IX, 'DONE' )
324 FORMAT (/ , IGX, 'L(19)=1, T25, 'SUBROUTINE-FOUT3 WILL BE EXECUTED',
2 F6.2, 2X, 'WITH', / , T25, 'A DETECTOR RESPONSE AT POSITION', 1X,
3 T25, 'DETECTOR WILL RESPOND TO NEUTRON ENERGY-EMAX=', 1X, 1PE10.3,
42 X, 'EMIN =', 1X, 1PE10.3, ' EV', / , T25, 'OUTPUT VIA DEVICE =', 1X,
5 I5)
2250 FORMAT (/ , IGX, 'RAMP FUNCTION-EXTERIOR SOURCE' )
3211 FORMAT (/ , IGX, 'L( 4)=1, T25, 'A SPECTRAL PLOT WILL BE DONE FOR',
12 X, I5, 2X, 'POINTS', / , T20, 'POINTS ARE', / , 10X, 5(2X, F8.5, 2X) )
3231 FORMAT (/ , IGX, 'L( 5)=1, T25, 'NO PRINTED LISTING OF MEAN ENERGY',
11 X, 'OR FLUX PROFILE WILL BE DONE' )
3232 FORMAT (/ , IGX, 'L( 7)=1, T25, 'OVER RIDE OF OUTPUT SEQUENCE',
11 X, 'ONLY FLUX PROFILE WILL BE PLOTTED' )
3233 FORMAT (/ , IGX, 'L( 8)=1, T25, 'OVER RIDE OF OUTPUT SEQUENCE',
11 X, 'ONLY MEAN ENERGY RATIO WILL BE PLOTTED' )
      RETURN
      END

```


SUBROUTINE ESTIM1
REVISED 2 JUNE 1971

INITIAL ESTIMATING PROGRAM TO PROVIDE GUIDANCE ON TOTAL
ESTIMATED EXECUTION TIME WHILE ERROR TESTING THE VERSION
OF MIL SIX INITIAL PROGRAM CHECKOUT TO INFORM THE USER OF THE
ANTICIPATED PROGRAM EXECUTION PARAMETERS
EXECUTE THIS ROUTINE IMMEDIATELY AFTER READIO OR READ 11
TO BEST ESTIMATE TOTAL EXECUTION TIME AFTER WE HAVE BETTER
KNOWLEDGE OF THE TOTAL PROGRAM OPERATING PARAMETERS

COMMON/COM1/ XBX(100), U(150), E(150), TIM(20), DUMMY(150), TXBX(100)
COMMON/COM4/ SOH, SA, S1, S2, S3, S4, SZERO, BUCKLE, EMAX, EMIN, DIFLEN, X
COMMON/COM7/ S01, S02, S03, S04, AK11, AK12, AK13, AK14, AN3A, AN3B,
1 AN3C, AN4, ANA4, ANB4, AN2X, AK15, AK5X
COMMON/COM2/ NT(20), NPROB(54), NAN\$ (6), L(20), LTAB(20)
COMMON/COM5/ MODES, IS, IDATA, N, NVIR, NI, NF, NCALL, NOUT, IBX,
1 INSTAT, MODE
COMMON/COM3/ FXT(100,20), FUTX(150,20), BUCK(6), FUT(6,71,20)
COMMON/COM6/ NTOP(6,20), NM(6,20), IKNT(6)
REAL ST(15)/15*0.0/
WRITE (6,76)
WRITE (6,77)

DEFINE TOTAL TIME AS A ZERO VALUE PRIOR TO CALCULATION

SUMT = 0.0

TIME FOR MAIN AND INCON1

ST(1) = 0.75

EXECUTION TIME FROM READ PROGRAMS

IF ((IDATA.EQ.1) .AND. (NOUT.EQ.1)) ST(2) = MODES*1.25 + 0.50

IF ((IDATA.EQ.1) .AND. (NOUT.EQ.3)) ST(2) = MODES*1.0 + 0.50

IF ((IDATA.EQ.2) .AND. (NOUT.EQ.1)) ST(3) = MODES*1.25 + 0.50

IF ((IDATA.EQ.2) .AND. (NOUT.EQ.3)) ST(3) = MODES*1.0 + 0.40

FOUT1 TIME

IF ((L(18).EQ.1) .AND. (NOUT.EQ.1)) ST(4) = MODES*IBX*12.50

IF ((L(18).EQ.1) .AND. (NOUT.EQ.3)) ST(4) = MODES*IBX*2.50

FOUT2 TIME ESTIMATE

IF (L(14).EQ.1) ST(5) = 10.0*MODES

FOUT3 TIME ESTIMATE

IF ((L(19).EQ.1) .AND. (NOUT.EQ.1)) ST(6) = MODES*20.0

IF ((L(19).EQ.1) .AND. (NOUT.EQ.3)) ST(6) = MODES*5.00

ESTIMATED FOUT4 TIME FOR EXECUTION

IF (L(20).EQ.1) ST(7) = MODES*3.0

FOUT5 TIME ESTIMATE

IF (L(17).EQ.1) ST(8) = MODES*4.0

ESTIMATED TIME - ESTIM1

ST(9) = L*25

ESTIMATED TIME - MODEL1

ST(10) = 0.25*MODES


```

AKN3=AK13*I MODE
AN3=AN3C*(AKN3*M66M*AN3B-AN3A)/(1.0+(AKN3*DIFLEN)**2)
PHI1=AN3*COS(AKN3*X)
GO TO 10
4 CONTINUE
C
IS=4
M44M=2*I MODE-1
M55M=(-1)**(I MODE+1)
AKN4=AK14*M44M
AK5=AKN4*S1
AK6=AKN4*DIFLEN
AK7=1.0+(AKN4*DIFLEN)**2
SIX5=SIN(AK5)
CIS6=COS(AK5)
AN4=((ANA4/AK7)*((CIS6-AK6*SIX5-(M55M*AK6*ANB4))
CN4=((ANA4/AK7)*((SIX5-AK6*CIS6+M55M*ANB4)
XPT=AKN4*X
PHI1=AN4*COS(XPT)+CN4*SIN(XPT)
GO TO 10
5 CONTINUE
C
IS = 5
PHI1=AK5X*SIN(I MODE*AK15*(XPNT+SOH))/(1.0*I MODE)
CONTINUE
RETURN
END

```

```

SUBROUTINE MODEL1
REvised 12 JUNE 1971

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```

THIS ROUTINE WILL:
1 CALCULATE THE CONSTANT ( NON-HARMONIC ) FACTORS IN THE
FOURIER POSITION EXPRESSIONS FOR THE INDICATED SOURCE
GEOMETRY( IS ) AS STATED IN THE PROBLEM
2 CALCULATE THE SUM OF THE INTEGRATED FOURIER EXPRESSION
TO DETERMINE THE TRUNCATION ERROR FOR THE REMAINING TERMS
3 OF THE FOURIER EXPANSION
ADJUST THE VALUE OF THE SOURCE STRENGTH ( SZFR0 ) TO OBTAIN
AN APPROXIMATION TO WITHIN + OR - 5% OF THE ORIGINAL SOURCE
A LIMIT OF 10 INTERACTIONS WILL BE DONE. IF THIS IS NOT
SPECIFIED BY THE VARIABLE NANS( 4 ) IN THE INPUT DATA
4 IF LOGICAL L( 9 ) = 1 A SIMPLE PLOT OF THE TOTAL SUM VS
POSITION ( X ) WILL BE DONE ON THE PRINTER
FUNCTION ON THE PRINTER - PERFORM SIMPLE PLOT OF SOURCE
EXTERNAL CONTROL - LIMIT TO ALLOWABLE NUMBER
NANS( 4 ) - OF INTERACTIONS IN ROUTINE MODEL1

```



```

203 C AK12= PI/(2.0*DELTA2)
      AK12=AK12*0.5
      AN2X= 4.0*S02/PI
      GO TO 215
      CONTINUE
      IS=3
      AK13= SQRT(BUCKLE)
      AK13=AK13*0.5
      AN3A= 1.0/DIFLEN
      AN3B=EXP(-1.0*SOH/DIFLEN)
      S03= SZERO/(2.0*DIFLEN*(1.0-AN3B))
      AN3C = (S03/SOH)*((DIFLEN)**2)
      GO TO 215
      CONTINUE
204 C IS=4
      S04=SZERO/(2.0*DIFLEN)*(1.0-EXP(-1.0*(SOH+S1)/DIFLEN))
      DELTA4= S1 + PI/(2.0*SQRT(BUCKLE))
      AK14= PI/(2.0*DELTA4)
      AK14=AK14*0.25
      ANC4= EXP(-1.0*SOH/DIFLEN)
      ANB4= SINH(S1/DIFLEN)
      ANA4= 2.0*S04*DIFLEN
      ANB4= ANB4*ANC4
      GO TO 215
      IS=5
      CONTINUE
205 C AK15=SQRT(BUCKLE)
      AK5X=2.0*SZERO/(SOH*PI)
      CONTINUE
215 C START APPROXIMATION ANALYSIS TO DETERMINE THE ERROR IN
      FOURIER APPROXIMATION FOR THE TOTAL MODES OF HARMONIC DATA TO
      CONSIDER FOR THIS PROBLEM
      GO TO (301,302,303,304,305),IS
301 C CONTINUE
      IS=1 CASE
      SUM1 = 0.0
      DO 95 MODE=1, MODES
      AKN1=(2*MODE-1)*AK11
      SUM1 = SUM1 + SIN(AKN1*SOH)*SIN(AKN1*SA)/((AKN1)**2)
      CONTINUE
95 C SUM1 = SUM1*2.0*SOI/SOH
      ER1= SZERO-SUM1
      COR1= ER1/SUM1
      CORF= 1.0+COR1
      WRITE (6,98) SZERO,IS, MODES,SUM1,ER1,COR1,CORF
      IF ((SUM1.LE.SPIPI).AND.(SUM1.GE.SPZP9)) GO TO 1015
      GO TO 315
302 C CONTINUE

```



```

C
IS=2 CASE
SUM2 = 0.0
DEL1 = SOH-S1
DO 295 MODE=1, MODES
AKN2= AK12*(2*MODE-1)
BX2 = 2.0*SIN(AKN2*SOH)*COS(AKN2*S1)
BX3 = SIN(AKN2*SA)/(2*MODE-1)**2)
BX4 = SIN(AKN2*DEL1)
BXSUM=BX3*(BX2-BX4)
SUM2=SUM2 + BXSUM
CONTINUE
295 SUM2*SO2/(SA*(PI**2))
ER2= SZERO-SUM2
COR2= ER2/SUM2
CORF=1.0+COR2
WRITE(6,98) SZERO, IS, MODES, SUM2, ER2, COR2, CORF
IF (( SUM2.LE.SP1PI).AND.(SUM2.GE.SPZP9)) GO TO 1015
GO TO 315
CONTINUE
303 IS=3 CASE
ABX3 = 4.0*S03*((DIFLEN)**2)/(SOH*AK13)
SUM3 = 0.0
DO 395 MODE=1, MODES
M35 = (-1)**(MODE+1)
M36 = 2*MODE - 1
AKN3 = AK13*M36
BX35 = M35*(AKN3*EXP(-1.0*SOH/DIFLEN)*M35-1.0/DIFLEN)
BX36 = M36*(1.0+((AKN3*DIFLEN)**2))
BX37 = BX35/BX36
SUM3 = SUM3 + BX37
CONTINUE
395 SUM3*ABX3
ERR3 = SZERO-SUM3
COR3= ERR3/SUM3
CORF=1.0+COR3
WRITE(6,98) SZERO, IS, MODES, SUM3, ERR3, COR3, CORF
IF (( SUM3.LE.SP1PI).AND.(SUM3.GE.SPZP9)) GO TO 1015
GO TO 315
CONTINUE
304 IS=4 CASE
SUM4 = 0.0
DO 495 MODE=1, MODES
AKN4=AK14*(2*MODE-1)
AX5=AKN4*S1
AX6=AKN4*DIFLEN
AX7=EXP(-1.0*SOH/DIFLEN)*SINH(S1/DIFLEN)
BX55=COS(AX5)-AX6*SIN(AX5)+((-1)**MODE)*AX6*AX7
BX56=SIN(AKN4*SOH)/(1.0+((AX6)**2))

```



```

DO 5 IKX=51,99
TXBX(IKX)= DELX*(IKX-50)
CONTINUE
DO 6 NPT=1,100
DO 6 IMODE=1,MODES
FX6(NPT)=FX6(NPT)+PHI1(IMODE,IS,TXBX(NPT))
MODC=0
WRITE (6,1006) IS,IMODE,MODES
CALL PLOTP(TXBX,FX6,NPNTS,MODC)
CONTINUE
SZERO=SPZERC
25 CONTINUE
5001 FORMAT ('1',//,T3C,'SOURCE MODEL CORRECTIONS FOR FINITE SERIES',)
96 FORMAT ('//,T20,'INTERACTION NUMBER',2X,I3,2X,'FOR SOURCE TYPE',I5)
97 FORMAT ('//,T30,'ESTIMATED ERROR IN SOURCE MODEL',//,
1T30,'STRENGTH',T65,IPE10.3,/,T30,'SOURCE TYPE',
22X,'GEOMETRY',T65,I5,/,T30,'TOTAL MODES-FOURIER EXPANSION',
3T65,I3,/,T3C,'SUM OF FOURIER MODES',T65,IPE10.3,//,T30,
4,'ESTIMATED ERROR',T65,IPE10.3,/,T30,'FRACTIONAL ERROR',
5T65,IPE10.3,/,T30,'SOURCE NORMALIZATION FACTOR',T65,IPE10.3)
1006 FORMAT ('//,2CX,'SOURCE TYPE =',T4C,I5,/,
119X,'SUM OF FIRST',2X,I3,2X,'TERMS OF',2X,I3,2X,'MODES',)
RETURN
END

```

```

SUBROUTINE READ10(LASTN)
REVISD 2 JUNE 1971

```

```

INPUT ALL BASIC STATE PARAMETERS FROM MOD-5 PROBLEM OUTPUT
ECHO PRINT OF STATE STRUCTURE

```

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READ10 PROCESSES ALL PUNCH CARD OUTPUT FROM MOD-5
BY READING A BLOCK OF SIX DATA SETS FUNCTION F(U,T)
FOR UP TO SIX MODES OF BUCKLING AND TWENTY TIME STEPS
OF DATA PER BUCKLING MODE
FOR THE INPUT DATA -- A PRINTED OUTPUT IS GIVEN OF THE
MOD 5 STATE STRUCTURE NU(IA) AND E(IA) FOR
ALL ENERGY/LEISURE STATES OF THE SYSTEM
THIS OUTPUT IS LISTED TIME STEP DATA I.E. NT AND TIM(NT)
A PRINT OUT AT THE END OF EACH DATA INPUT SEQUENCE
IS GIVEN AT THE END OF THE PARTICULAR MODE IS DONE
AFTER ALL DATA FOR THE PART TO DETERMINE THOSE PROGRAMS
ROUTINE DOES A TEST COUNT BY CHECKING VALUES OF NT
DATA SETS WILL BE PROCESSED THE SAME VALUES OF NT
WHICH GAVE MOD 5 OUTPUT AT THE END OF THE READ SEQUENCE
A LISTING IS GIVEN AT THE END OF THE READ SEQUENCE
OF THE NUMBER OF TIME STEPS THAT WILL BE RETAINED.

```

CCCCCCCCCCCCCCCCCCCC


```

C THEN DEFAULT BY PRINTING OUT ALL REMAINING DATA THAT
C CANNOT BE PROCESSED INTO MILL SIX
C MOD FIVE PRODUCES THE FUNCTIONS F(U,T)
C DUMMY PARAMETER NRUN DEFINES INTERMEDIATE OR LAST DATA SET
C IF NRUN = 1 THIS IS INTERMEDIATE DATA SET
C IF NRUN = 2 THIS IS LAST DATA SET

IA=ISET
NTEST = NCALL
NCALL = 0
IKNT(ISET) = 0
WRITE (6,23)
21 NCALL = NCALL + 1
   IF (NCALL.EQ.21) GO TO 8
   READ (5,1,END=13,ERR=16) NRUN,TIM(NCALL),NT(NCALL),
   1 NTOP(IA,NCALL),NM(IA,NCALL)
   MOD=5 ALWAYS PUNCHES CARDS FROM THE DENSITY VECTOR IN THE
   FORM =NM(IA,1),IA=1,NM) IN A FORMAT IPELO.3 FOR 4 FIGURE DATA
   MILL=NM(IA,NCALL)
   IF (NRUN.EQ.1) IKNT(ISET) = IKNT(ISET) + 1

C TEST SERIES TO DETERMINE THE TOTAL DATA SETS WITH NRUN = 1
C THIS SHOULD REDUCE THE PROCESSING OF USELESS DATA LATER IN THE
C PROGRAM CALCULATIONS
C OUTPUT NEWS NOTE TO WRITE AT END OF THE CARD READING ROUTINE
C READ (5,2,END=15,ERR=16) (FUT(IA,IB,NCALL),IB=1,MILL)
C WRITE (6,5) NRUN,TIM(NCALL),NT(NCALL),NTOP(IA,NCALL),NM(IA,NCALL)
C IF (NRUN.EQ.1) GO TO 21
3 CONTINUE
IF (NRUN.EQ.2) WRITE(6,7)

C GO TO 15
8 CONTINUE

C THIS FIXES DATA SET 20 AS THE MAXIMUM AMOUNT FOR THE PROGRAM

MTEST = 0
9 WRITE (6,10)
11 READ (5,1,END=13,ERR=16) NNRUN,TIMN,NTTN,NTOPN,NMM
   IF (NNRUN.EQ.1) MTEST = MTEST + 1
12 READ (5,2) (DUMMY(IA),IA=1,NVIR)
13 IF (NNRUN.EQ.1) GO TO 11
   WRITE (6,7)
15 IF (NTEST.GE.NCALL) NCALL=NTEST
   CONTINUE
GO TO 19
20 CONTINUE

```



```

C 16 LASTN=1
C 19 CONTINUE
C 19 WRITE (6,91) MODES
C 19 DO 90 IX1 = 1, MODES
C 19 WRITE (6,92) IX1, IKNT(IX1)
C 90 CONTINUE
C
1  FORMAT (I5, E15.8, 3I5)
2  FORMAT (8E10.3)
5  FORMAT (10X, I5, 5X, 1PE10.3, 5X, 3I5)
7  FORMAT (///, 10X, 'END OF INPUT DATA FILE - F(U,T)')
10 FORMAT (///, 10X, 'NCALL EXCEEDS 20 - MAXIMUM SPACE IN PROGRAM')
23 FORMAT (///, 10X, 'F(U,T)-TIME STEP DATA', //, 11X, 'NRUN',
19X, 'TIM(NT)', 6X, 'NT', 2X, 'NTOPI', 3X, 'NM')
43 FORMAT (I5)
44 FORMAT (18A4)
45 FORMAT (4I5, E15.8)
46 FORMAT (8(1PE10.3))
49 FORMAT (10X, 2(I5, 5X, 2(1PE10.3, 8X), 6X))
50 FORMAT (10X, 18A4)
51 FORMAT (10X, 'NUMBER OF REAL STATES =', 3X, I5, /,
110X, 'NUMBER OF VIRTUAL STATES =', I5, /,
210X, 'NUMBER OF ISOTOPES =', 9X, I5, /,
310X, 'NUMBER OF FISSILE ISOTOPES =', 2X, I5, /,
410X, 'BUCKLING =', 15X, E10.5)
53 FORMAT ('1', 10X, 'INPUT DATA FROM MOD 5')
55 FORMAT (//, 8X, (2('STATE', 9X, 'U-LETHARGY', 6X, 'ENERGY--EV', 5X)))
56 FORMAT (10X, I5)
57 FORMAT (10X, I5)
91 115X, 'DATA SET', 5X, 'SETS')
92 FORMAT (17X, I3, 7X, I3)
92 RETURN
END
SUBROUTINE FOUT1(XPNT)
REVISD 2 JUNE 1971
SUBROUTINE FOUT1
CALCULATES THE POSITION FUNCTION F(U,T,X) FOR A SERIES OF SIX
BUCKLING SETS AT A TIME.
IF ONLY SIX BUCKLING MODES ARE CONSIDERED, THIS ROUTINE CAN DO
THE TIME-ENERGY-POSITION FUNCTION FOR AS MANY POINTS AS DESIRED
OF THIS ONE DIMENSIONAL MODEL -- SLAB GEOMETRY

```



```

C      CALCULATE THE FUNCTION F(U,T,X) = F(U,T)*PHI1(X) FOR ALL FUT
C      N5= NSTAT
C      X1 = XPNT
C      CHANGE TO FOUT1 CHANGE AFTER ROUTINE IS SET TO PROCESS MORE
C      THAN SIX INPUT DATA DECKS OR TAPE BY PROVIDING PROTECTED
C      STORAGE FOR THE FIRST MODE OF BUCKLING
C      BUCK1 = BUCK(1)
C      DO 2 IMOD=1,MODES
C      NCALL1 = IKNT(IMOD)
C      DO 2 NCX=1,NCALL1
C      MIL1 = NTOP(IMOD,NCX) + 1
C      MIL2 = MIL1 + NM(IMOD,NCX)
C      DO 2 M=MIL1,MIL2
C      FUTX(M,NCX)=FUT(IMOD,M,NCX)*PHI1(IMOD,IS,X1)
C      1 + FUTX(M,NCX)
C      2 CONTINUE
C      OVER RIDE STEP TO BY PASS THE OUTPUT ROUTINES ON CALLS FROM
C      OTHER SUPROUTINE AND SUB PROGRAMS OF MOD 2
C      OVER RIDE PARAMETER IS LOGICAL -- L( 16)
C      IF L(16) = 1 THE PRINT ROUTINE AND ALL OTHERS ARE SKIPPED
C      IF (L(16).EQ.1) GO TO 25
C      WRITE OUT ALL VALUES OF FUTX(M,NCX)
C      BEFORE CALLING THE CAL COMP PLOTTER
C      NCALL = IKNT(1)
C      WRITE (6,90) X,NVIR,NCALL
C      PLOT P SEQUENCE
C      DO 93 I12=1,NCALL
C      WRITE (6,94) I12,TIM(I12),NT(I12)
C      WRITE (6,95) (FUTX(I13,I12),I13=1,NVIR)
C      CONTINUE
C      93 IF(NCUT.NE.3) GO TO 25
C      IF THE TEST VARIABLE NOUT NOT EQ 3 BY PASS REST OF SEQUENCE
C      DEFINE SCALE PARAMETERS FOR THE DRAW PLOT ROUTINE
C      PAGES TO BE PLOTTED VIA THE CAL COMP
C      TOTAL LIMITS ON GRAPH PLOTTING FOR THIS ROUTINE IS DEFINED
C      VIA THE INPUT CONTROL VARIABLE NANS(1) = (( THE TOTAL GRAPH
C      DEFINE A DEFAULT VALUE FOR NANS( 1 ) OF 1
C      IF (NANS(1).LT.1) NANS(1) = 1
C      EXSCALE = U(NVIR-1)/8.
C      YSCALE = FUTX(2,1)/8.
C      IXUP = C
C      IYRIGH = C
C      MODXAX = C

```



```

MODYAX = 0
IWIDTH = 9
IHIGH = 10
IGRID = 1
KKNT = NCALL
KTEST = NCALL/5 + 1
KK1 = 0
DO 8 IKT = 1, KTEST
MODIFICATION OF GRAPH TITLING TO INDICATE DATA SET PLOTTED
INAME(10) = BMODES(MODES)
NPXNT = NANS(2)
INAME(11) = APNT(NPXNT)
IF(IKT.LE.4) INAME(12) = TIMST(IKT)
IF(IKT.GT.4) INAME(12) = TIMST(4)
DO 5 IXM = 1,5
KK1 = KK1 + 1
MODCUR = 2
IF (IXM.EQ.1) MODCUR = 1
IF (IXM.EQ.5) MODCUR = 3
IF (KK1.EQ.NCALL) MODCUR=3
IF (KK1.GT.20) GO TO 3
INMM = LABEL(KK1)
3 IF(KK1.GT.20) INMM = LABEL(20)

CALL DRAW USING THE WORKING VECTOR FFUTX
USE OF THE INPUT CONTROL PARAMETER LIMITS GRAPH
TO SOME MULTIPLE OF FIVE DATA PLOTS IF POINT PLOTTING
PARAMETER USED IS NANS(1)
TEST QUESTION TO CHECK : IF (NANS(1).EQ.1.AND.IKT.EQ.6) GT TO10
INITIAL TEST QUESTION FORM::: THE FUTURE:::
SECOND FORM TO USE IN TI THE FUTURE:::
IF (NTST2.EQ.IKT) GO TO 10
PARAMETER NTST2 DEFINES THE TOTAL NUMBER OF
TIMES MOD5 THAT PLOT WILL BE CALLED

THREE DEFAULT OPTIONS TO DROP OUT OF THE CALL DRAW PROGRAM

DO 4 ICX = 1, NVIR
FFUTX(ICX) = FUTX(ICX, KK1)
4 IF ((NANS(1).EQ.1).AND.(IKT.EQ.6)) GO TO 10
IKT1 = IKT
IF (NVIR.GE.30) IKT1 = 0

CALL DRAW (NVIR, U, FFUTX, MODCUR, IKT1, INMM, INAME, EXSCAL, YSCALE,
1 IXUP, IYRIGH, MODXAX, MODYAX, IWIDTH, IHIGH, IGRID, LAST)
IF (LAST.EQ.3) GO TO 10

```



```

C
IF (KK1.EQ.NCALL) GO TO 10
ON REACHING MAXIMUM GRAPH OUTPUT-DROP OUT OF THE LOOP
IF (NANS(1).GE.IKT) GO TO 10
CONTINUE
CONTINUE
CONTINUE
CONTINUE
CONTINUE
90 FORMAT ('1',10X,'CALCULATION OF THE POSITION FUNCTIONF(U,T,X)',
1/,20X,'FOR X=',2X,F10.5,2X,'WITH',2X,I5,2X,'STATES AT',2X,
2 I5,2X,'TIMES,')
94 FORMAT ('//',10X,'RUN NUMBER',2X,I5,2X,'TIME =',2X,E15.8,
12X,'SECONDS',5X,'NT =',2X,I5)
95 FORMAT ('41(',10X,5(1PE10.3,5X))
RETURN
END

```

```

C
SUBROUTINE FOUT2
REVISD 16 JUNE 1971
CORRECTED ROUTINE TO CALCULATE THE MEAN ENERGY OF THE
NEUTRONS AT A POINTX(I) AND THE SUM OVER ALL LETHARGY STATES
OF THE NEUTRON DISTRIBUTION AT THE SAME POINT AND PLOT BOTH
RESULTANT ON A SINGLE GRAPH FOR EACH TIME STEP VALUE
LETHARGY SUM
F(U,T,X,MODE) = F(U,T,MODE)*PHI(X,MODE)
F(U,T,X) = SUM OVER MODES--F(U,T,X,MODE)
OVER ALL STATES U(LETHERGY) TO GIVE THE RESULTANT
FUNCTION -- F ( X,T )
MEAN ENERGY CALCULATION
F(X,T) = SUM OVER LETHARGY - F(U,T,X)
EMEAN(X,T) = SUM OVER LETHARGY-F(U,T,X)*E(U)
COMMON BLOCKS 3 AND 6 MUST BE REVISED TO HANDLE MORE THAN
SIX HARMONIC MODES OF THE FUNCTIONAL DATA INPUT
DIMENSION VSIG(150)
EQUIVALENCE (DUMMY(1),VSIG(1))
REAL*8 ENGY,'ENERGY',ENGX,'RATIO',
REAL*8 FNGY,'NEUTRON',FNGX,'DENSITY',
REAL*8 ARUNT(9),'SYSTEM-1',SYSTEM-2,'SYSTEM-3',SYSTEM-4',
1 REAL*8 MTT(12),'FUNCTION',F(X,T),
1 REAL*8 MTT(12),'FUNCTION',F(X,T),
REAL*8 TIMX(20),'NT - 1',NT - 2,'NT - 3',NT - 4,'NT - 5',
2,'NT - 6',NT - 7,'NT - 8',NT - 9,'NT - 10',NT - 11,'NT - 12',
12,'NT - 13',NT - 14,'NT - 15',NT - 16,'NT - 17',NT - 18',
3,'NT - 19',NT - 20',
REAL*8 BMODES(6),'1-MODE', '2-MODES', '3-MODES', '4-MODES',

```



```

NPNTS=99
DELX2=SOH/49.0
TXBX(1)=(-1.)*SOH
TXBX(99)=SOH
DO 3 I2=2,98
TXBX(I2)=TXBX(1)+DELX2*(I2-1)
CONTINUE
TXRX(50)=0.0
3 CONTINUE
C
CP2
WRITE (6,1002)
DEFINE INITIAL CONSTANTS FOR THE ROUTINE
C
NONE = 1
BKI=BUCK(1)
NTMAX=IKNT(1)
NMAX = NVIR-1
IMPROVEMENT TO INCREASE EXECUTION SPEED OF CALCULATION
DO 321 NXPT=1,NPNTS
DO 321 KMOD=1,MODES
FUTX(NXPT,KMOD)=PHI1(KMOD,IS,TXBX(NXPT))
321 CONTINUE
C
CP3
WRITE (6,1003)
START OF CALCULATIONS
DO 250 INT=1,NTMAX
DO 245 IXPT=1,NPNTS
AX=0.0
BX=0.0
CX5=C.0
DO 240 IMODE=1,MODES
PROTECTION $STEP WITHIN CALCULATION LOOP TO REDUCE UNNECESSARY
CALCULATIONS
C
NP3=IKNT(IMODE)
IF(INT.GT.NP3) GO TO 240
MIL1=NTOP(IMODE,INT)
MIL2=NM(IMODE,INT)
MIL3= MIL1+MIL2
DO 235 NUI=1,MIL2
MIL4=MIL1+NUI
IF (MIL4.GT.NMAX) GO TO 235
BX1=FUT(IMODE,NUI,INT)*FUTX(IXPT,IMODE)
C
CX5=CX5+BX1
AX=AX+BX1*VSIG(MIL4)
BX=BX+BX1*E(MIL4)
235 CONTINUE
240 CONTINUE
C
DEFAULT PROTECTION STEP FOR PARAMETER CX5

```



```

C THIS VARIABLE IS TO REPRESENT THE NEUTRON POPULATION DENSITY
C FOR A PARTICULAR TIME STEP VALUE
C IN GENERAL, IT SHOULD BE A VALUE BETWEEN 0.001 AND 1.000
C FOLLOWING IS TO PROTECT AGAINST AN ACCIDENTAL ATTEMPT TO DIVIDE
C BY ZERO OR SOME OTHER UNUSUAL VALUE THAT MIGHT ACCIDENTLY
C COME UP.
C TEST CHECK AND DEFAULT SET:
C IF ((CX5.LE.0.001).OR.(CX5.GE.1.0001)) CX5=1.000
C
C EMEAN(IXPT,INT)=BX/CX5
C FXT(IXPT,INT)=AX/CX5
C CONTINUE
C CONTINUE
C CP4
C WRITE (6,1004)
C DO 253 NT4=1,NTMAX
C IF ((IS.EQ.1).OR.(IS.EQ.3)) ZA=EMEAN(1,NT4)
C IF ((IS.EQ.2).OR.(IS.EQ.4).OR.(IS.EQ.5)) ZA=EMEAN(50,NT4)
C DEFAULT TEST CHECK CN ZA VALUE
C WRITE (6,6001) NT4,ZA
C
C IF (ZA.LE.0.0) GO TO 255
C DO 252 IX4=2,NPNTS
C FUTX(IX4,NT4) = EMEAN(IX4,NT4)/ZA
C CONTINUE
C FUTX(1,NT4) = 1.0
C GO TO 253
C 252 WRITE (6,271) NT4,ZA
C DO 256 IX5=1,NPNTS
C FUTX(IX5,NT4) = .90
C CONTINUE
C CP5
C WRITE (6,1005)
C BY PASS THE CALL COMP ROUTINE PORTION IF NOUT NOT EQUAL 3
C IF (NOUT.NE.3) GO TO 12
C START PREPARATION OF DATA FOR GRAPH PLOTTING ROUTINES
C TWO ITEMS OF DATA WILL BE PLOTTED FOR EACH TIME STEP
C LAST=C
C ITYPE=C
C XCAL=SOH/8.0
C IHIGH=10
C IWIDE=8
C IYRT=C
C IGRID=1
C FXT(I,J) = FLUX PROFILE FOR X AND T
C EMEAN(I,J) = MEAN ENERGY DATA VERSUS X AND T
C FUTX(I,J) = RATIO MEAN ENERGY(X,T)/MEAN ENERGY(0.0,T)

```



```

C      L(K) = 1
C      L(7) = 1
C      L(8) = 1
C      ADJUSTMENT = 1
C      MTIT(10) = BMODES(MODES)
C      LOAD DUMMY WITH FXT(NX,NT) DATA
C      DO 155 NXT1=1, NTMAX
C      LABEL=LTAB(NXT1)
C      MTIT(12) = TIMX(NXT1)
C      IXUP1=0
C      YSCAL1=FXT(1,NXT1)/4.0
C      MODC1=1
C      MODX1=0
C      IXUP1=C
C      IYRT=0
C      MODY1=0
C      PROTECT=0
C      PLOTTING PROTECTION STEP TO MAINTAIN ONLY POSITIVE VALUES
C      IF (L(7).EQ.C) GO TO 1175
C      DO 162 NP5=1, NPNTS
C      DUMMY(NP5)=FXT(NP5, NXT1)
C      IF (DUMMY(NP5).LE.C.001) DUMMY(NP5)=0.0
C      CONTINUE
C      CALL DRAW TO PLOT FIRST GRAPH FOR THIS TIME STEP
C      MTIT(2)=FNGY
C      MTIT(3)=FNGX
C      CALL DPAW(NPNTS, IXBX, DUMMY, MODC1, ITYPE, LABEL, MTIT, XCAL,
1175  YSCAL1, IXUP1, IYRT, MODX1, MODY1, IWIDE, IHIGH, IGRID, LAST)
C      IF (LAST.EQ.C) GO TO 155
C      IF (L(8).EQ.C) GO TO 1176
C      MODX2=2
C      YSCAL2=0.50
C      IXUP2=6
C      MODC2=3
C      MODY2=0
C      MUDY2=0
C      DU 163 NP6=1, NPNTS
C      DUMMY(NP6)=FUTX(NP6, NXT1)
C      CALL TO DRAW TO PLOT THE MEAN ENERGY RATIO
C      MTIT(2)=FNGY
C      MTIT(3)=FNGX
C      CALL DRAW(NPNTS, TXBX, DUMMY, MODC2, ITYPE, LABEL, MTIT, XCAL,
1176  YSCAL2, IXUP2, IYRT, MODX2, MODY2, IWIDE, IHIGH, IGRID, LAST)
C      CP6
C      INTERNAL PROTECTION STEP TO LIMIT TOTAL GRAPHICAL OUTPUT TO
C      THE CAL COMP PLOTTER VIA VARIABLE NANS ( 3 ) TO COMPARE
C      WITH THE INDEXING VARIABLE ---NXT1
C      KPLOTS IS A DUMMY VARIABLE DEFINED WITH A DEFAULT VALUE OF 1
C      CONTINUE

```



```

155 IF ( NXT1.GE.KPLOTS ) GO TO 156
156 CONTINUE
C WRITE (6,IC6)
PLOTP OUTPUT ROUTINE USING THE PRINTER
12 IF (NOUT.EQ.3) GO TO 16
NCALL = IKNT(1)
DO 15 J=2,NCALL
DO 13 I=1,NPNTS
13 DUMMY(I) = FXT(I,J)
MODC = 0
CALL PLOTP(TXBX,DUMMY,NPNTS,MODC)
WRITE (6,14) J,NT(J),TIM(J)
15 CONTINUE
PRINT-PLOT THE MEAN ENERGY VALUES OF THE PRINTER
DO 615 J=2,NCALL
DO 613 I=1,NPNTS
613 DUMMY(I) = EMEAN(I,J)
MODC=0
CALL PLOTP(TXBX,DUMMY,NPNTS,MODC)
WRITE (6,617) J,NT(J),TIM(J)
615 CONTINUE
PRINT-PLOT THE MEAN ENERGY RATIO
DO 715 J=2,NCALL
DO 713 I=1,NPNTS
713 DUMMY(I) = FUTX(I,J)
MODC=0
CALL PLOTP(TXBX,DUMMY,NPNTS,MODC)
WRITE (6,717) J,NT(4),TIM(J)
715 CONTINUE
PRINT OUT RESULTS ON THE PRINTER ONLY
PRINT OUT ALL VALUES OUT IN COLUMNS OF SIX TIME STEPS
C L(5) = 1
C IF(L(5).EQ.1) GO TO 23
C OVER-RIDE CONTROL TO SKIP THE PRINT SEQUENCE
35 CONTINUE
NCALL = IKNT(1)
IF(NCALL.GT.6) NTT=6
IF(NCALL.LE.6) NTT=NCALL
WRITE (6,22) (NT(I),I=1,NTT),(TIM(I),I=1,NTT)
WRITE (6,501)
DO 17 I=1,NPNTS
17 WRITE (6,21) TXBX(I),(FXT(I,J),J=1,NTT)
WRITE (6,502)
DO 67 I2=1,NPNTS
67 WRITE (6,21) TXBX(I2),(EMEAN(I2,J2),J2=1,NTT)
WRITE (6,503)
DO 93 I2=1,NPNTS

```



```

93 WRITE (6,21) TXBX(I2), (FUTX(I2,J2), J2=1, NTT)
   IF (NCALL.LE.6) GO TO 23
   IF (NCALL.GT.12) NTT=12
   IF (NCALL.LE.12) NTT = NCALL
   WRITE (6,22) (NT(I), I=7, NTT), (TIM(J), J=7, NTT)
   WRITE (6,501)
DO 18 I=1, NPNTS
   WRITE (6,21) TXBX(I), (FXT(I,J), J=7, NTT)
   WRITE (6,502)
DO 68 I2=1, NPNTS
   WRITE (6,21) TXBX(I2), (EMEAN(I2,J2), J2=7, NTT)
   WRITE (6,503)
DO 94 I2=1, NPNTS
   WRITE (6,21) TXBX(I2), (FUTX(I2,J2), J2=7, NTT)
   IF (NCALL.LE.12) GO TO 23
   IF (NCALL.GT.18) NTT=18
   IF (NCALL.LE.18) NTT = NCALL
   WRITE (6,22) (NT(I), I=13, NTT), (TIM(J), J=13, NTT)
   WRITE (6,501)
DO 19 I=1, NPNTS
   WRITE (6,21) TXBX(I), (FXT(I,J), J=13, NTT)
   WRITE (6,502)
DO 69 I2=1, NPNTS
   WRITE (6,21) TXBX(I2), (EMEAN(I2,J2), J2=13, NTT)
   WRITE (6,503)
DO 95 I2=1, NPNTS
   WRITE (6,21) TXBX(I2), (FUTX(I2,J2), J2=13, NTT)
   IF (NCALL.LE.18) GO TO 23
   NEXT = C
   WRITE (6,22) NT(19), NT(20), NXT, NXT, NXT, TIM(19), TIM(20)
   WRITE (6,501)
DO 20 I=1, NPNTS
   WRITE (6,21) TXBX(I), FXT(I,19), FXT(I,20)
   NTT=NCALL
DO 70 I2=1, NPNTS
   WRITE (6,21) TXBX(I2), (EMEAN(I2,J2), J2=19, NTT)
   WRITE (6,503)
DO 96 I2=1, NPNTS
   WRITE (6,21) TXBX(I2), (FUTX(I2,J2), J2=19, NTT)
CONTINUE
11 FORMAT (//, 10X, I5, 2X, 'TOTAL POINT S PLOTTED USING DRAW', /,
14 FORMAT (//, 10X, I5, 2X, 'TOTAL CALLS MADE TO DRAW', /,
12X, I5, /, 10X, 'DATA PLOT NUMBER', 2X, I5, /, 10X, 'TEST RUN NT =',
21 FORMAT (11X, F6.2, 6X, 6(1PE10.3, 4X))
22 FORMAT (11X, //, 11X, 'X =', 8X, 6(4X, 'NT =', I3, 1X), //,
122X, 6(1X, E10.3, 3X), /, 22X, 6(5X, 'SEC', 6X))

```



```

271 FORMAT (10X,'ERROR CHECK',2X,'NT = ',I4,2X,'EMEAN(1,NT) = ',
12X,1PE10.3)
501 FORMAT (//,10X,'FUNCTION-F( X,T )--SUM OVER LETHARGY')
502 FORMAT (//,10X,'OUTPUT DATA FROM FOUT2',//,
110X,'MEAN ENERGY VS POSITION')
503 FORMAT (//,10X,'MEAN ENERGY RATIO')
617 FORMAT (//,10X,'MEAN ENERGY VS POSITION',//,
120X,'PLOT TIME = ',I5,3X,'NT = ',I5,10X,'TIME = ',2X,1PE10.3,
22X,'SECONDS')
717 FORMAT (//,10X,'MEAN ENERGY RATIO',/,20X,'PLOT NUMBER',
12X,15,3X,'NT = ',I5,10X,'TIME = ',2X,1PE10.3,2X,'SECONDS')
1001 FORMAT (30X,'FOUT2 ---- CHECK POINT 1')
1002 FORMAT (30X,'FOUT2 ---- CHECK POINT 2')
1003 FORMAT (30X,'FOUT2 ---- CHECK POINT 3')
1004 FORMAT (30X,'FOUT2 ---- CHECK POINT 4')
1005 FORMAT (30X,'FOUT2 ---- CHECK POINT 5')
1006 FORMAT (30X,'FOUT2 ---- CHECK POINT 6')
6001 FORMAT (/,10X,'TIME STEP VALUE',2X,I3,5X,'MEAN ENERGY = ',
15X,1PE10.3,2X,'ELECTRON VOLTS')
      RETURN
      END

```

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SUBROUTINE FOUT3
  REVISED 2 JUNE 1971

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5 MAY 1971 FIRST DRAFT OF DETECTOR RESPONSE FUNCTION
THIS ROUTINE WILL REQUIRE THE MODIFICATION FOR THE NAMELIST
DATA ENTERED IN SUBROUTINE INCONI TO REFLECT THE ADDITION
OF TWO ITEMS OF INFORMATION --- THE MAXIMUM AND MINIMUM ENERGY
THAT THE DETECTOR WILL RECOGNIZE
WE NOW MUST CONSIDER THE MINIMUM INTENSITY THE
DETECTOR CAN RESPOND TO-- AS WITH SEVERAL OTHER OF THE
ROUTINES, WE WILL SELECT AN ARBITRARY VALUE OF 0.001 OF THE
NORMALIZED VALUE
THIS NORMALIZED LIMIT IS SELECTED AS A RESULT OF THE AVAILABLE
ACCURACY AND PRECISION OF THE CAL COMP PLOTTER WHICH CAN PLOT
TO 0.001 OF AN INCH OR SOME VALUE WITHIN THAT RANGE
COMMON BLOCKS 3 AND 6 MUST BE REVISED TO HANDLE MORE THAN
SIX HARMONIC MODES OF THE FUNCTIONAL DATA INPUT
REAL ESIG(27)/1.0E+7,6.5E06,4.0E06,2.5E06,1.4E06,8.0E05,4.0E05,
24.0E05,1.0E05,4.65E04,2.15E04,1.0E04,4.65E03,2.15E03,1.0E03,
24.0E02,2.15E02,1.0E02,4.65E01,2.15E01,10.0,4.65,2.15,1.0,0.465,
30.0,215,2.5E-02/
REAL SIGMAS(26)/1.2,1.65,2.20,3.0,4.10,5.70,8.10,11.0,14.0,
116.6,18.5,19.3,19.7,20.0,20.1,20.2,20.2,9*20.3/
REAL LABEL
REAL LABEL(2)/'F(X)':'FXT'/

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CCCCCCCCCCCCCCCC

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1015 ATIMX(IK6)=SC*(ALOG10(TIM(IK6))+10.0)
      CONTINUE
      CALL PLOT(PDUM2,ATIMX,NTCALL,MODC)
      WRITE (6,155) ((DUMMY(IA),TIM(IA)),IA=1,NTCALL)
      CALL PLOT(PDUM1,ATIMX,NTCALL,MODC)
      DEFINE SCALE VARIABLES FOR THE DRAW--CAL COMP PLOTTER
      OVER RIDE CALL DRAW SEQUENCE UNLESS NOUT = 3
      IF (NOUT.NE.3) GO TO 1003
      MCDX=0
      MODY=0
      IYRT=9
      IXUP=0
      ITYPE=1
      EXSCL=1.0/9.0
      YSCAL=TIM(NTCALL)/14.0
      PREPARE DATA TO PLOT REFLECTED AND REVERSED AXIS ON THE PLOTTER
      DO 151 KPNT=1,NTCALL
      IF (TXBX(KPNT).LE.0.001) TXBX(KPNT)=0.0
      TXBX(KPNT)=(-1.0)*TXBX(KPNT)
151  CONTINUE
      IWIDE=9
      IHIGH=15
      IGRID=1
      PLOT F(X,T) ON THE GRAPH FIRST
      MODC1=1
      CALL DRAW(NTCALL,TXBX,TIM,MODC1,ITYPE1,LABL(1),TITLE,EXSCL,
1  YSCAL,IXUP,IYRT,MODX,MODY,IWIDE,IHIGH,IGRID,LAST)
      PLOT F(X) VS TIME
      PREPARE DATA BY REFLECTING ABOUT Y-AXIS
      DO 152 KPNT=1,NTCALL
      IF ( DUMMY(KPNT2).LE.0.001) DUMMY(KPNT2)=0.0
      DUMMY(KPNT2)=(-1.0)*DUMMY(KPNT2)
152  CONTINUE
      ITYPE2=2
      MODC2=3
      CALL DRAW(NTCALL,DUMMY,TIM,MODC2,ITYPE2,LABL(2),TITLE,EXSCL,
1  YSCAL,IXUP,IYRT,MODX,MODY,IWIDE,IHIGH,IGRID,LAST)
1003  CONTINUE
154  FORMAT (//,10X,'PRELIMINARY DATA PLOTS',/,10X,'TIME INTERGRATED',
1  1X,'FUNCTION--F(X) VS TIME',/,20X,'F ( X )',T35,'TIME--SECONDS',/,
2  2C(//,18X,1PE10.3,8X,1PE10.3))
155  FORMAT (//,10X,'PLOT OF DETECTOR RESPONSE FUNCTION',/,
1  120X,'F ( X,T ) VS TIME',/,20X,'F(X,T)',T35,'TIME--SECONDS',
2  20(//,18X,1PE10.3,8X,1PE10.3))
      RETURN
      END

```



```

YSCALE = U(NMAX)/14.0
IXUP = 0
IYRT = 0
MODX = 0
MODY = 0
IWIDE = 9
IHIGH = 15
IGRID = 1
DO 300 J=1, MODES
  TITLE(6) = RMOD(J)
DO 92 NXX = 1, 150
  DUMMY(NXX) = 0.0
  FX=0.0
  NCALL1 = IKNT(J)
DO 55 K=2, NCALL1
  NNU=NM(J,K)
  NNU DEFINES THE MAXIMUM STEP VALUES INPUT FROM MOD 5
  FX = 0.0
DO 50 NU=1, NNU
  IF(FUT(J,NU,K) .GT. FX) FX=FUT(J,NU,K)
  THIS SCAN WILL DETERMINE THE MAX VALUE IN THE INPUT DATA
  CX = BX*FX
  THIS WILL DETERMINE THE MINIMUM VALUES TO CONSIDER
  WRITE(6,1C8) CX, BX, FX, K
  FORMAT(//,10X,'CX = ',2X,1PE10.3,2X,'BX = ',2X,1PE10.3,
12X,'FX = ',2X,1PE10.3,2X,'TIME STEP=',2X,I3)
C
MIL1 = NTOP(J, K)
MIL2 = NM(J, K)
MIL3 = MIL1 + MIL2
DO 51 NU=1, NNU
  AX = FUT(J, NU, K)
  SET MINIMUM VALUE FOR LOG CONVERSION
  IF (AX.LT.BX) AX = BX
  MIL4 = MIL1 + NU
  FUTX(MIL4, K) = SC*ALOG10(AX) + 9.0
C
51 CONTINUE
  THIS STEP WILL AUTOMATICALLY ZERO OUT ANY POINT AT THE
  LOW END OF THE SPECTRUM(HIGH-ENERGY; LOW - LETHARGY) TO MAINTAIN
  THE CORRECT PROFILE OF THE SPECTRUM
  IF (MIL1.LE.2) GO TO 53
DO 52 KN = 1, MIL1
  FUTX(KN, K) = 0.0
52 CONTINUE
53 CONTINUE
55 CONTINUE

```



```

C START OF THE CALL TO DRAW SERIES
C NCALL1 = IKNT(J) - 1
C THIS SHOULD REMOVE ERRORS IN DRAW CALL
LAST = 0
DO 60 J6 = 2, NCALL1
IF (J6.EQ.2) MODC = 1
IF ((J6.GT.2).AND.(LAST.EQ.0)) MODC = 2
IF (J6.EQ.NCALL1) MODC = 3
IF (LAST.EQ.2) MODC = 1
DO 59 K1 = 1, NPNTS
59 DUMMY(K1) = FUTX(K1, J6)

C DO FINAL DATA CHECK TO INSURE THAT PLOT POINTS ARE WITH
C IN THE LIMITS OF +0.001 AND 1.0
C DO 73 K2 = 1, NPNTS
IF (DUMMY(K2).GT.9.0) DUMMY(K2) = 9.0
IF (DUMMY(K2).LT.0X) DUMMY(K2) = 0.0
73 CONTINUE
END OF DATA PLOT CHECK ROUTINE
LABEL = 0
LTAB(J6)
LAST = 0
CALL DRAW(NPNTS, DUMMY, U, MODC, ITYPE, LABEL, TITLE, EXSCAL,
1 YSCALE, IXUP, IYRT, MODX, MODY, IWIDTH, IGRID, LAST)
IF (LAST.EQ.1) GO TO 105
IF (LAST.EQ.3) GO TO 105
60 CONTINUE
105 CONTINUE
IF (LAST.EQ.1) WRITE(6, 106)
IF (LAST.EQ.3) WRITE(6, 107)
106 FORMAT (//, 10X, 'LAST = 1, 3X, 'RECHECK DATA')
107 FORMAT (//, 10X, 'LAST=3, 3X, 'RECHECK DATA')

C CONTINUE
C RETURN
C END

SUBROUTINE FOUT5
REVISED 15 JUNE 1971

SUBROUTINE FOUT5 PROVIDES A FOURIER POSITION FUNCTION
WEIGHTED SPECTRAL PLOT OF THE INPUT DATA FROM MOD-5 AND
PLOTS F(I,MODE,U,I)*PHI(IMODE,XI) VS U(LETHAGY) FOR EACH
MODE OF INPUT DATA AND A SUMMED SPECTRUM OF ALL MODES OF INPUT
( 4 ) = 1 DEFINES THE EXCEPTIONAL CASE WHERE THE USER WOULD
DESIRE TO PLOT THE SPECTRAL RESPONSE FUNCTION WEIGHTED BY THE

```



```

C FOURIER EXPANSION COEFFICIENT FOR EACH MODE AND POSITION OTHER
C THAN AT THE DEFAULT CASE OF X = C.0
C FACTORS ( POSITION ) THAT THE PROGRAM WILL CONSIDER IS
C INPUT USING THE CONTROL VARIABLE NANS ( 6 ) FOR THE NUMBER OF
C POINTS AND THE INPUT ARRAY XBXC ( I ) FOR THE ACTUAL POINTS
C PRE CALCULATE THE FOURIER SPACE DEPENDENT VALUES FOR EACH MODE
C FOR EACH POINT XI
COMMON/COM1/ XBXC(100),U(150),E(150),TIM(20),DUMMY(150),TXBX(100)
COMMON/COM4/ SOH,SA,S1,S2,S3,S4,SZERO,BUCKLE,EMAX,EMIN,DIFFLEN,X
COMMON/COM7/ SOL,S02,S03,S04,AKI1,AKI2,AKI3,AKI4,AN3A,AN3B,
IAN3C,ANC4,ANA4,ANB4,AN2X,AK15,AK5X
COMMON/COM2/ NT(20),NPROB(54),NANS(6),L(20),LTAB(20)
COMMON/COM5/ MODES,IS,IDATA,N,NVIR,NI,NF,NCALL,NOUT,IBX,
INSTAT,MODE
COMMON/COM3/ FXT(150,20),FUTX(150,20),BUCK(6),FUT(6,71,20)
COMMON/COM6/ NTOP(6,20),NM(6,20),IKNT(6)
REAL*8 RTITLE(12),MILLS,BOX,M,,'FOUT5',,' ',,' ',
1,TOTAL,SUM OF,ALL,MODES,,' ',,' ',,' ',,' ',,' ',
1,REAL*8 ANAME(12),MILLS,BOX,M,,'FOUT5',,' ',,' ',,' ',,' ',
1,REAL*8 BMODE(5),,' 1 - 2',,' 1-2-3',,' 1-2-3-4',,' 12345',,' 123456',/
REAL*8 APNT(9),POINT-1,POINT-2,POINT-3,POINT-4,
1,POINT-5,POINT-6,POINT-7,POINT-8,POINT-9,
1,REAL*8 AMODE(6),MODE-1,MODE-2,MODE-3,MODE-4,MODE-5,
1,MODE-6,/,
DIMENSION PHIMX(6)
START EXECUTION DEFINE PARAMETERS--SCALE FACTORS,ETC FOR EXEC
C
IPNTR=NANS(2)
IF ((IPNTR.LE.0).OR.(IPNTR.GE.9)) IPNTR=9
RTITLE(12)=APNT(IPNTR)
ANAME(12)=APNT(IPNTR)
NRUNS=NANS(6)
THIS DEFINES THE TOTAL NUMBER OF DATA POINTS THAT WILL
HAVE SPECTRAL RESPONSE DATA POINTS PLOTTED
IF ((NRUNS.LE.0).OR.(NRUNS.GE.3)) GO TO 5500
IF (L(4).EQ.0) NRUNS = 1
DO 5001 NTIMX = 1,NRUNS
X1= XBXC(NTIMX)
IF (L(4).EQ.0) X1=0.0
WRITE (6,1013) X1
FORMAT (//,10X,POSITION POINT FOR SPECTRAL PLOT =',5X,F8.2)
DO 1011 JXX = 1,MODES
PHIMX(JXX)=PHI1(JXX,IS,X1)
WRITE(6,1012) JXX,PHIMX(JXX)
FORMAT (//,10X,MODE =',2X,15,PHIX = ',E15.8)
1012 CONTINUE
1011

```



```

C          SC = 3.0
C          IXTYPE = 0
C          NPNTS = NVIR-1
C          TEST SERIES CORRECTION FOR MIN SCALING VALUE TO PLOT
C          THIS SHOULD BE ADJUSTED TO BECOME PART OF THE INPUT
C          NAMELIST DATA-- DATAIN --- TO ALLOW THE USER TO SELECT
C          AN ESTIMATED BEST VALUE ---- THIS WILL REQUIRE ADDITION
C          TO ONE OF THE PRINCIPAL COMMON BLOCKS TO TRANSFER
C          PLUS THE ESTABLISHMENT OF A DEFAULT OPTION IN INCON1
C          REVISED 15 JUNE 1971
C          BX=1.0E-05
C          YSCALE = U(NPNTS)/14.0
C          IXUP = 0
C          IYRT = 0
C          MODX = 0
C          MCDY = 0
C          IWIDE = 9
C          IHIGH = 15
C          IGRID = 1
C          DO 92 NXX=1,150
C          DO 92 NTT = 1,20
C          FUTX(NXX,NTT) = 0.0
C          CONTINUE
C          DO 300 IX=1,MODES
C          ANAME(6) = AMODE(IX)
C          NCALL1 = IKNT(IX)
C          DO 55 K=2,NCALL1
C          DO 49 IXX = 1,150
C          DUMMY(IXX) = 0.0
C          FX = 0.0
C          MIL1 = NTOP(IX,K)
C          MIL2 = NM(IX,K)
C          MIL3 = MIL1 + MIL3
C          DO 51 NU = 1, MIL2
C          MIL4 = MIL1 + NU
C          AX = PHIMX(IX)*FUT(IX,NU,K)
C          FUTX(MIL4,K) = AX + FUTX(MIL4,K)
C          IF ( AX.LE.BX) AX = BX
C          PLOTTING FACTOR ADJUSTMENT
C          CX=SC*ALOG10(AX) + 15.0
C          IF ( CX.GT.9.0) CX = 9.0
C          IF ( CX.LT.BX) CX=0.0
C          DUMMY(MIL4) = CX
C          CONTINUE
C          LABEL = LTAB(K)

```



```

IF ( K.EQ.2) MODC = 1
IF ( (K.GT.2).AND.(LAST.EQ.C)) MODC = 2
IF ( (K.EQ.NCALL1) MODC = 3
LAST = K
CALL DRAW(NPNTS,DUMMY,U,MODC,ITYPE,LABEL,ANAME,EXSCAL,YSCALE,
1 IXUP,IYRT,MODX,MODY,IWIDE,IHIGH,IGRID,LAST)
IF (LAST.EQ.1) GO TO 105
IF (LAST.EQ.3) GO TO 105
CONTINUE
55 CONTINUE EQ.1) WRITE (6,106)
105 IF (LAST.EQ.3) WRITE (6,107)
300 CONTINUE
LAST = 0
C SECOND PART OF THE DRAW PLOT ROUTINE TO PLOT THE SUM OF THE TOTAL
C HARMONIC EXPANSION OF THE NEUTRON DENSITY FUNCTIONS
NKNT = IKNT(1)
MX=MCDES-1
IF (MX.EQ.0) MX=1
RTITLE(6) = BMODE(MX)
DO 310 NXT = 2,NKNT
DO 309 NTNX =1,NPNTS
AX = FUTX(NTNX,NXT)
IF (AX.LI.BX) AX = BX
DUMY(NTNX) = SC*ALOG10(AX) + 9.C
CONTINUE
309 LABEL = LTAB(NXT)
IF (NXT.EQ.2) MODC = 1
IF ( (NXT.GT.2).AND.(LAST.EQ.C)) MODC = 2
IF (NXT.EQ.NKNT) MODC = 3
LAST = 0
C PROTECTION STEPS TO PREVENT EXCESSIVE PLOTTING ON THE CAL COMP
C TO OVERRIDE THIS PROTECTION,THE ROUTINE MUST BE REWRITTEN
NFLAG=NFLAG+1
IF (NFLAG.EQ.4) GO TO 5500
CALL DRAW(NPNTS,DUMMY,U,MODC,ITYPE,LABEL,RTITLE,EXSCAL,YSCALE,
1 IXUP,IYRT,MODX,MODY,IWIDE,IHIGH,IGRID,LAST)
IF (LAST.EQ.1) GO TO 320
IF (LAST.EQ.3) GO TO 320
CONTINUE
310 CONTINUE EQ.1) WRITE (6,106)
320 IF (LAST.EQ.3) WRITE (6,107)
106 IF (LAST.EQ.3) WRITE (6,107)
107 FORMAT (//,1GX,'LAST = 1',3X,'RECHECK DATA')
5001 FORMAT (//,1GX,'LAST=3',3X,'RECHECK DATA')
5500 CONTINUE
CONTINUE
RETURN
END

```



```

SUBROUTINE READ11(LASTN)
REVISED 5 MAY 1971
THIS SUBROUTINE HAS NOT BEEN COMPLETELY TESTED AND DEBUGGED

READ11 INPUTS THE SAME TEST DATA AS ROUTINE READ10
BY PROCESSING NINE(9) TRACK MAGNETIC TAPE
INITIALLY AT A DENSITY OF 800 BPI
IDENTICAL INPUT/OUTPUT MESSAGES AND DATA ARE FURNISHED
AS BY READ10

REAL MTITLE (54)
DIMENSION NNRUN(20), NNT(20)

COMMON/COM1/ XBX(100), U(150), E(150), TIM(20), DUMMY(150), TXBX(100)
COMMON/COM4/ SOH,SA,S1,S2,S3,S4,SZERO,BUCKLE,EMAX,EMIN,DIFLEN,X
COMMON/COM7/ S01,S02,S03,S04,AK11,AK12,AK13,AK14,AN3A,AN3B,
1 AN3C,AN34,AN4,ANB4,AN2X
COMMON/COM2/ NT(20),NPROB(54),NANS(6),L(20),LTAB(20)
COMMON/COM5/ MODES,IS,IDATA,N,NVIR,NI,NF,NCALL,NOUT,IBX,
1 INSTAT,MODE
COMMON/COM3/ FXT(100,20),FUTX(150,20),BUCK(6),FUT(6,71,20)
COMMON/COM6/ NTOP(6,20),NM(6,20),IKNT(6)

SET DEFAULT OPTION FOR LASTN = 0
LASTN=0
DO 8 ISET=1,MODES
READ (4,20,END=8,ERR=9) (MTITLE(IA),IA=1,54)
READ (4,21,END=8,ERR=9) N,NVIR,NI,NF,BUCK(ISET)
IF (ISET.GT.1) GO TO 1
READ (4,22) (U(IA),IA=1,NVIR)
READ (4,22) (E(IA),IA=1,NVIR)
1 NCOUNT = 0
2 NCALL1 = 0
IF (NCALL1 = NCALL1 + 1 GO TO 18
IF (NCALL1.GE.21) GO TO 18
READ (4,23,END=8,ERR=9) IRUN,TIM(NCALL1),NNT(NCALL1),
1 NTRUN(ISET,NCALL1),NM(ISET,NCALL1)
NNRUN(NCALL1) = IRUN
MILLI = NM(ISET,NCALL1) + 1
READ (4,24,END=8,ERR=9) (FUT(ISET,IB,NCALL1),IB=1,MILL)
IF (IRUN.EQ.1) GO TO 2
END OF TAPE READ IN LOOP FOR NORMAL TIME STEP INPUT DATA

DEFAULT OPTION IF MORE THAN 20 DATA SETS IS PRODUCED PER
TIME STEP IN MOD 5
WRITE (6,27)
18 WKNT = 0

```



```

19 READ (4,23,ERR=9) IXRUN,XTIM,NXNT,NXNM
   NKNT = NKNT + 1
   MIL2 = NXNM + 1
   READ (4,24,END=8,ERR=9) (DUMMY(IA),IA=1,MIL2)
   WRITE (6,28) IXRUN,XTIM,NXNT,NXTOP,NXNM
   IF(NKNT.GT.20) GO TO 9
   PROTECTICN STEP TO GET OUT OF INFINITE TAPE READ LOOP
   IF (IXRUN.EQ.1) GO TO 19
   WRITE(6,29) NKNT
   END OF TAPE READ DEFAULT OPTIONS
   DATA PRINT OUT ROUTINE TO INFORM USER OF DATA SETS INPUT
   IF (ISET.GT.1) GO TO 4
   DO 3 IX=1,NCALL1
   NT(IX) = NNT(IX)
   3 WRITE (6,25) (MTITLE(IA),IA=1,54)
   4 WRITE (6,26) N,NVIR,NI,NF,BUCK(ISET)
   IF (ISET.GT.1) GO TO 6
   NPX=(NVIR/2) + 1
   DO 5 IK1=1,NPX
   IKIP = NPX + IK1
   5 WRITE (6,30) IK1,U(IK1),E(IK1),IKIP,U(IKIP),E(IKIP)
   6 WRITE (6,31)
   DO 7 IX=1,NCALL1
   7 WRITE (6,29) IX,TIM(IX),NNT(IX),NTOPI(ISET,IX),NM(ISET,IX)
   8 CONTINUE
   8 WRITE (6,32)
   9 CONTINUE
   9 LASTIN=1
   50 GO TO 50
   20 FORMAT (18A4)
   21 FORMAT (4I5,E15.8)
   22 FORMAT (8E10.3)
   23 FORMAT (I5,E15.8,3I5)
   24 FORMAT (8E10.3)
   25 FORMAT (/,10X,NF,18A4,/)
   26 18X,I5,/,10X,NF,10X,NVIR=' ',6X,I5,/,10X,NI=' ',
   27 10X,I5,/,10X,NVIR=' ',3X,F10.5)
   27 10X,INPUT TIME DATA FROM MOD-5,/,
   28 10X,EXCEEDS 20 DATA SETS,/)
   28 FORMAT (/,10X,I5,E10.3,3I5)
   29 FORMAT (/,10X,/,TOTAL EXCESS DATA SETS =',2X,I5)
   30 FORMAT (10X,2(I3,5X,2(IPEIC.3,5X))})
   31 10X,INPUT DATA FROM MOD 5,/,/,
   31 10X,INPUT TIME,3X,NT,3X,NTOPI,2X,NM,/)
   32 13X,SET,3X,I5,/,3X,/,***,3X,/,END OF INPUT DATA SET',/,
   32 10X,3(8X,/,10X,/,END OF TAPE READ CYCLE')
   34 FORMAT (/,10X,/,END OF TAPE READ CYCLE')

```


RETURN
END

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(Security classification of title, body of abstract and indexing annotation must be entered when the overall report is classified)

1. ORIGINATING ACTIVITY (Corporate author)		2a. REPORT SECURITY CLASSIFICATION	
Naval Postgraduate School Monterey, California 93940		UNCLASSIFIED	
		2b. GROUP	
3. REPORT TITLE			
SPACE DEPENDENT MODEL FOR THE SLOWING DOWN OF FAST NEUTRONS			
4. DESCRIPTIVE NOTES (Type of report and, inclusive dates)			
Master's Thesis; June 1971			
5. AUTHOR(S) (First name, middle initial, last name)			
Thomas M. Mills			
6. REPORT DATE		7a. TOTAL NO. OF PAGES	7b. NO. OF REFS
June 1971		151	15
8a. CONTRACT OR GRANT NO.		9a. ORIGINATOR'S REPORT NUMBER(S)	
b. PROJECT NO.			
c.		9b. OTHER REPORT NO(S) (Any other numbers that may be assigned this report)	
d.			
10. DISTRIBUTION STATEMENT			
Approved for public release; distribution unlimited.			
11. SUPPLEMENTARY NOTES		12. SPONSORING MILITARY ACTIVITY	
		Naval Postgraduate School Monterey, California 93940	
13. ABSTRACT			
<p>The slowing down of fast neutrons was analyzed by a multi-group method of discrete time and energy states coupled with a spatial harmonic expansion method to determine the neutron density in a homogeneous, isotropically scattering slab. Five neutron source geometries were studied for both a fissioning and a non-fissioning system.</p> <p>Numerical results were obtained for the neutron flux, mean neutron energy and the neutron spectra for the one dimensional system using a harmonic mode expansion of up to six terms to determine the time-energy-space dependence.</p>			

KEY WORDS	LINK A		LINK B		LINK C	
	ROLE	WT	ROLE	WT	ROLE	WT
Neutron Diffusion						
Boltzman Transport						
Neutron Slowing Down						
Multi-Group						

Neutron Diffusion

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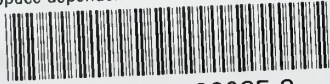
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